

Review: Hybrid elements for incompressible materials

- Re-derive equations for fully incompressible linear elasticity ab-initio

Field equations:

$$\sigma_{ij} = S_{ij} + \frac{P}{3} \delta_{ij} \quad \varepsilon_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \quad \frac{1}{3} \varepsilon_{kkk} = \frac{1}{3} \frac{\partial u_k}{\partial x_k} = 0$$

$$S_{ij} = 2\mu \varepsilon_{ij} \quad \frac{\partial \sigma_{ij}}{\partial x_i} = 0 \quad u_i = u_i^* \quad x_i \in S_1 \quad \sigma_{ij} n_i = t_j^* \quad x_i \in S_2$$

Obtain weak form:

$$\int_V \left(2\mu \frac{\partial \varepsilon_{ij}}{\partial x_j} + \frac{1}{3} \frac{\partial p}{\partial x_i} \right) \eta_i dV = 0 \quad \int_V \frac{1}{3} \frac{\partial u_k}{\partial x_k} q dV = 0$$

Integrate by parts:

$$\int_V \left(2\mu \varepsilon_{ij} \frac{\partial \eta_i}{\partial x_j} + \frac{p}{3} \frac{\partial \eta_i}{\partial x_i} \right) dV - \int_V b_i \eta_i dV - \int_{S_2} t_i^* \eta_i dA = 0 \quad \int_V \frac{1}{3} \frac{\partial u_k}{\partial x_k} q dV = 0$$

FE interpolation

$$u_i = N^a(\mathbf{x}) u_i^a \quad \eta_i = N^a(\mathbf{x}) \eta_i^a$$

$$p = M^a(\mathbf{x}) p^a \quad q = M^a(\mathbf{x}) q^a$$

$M^a(\mathbf{x})$ are interpolation functions for pressure
 Must be one order lower than interpolation for displacement
 Need not be continuous across element boundaries

Insert interpolation into weak form to get FE eqs

Define $C_{ijke}^{DEV} = \mu (\delta_{ie} \delta_{jk} + \delta_{ik} \delta_{je})$

$$\left\{ \int_{\mathbb{R}} \underbrace{C_{ijke}^{DEV} \frac{\partial N^b}{\partial x_e} \frac{\partial N^a}{\partial x_j} u_k^b}_{K_{abbb}^{uu}} + \underbrace{\frac{1}{3} \frac{\partial N^a}{\partial x_i} m_b^a p^b}_{K_{aib}^{up}} \right\} dV - \int_{\Gamma_i} t_i^* N^a \Big|_{\Gamma_i} \eta_i^a = 0$$

\(\forall\) admis η_i^a

$$\left\{ \int_{\mathbb{R}} \underbrace{\frac{1}{3} \frac{\partial N^b}{\partial x_i} m^a u_i^b}_{K_{aib}^{pu}} \right\} q^a = 0 \quad \forall \text{ admis } q^b$$

FE equations follow

$$K_{aibb}^{uu} U_k^b + K_{aib}^{up} p^b = f_i^a$$

$$K_{aib}^{pu} U_i^b = 0$$

Assembly is done element-by-element as before

$$[K] \begin{bmatrix} \underline{u} \\ \underline{p} \end{bmatrix} = \begin{bmatrix} \underline{f} \\ 0 \end{bmatrix}$$

Focus on element level operations

$$[k^{el}] = \begin{bmatrix} K^{uu} & K^{up} \\ K^{pu} & 0 \end{bmatrix} \begin{bmatrix} \underline{u} \\ \underline{p} \end{bmatrix}$$

$$[K^{uu}] = \int_{\Omega_{el}} [B]^T [D^{dev}] [B] dV$$

$$[K^{up}] = \int_{\Omega_{el}} \frac{1}{3} [B^T] [M] dV$$

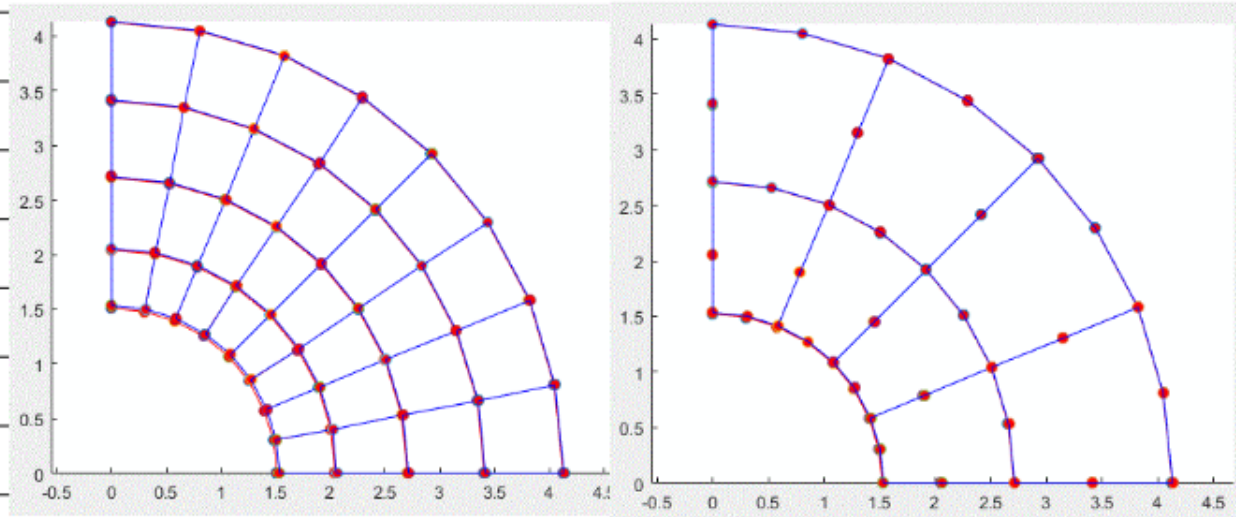
$$[K^{pu}] = [K^{up}]^T \quad [M] \text{ maps } \underline{p} \text{ onto } \underline{\sigma} = \underline{\sigma} = [M] \underline{p}$$

Compute with
full Gauss quadrature

Test:

Linear

Quadratic



Advantages : Works for fully incompressible materials

Disadvantages: Can't handle compressible solids
 Solve for more global DOF
 Zeros on diagonal can be a problem for some solvers

6.6 Augmented Lagrangean hybrid elements

Goal: Develop hybrid scheme for compressible materials

Approach: Solve separately for pressure-volume & deviatoric relations

$$\sigma_{ij} = S_{ij} + p \delta_{ij} \quad p = \sigma_{kk} / 3$$

$$S_{ij} = C_{ijkl}^{DEV} E_{kl}$$

$$p = K E_{kk} = K \frac{dU_k}{dX_k}$$

$$K = \text{Bulk modulus} = \frac{E}{3(1-2\nu)}$$

New equation

$$C_{ijkl} = \frac{E}{2(1+\nu)} (\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk}) + \frac{E\nu}{(1+\nu)(1-2\nu)} \delta_{ij}\delta_{kl}$$

$$\Rightarrow C_{ijkl}^{DEV} = C_{ijkl} - \frac{E}{3(1-2\nu)} \delta_{ij}\delta_{kl} = \frac{E}{2(1+\nu)} (\delta_{ik}\delta_{jl} + \delta_{il}\delta_{jk}) - \frac{E}{3(1+\nu)} \delta_{ij}\delta_{kl}$$

Interpolated weak form of governing equations

$$\left\{ \int_{\mathcal{R}} \underbrace{C_{ijkl}^{DEV} \frac{\partial N^b}{\partial x_l} \frac{\partial N^a}{\partial x_j} U_k^b}_{K^{uu}} + \underbrace{\frac{\partial N^a}{\partial x_i} M^b p^b}_{K^{up}} \right\} \eta_i^a = 0$$

Also enforce $p = \kappa \frac{\partial U_k}{\partial x_k}$ in weak form $\left(\frac{\partial U_k}{\partial x_k} - \frac{p}{\kappa} \right) = 0$

$$\left\{ \int_{\mathcal{R}} \underbrace{\frac{\partial N^b}{\partial x_i} U_i^b M^a}_{K^{pu}} - \underbrace{M^a M^b p^b}_{K^{pp}} \right\} q^a = 0$$

Assembly of global system is the same as for hybrid element

Focus on element computation

$$[k^{el}] = \begin{bmatrix} k^{uu} & k^{up} \\ k^{pu} & k^{pp} \end{bmatrix} \begin{bmatrix} \underline{u} \\ \underline{p} \end{bmatrix}$$

$$[k^{uu}] = \int_{\Omega_{el}} [B]^T [D^{Der}] [B] dV$$

$$[k^{up}] = [k^{pu}]^T = \int_{\Omega_{el}} [B]^T [M] dV$$

$$[k^{pp}] = \int_{\Omega_{el}} -\frac{1}{K} \underline{m} \otimes \underline{m} dV$$

Compute
with full
integration

$\underline{m} = [1]$ for linear element

$\underline{m} = [m^1 \ m^2 \ m^3 \ m^4]$ for quadratic quad

This formulation works for all materials

Note that as long as $K \neq \infty$ we can eliminate p inside each element

$$K^{pu} \underline{u} + K^{pp} p = \underline{0}$$

$$\Rightarrow p = -K^{pp-1} K^{pu} \underline{u}$$

$$K^{uu} \underline{u} + K^{up} p = \underline{f}$$

$$\Rightarrow \underbrace{\left[K^{uu} - K^{up} K^{pp-1} K^{pu} \right]}_{\text{Reduced } [K^{e1}]} \underline{u} = \underline{f}$$

In practice p is usually left in global system

This version is used as hybrid element in ABAQUS