### **Course Outline**

- MATLAB tutorial
- 2. Motion of systems that can be idealized as particles
  - Description of motion; Newton's laws;
  - Calculating forces required to induce prescribed motion;
  - Deriving and solving equations of motion
- 3. Conservation laws for systems of particles
  - Work, power and energy;
  - Linear impulse and momentum
  - Angular momentum
- 4. Vibrations
  - Characteristics of vibrations; vibration of free 1 DOF systems
  - Vibration of damped 1 DOF systems
  - Forced Vibrations
- 5. Motion of systems that can be idealized as rigid bodies
  - Description of rotational motion; kinematics formulas
  - Dynamics formulas for rigid bodies; calculating moments of inertia
  - Motion of systems of rigid bodies
  - Energy and momentum for rigid bodies

Exam topics

# **Particle Dynamics: Concept Checklist**

- Understand the concept of an 'inertial frame'
- Be able to idealize an engineering design as a set of particles, and know when this idealization will give accurate results
- Describe the motion of a system of particles (eg components in a fixed coordinate system; components in a polar coordinate system, etc)
- Be able to differentiate position vectors (with proper use of the chain rule!) to determine velocity and acceleration; and be able to integrate acceleration or velocity to determine position vector.
- Be able to describe motion in normal-tangential and polar coordinates (eg be able to write down vector components of velocity and acceleration in terms of speed, radius of curvature of path, or coordinates in the cylindrical-polar system).
- Be able to convert between Cartesian to normal-tangential or polar coordinate descriptions of motion
- Be able to draw a correct free body diagram showing forces acting on system idealized as particles
- Be able to write down Newton's laws of motion in rectangular, normal-tangential, and polar coordinate systems
- Be able to obtain an additional moment balance equation for a rigid body moving without rotation or rotating about a fixed axis at constant rate.
- Be able to use Newton's laws of motion to solve for unknown accelerations or forces in a system of particles
- Use Newton's laws of motion to derive differential equations governing the motion of a system of particles
- Be able to re-write second order differential equations as a pair of first-order differential equations in a form that MATLAB can solve

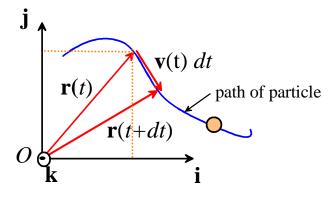
Inertial frame – non accelerating, non rotating reference frame Particle – point mass at some position in space

Position Vector 
$$\mathbf{r}(t) = x(t)\mathbf{i} + y(t)\mathbf{j} + z(t)\mathbf{k}$$

Velocity Vector 
$$\mathbf{v}(t) = v_x(t)\mathbf{i} + v_y(t)\mathbf{j} + v_z(t)\mathbf{k}$$
  

$$= \frac{d}{dt}(x\mathbf{i} + y\mathbf{j} + z\mathbf{k}) = \frac{dx}{dt}\mathbf{i} + \frac{dy}{dt}\mathbf{j} + \frac{dz}{dt}\mathbf{k}$$

$$\Rightarrow v_x(t) = \frac{dx}{dt} \qquad v_y(t) = \frac{dy}{dt} \qquad v_z(t) = \frac{dz}{dt}$$



- Direction of velocity vector is parallel to path
- Magnitude of velocity vector is distance traveled / time

#### **Acceleration Vector**

$$\mathbf{a}(t) = a_x(t)\mathbf{i} + a_y(t)\mathbf{j} + a_z(t)\mathbf{k} = \frac{d}{dt}\left(v_x\mathbf{i} + v_y\mathbf{j} + v_z\mathbf{k}\right) = \frac{dv_x}{dt}\mathbf{i} + \frac{dv_y}{dt}\mathbf{j} + \frac{dv_z}{dt}\mathbf{k}$$

$$\Rightarrow a_x(t) = \frac{dv_x}{dt} = \frac{d^2x}{dt^2} \quad a_y(t) = \frac{dv_y}{dt} = \frac{d^2y}{dt^2} \quad a_z(t) = \frac{dv_z}{dt} = \frac{d^2z}{dt^2}$$

Straight line motion with constant acceleration

$$\mathbf{r} = \left[ X_0 + V_0 t + \frac{1}{2} a t^2 \right] \mathbf{i} \qquad \mathbf{v} = (V_0 + a t) \mathbf{i} \qquad \mathbf{a} = a \mathbf{i}$$

Time/velocity/position dependent acceleration – use calculus

$$\mathbf{r} = \left(X_0 + \int_0^t v(t)dt\right)\mathbf{i} \qquad \mathbf{v} = \left(V_0 + \int_0^t a(t)dt\right)\mathbf{i}$$

$$v(t) \qquad x(t) \qquad x(t)$$

$$\int_0^t vdv = \int_0^t a(x)dx$$

$$V_0 \qquad 0$$

$$a = \frac{dv}{dt} = \frac{g(t)}{f(v)} \Rightarrow \int_{V_0}^v f(v)dv = \int_0^t g(t)dt$$

$$v = \frac{dx}{dt} = \frac{g(t)}{f(v)} \Rightarrow \int_{V_0}^{x(t)} f(x)dv = \int_0^t v(t)dt$$

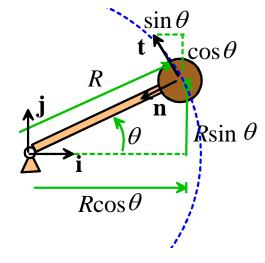
Circular Motion at const speed

$$\theta = \omega t \quad s = R\theta \quad V = \omega R$$

$$\mathbf{r} = R\left(\cos\theta \mathbf{i} + \sin\theta \mathbf{j}\right)$$

$$\mathbf{v} = \omega R\left(-\sin\theta \mathbf{i} + \cos\theta \mathbf{j}\right) = V\mathbf{t}$$

$$\mathbf{a} = -\omega^2 R(\cos\theta \mathbf{i} + \sin\theta \mathbf{j}) = \omega^2 R\mathbf{n} = \frac{V^2}{R}\mathbf{n}$$



General circular motion

$$\omega = d\theta / dt \quad \alpha = d\omega / dt = d^{2}\theta / dt^{2}$$

$$s = R\theta \quad V = ds / dt = R\omega$$

$$\mathbf{r} = R(\cos\theta \mathbf{i} + \sin\theta \mathbf{j})$$

$$\mathbf{v} = \omega R(-\sin\theta \mathbf{i} + \cos\theta \mathbf{j}) = V\mathbf{t}$$

$$\mathbf{a} = R\alpha(-\sin\theta \mathbf{i} + \cos\theta \mathbf{j}) - R\omega^{2}(\cos\theta \mathbf{i} + \sin\theta \mathbf{j})$$

$$= \alpha R\mathbf{t} + \omega^{2}R\mathbf{n} = \frac{dV}{dt}\mathbf{t} + \frac{V^{2}}{R}\mathbf{n}$$

Arbitrary path

$$\mathbf{v} = V\mathbf{t}$$

$$\mathbf{a} = \frac{dV}{dt}\mathbf{t} + \frac{V^2}{R}\mathbf{n}$$

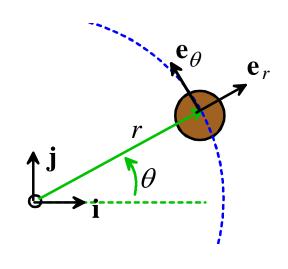
$$\mathbf{r} = x(\lambda)\mathbf{i} + y(\lambda)\mathbf{j}$$

$$\frac{1}{R} = \frac{\left| \frac{dx}{d\lambda} \frac{d^2y}{d\lambda^2} - \frac{dy}{d\lambda} \frac{d^2x}{d\lambda^2} \right|}{\left\{ \left( \frac{dx}{d\lambda} \right)^2 + \left( \frac{dy}{d\lambda} \right)^2 \right\}^{3/2}}$$

Polar Coordinates

$$\mathbf{v} = \frac{dr}{dt}\mathbf{e}_r + r\frac{d\theta}{dt}\mathbf{e}_{\theta}$$

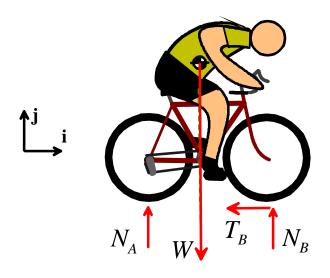
$$\mathbf{a} = \left(\frac{d^2r}{dt^2} - r\left(\frac{d\theta}{dt}\right)^2\right)\mathbf{e}_r + \left(r\frac{d^2\theta}{dt^2} + 2\frac{dr}{dt}\frac{d\theta}{dt}\right)\mathbf{e}_{\theta}$$



### **Newton's laws**

• For a particle  $\mathbf{F} = m\mathbf{a}$ 

 For a rigid body in motion without rotation, or a particle on a massless frame



$$\mathbf{M}_{c} = \mathbf{0}$$

You MUST take moments about center of mass

# Calculating forces required to cause prescribed motion of a particle

- Idealize system
- Free body diagram
- Kinematics
- F=ma for each particle.
- $\mathbf{M}_c = \mathbf{0}$  (for rigid bodies or frames only)
- Solve for unknown forces or accelerations

# **Deriving Equations of Motion for particles**

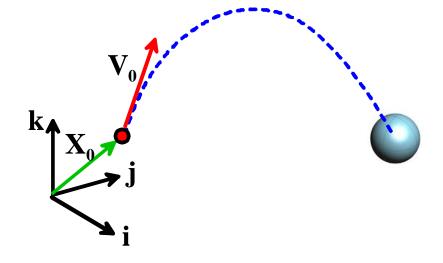
- 1. Idealize system
- 2. Introduce variables to describe motion (often *x,y* coords, but we will see other examples)
- 3. Write down **r**, differentiate to get **a**
- 4. Draw FBD
- 5. F = ma
- 6. If necessary, eliminate reaction forces
- 7. Result will be differential equations for coords defined in (2), e.g.  $m \frac{d^2x}{dt^2} + \lambda \frac{dx}{dt} + kx = kY_0 \sin \omega t$
- 8. Identify initial conditions, and solve ODE

# Motion of a projectile

$$\mathbf{r} = X_0 \mathbf{i} + Y_0 \mathbf{j} + Z_0 \mathbf{k}$$

$$\frac{d\mathbf{r}}{dt} = V_{x0} \mathbf{i} + V_{y0} \mathbf{j} + V_{z0} \mathbf{k}$$

$$t = 0$$



$$\mathbf{r} = (X_0 + V_{x0}t)\mathbf{i} + (Y_0 + V_{y0}t)\mathbf{j} + (Z_0 + V_{z0}t - \frac{1}{2}gt^2)\mathbf{k}$$

$$\mathbf{v} = (V_{x0})\mathbf{i} + (V_{y0})\mathbf{j} + (V_{z0} - gt)\mathbf{k}$$

$$\mathbf{a} = -g\mathbf{k}$$

# Rearranging differential equations for MATLAB

• Example 
$$\frac{d^2y}{dt^2} + 2\zeta\omega_n \frac{dy}{dt} + \omega_n^2 y = 0$$

- Introduce v = dy / dt
- Then  $\frac{d}{dt} \begin{bmatrix} y \\ v \end{bmatrix} = \begin{bmatrix} v \\ -2\zeta\omega_n v \omega_n^2 y \end{bmatrix}$
- This has form  $\frac{d\mathbf{w}}{dt} = f(t, \mathbf{w})$   $\mathbf{w} = \begin{vmatrix} y \\ v \end{vmatrix}$

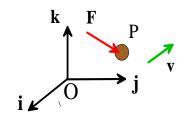
# **Conservation laws for particles: Concept Checklist**

- Know the definitions of power (or rate of work) of a force, and work done by a force
- Know the definition of kinetic energy of a particle
- Understand power-work-kinetic energy relations for a particle
- Be able to use work/power/kinetic energy to solve problems involving particle motion
- Be able to distinguish between conservative and non-conservative forces
- Be able to calculate the potential energy of a conservative force
- Be able to calculate the force associated with a potential energy function
- Know the work-energy relation for a system of particles; (energy conservation for a closed system)
- Use energy conservation to analyze motion of conservative systems of particles
- Know the definition of the linear impulse of a force
- Know the definition of linear momentum of a particle
- Understand the impulse-momentum (and force-momentum) relations for a particle
- Understand impulse-momentum relations for a system of particles (momentum conservation for a closed system)
- Be able to use impulse-momentum to analyze motion of particles and systems of particles
- Know the definition of restitution coefficient for a collision
- Predict changes in velocity of two colliding particles in 2D and 3D using momentum and the restitution formula
- Know the definition of angular impulse of a force
- Know the definition of angular momentum of a particle
- Understand the angular impulse-momentum relation
- Be able to use angular momentum to solve central force problems

# Work and Energy relations for particles

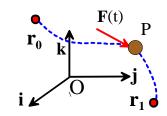
Rate of work done by a force (power developed by force)

$$P = \mathbf{F} \cdot \mathbf{v}$$



Total work done by a force

$$W = \int_{0}^{t_{1}} \mathbf{F} \cdot \mathbf{v} dt \quad W = \int_{\mathbf{r}_{0}}^{\mathbf{r}_{1}} \mathbf{F} \cdot d\mathbf{r}$$



Kinetic energy

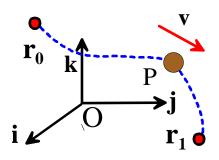
$$T = \frac{1}{2}m|\mathbf{v}|^2 = \frac{1}{2}m(v_x^2 + v_y^2 + v_z^2)$$

Power-kinetic energy relation

Work-kinetic energy relation

$$P = \frac{dT}{dt}$$

$$W = \int_{\mathbf{r}_{0}}^{\mathbf{r}_{1}} \mathbf{F} \cdot d\mathbf{r} = T - T_{0}$$

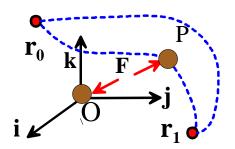


# **Potential energy**

Potential energy of a conservative force (pair)

$$V(\mathbf{r}) = -\int_{\mathbf{r}_0}^{\mathbf{r}} \mathbf{F} \cdot d\mathbf{r} + \text{constant}$$

$$\mathbf{F} = -\operatorname{grad}(V)$$



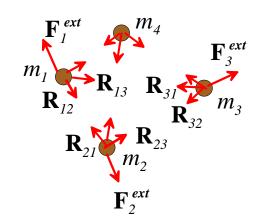
Type of force	Potential energy	
Gravity acting on a particle near earths surface	V = mgy	j h
Gravitational force exerted on mass <i>m</i> by mass <i>M</i> at the origin	$V = -\frac{GMm}{r}$	F m
Force exerted by a spring with stiffness $k$ and unstretched length $L_0$	$V = \frac{1}{2}k(r - L_0)^2$	$ \begin{array}{cccc} j & r & O \\ \uparrow & & & & & \\ \downarrow & & \\ \downarrow$
Force acting between two charged particles	$V = \frac{Q_1 Q_2}{4\pi\varepsilon r}$	$Q_1$ $\mathbf{j}$ $Q_2$ $\mathbf{F}$
Force exerted by one molecule of a noble gas (e.g. He, Ar, etc) on another (Lennard Jones potential). <i>a</i> is the equilibrium spacing between molecules, and <i>E</i> is the energy of the bond.	$E\left[\left(\frac{a}{r}\right)^{12} - 2\left(\frac{a}{r}\right)^{6}\right]$	j 2 F

# Energy relations for conservative systems subjected to external forces

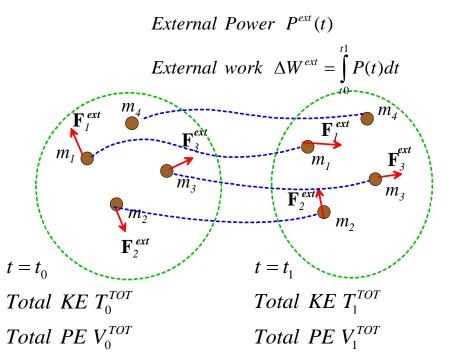
Internal Forces: (forces exerted by one part of the system on another)  $\mathbf{R}_{ij}$ 

**External Forces:** (any other forces)  $\mathbf{F}_{i}^{ext}$ 

System is conservative if all internal forces are conservative forces (or constraint forces)



#### **Energy relation for a conservative system**



$$\Delta W_{ext} = T_1^{TOT} + V_1^{TOT} - \left(T_0^{TOT} + V_0^{TOT}\right)$$

#### Special case – zero external work:

$$T_1^{TOT} + V_1^{TOT} = T_0^{TOT} + V_0^{TOT}$$

KE+PE = constant

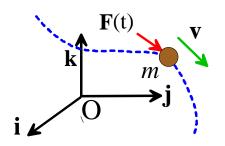
# Impulse-momentum for a single particle

### **Definitions**

Linear Impulse of a force  $\mathbf{I} = \int_{t}^{t_1} \mathbf{F}(t)dt$ 

Linear momentum of a particle

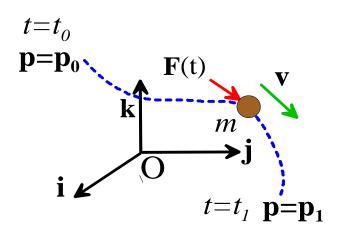
$$\mathbf{p} = m\mathbf{v}$$



### **Impulse-Momentum relations**

$$\mathbf{F} = \frac{d\mathbf{p}}{dt}$$

$$\mathbf{I} = \mathbf{p}_1 - \mathbf{p}_0$$

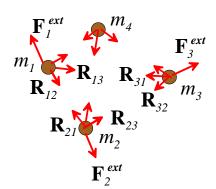


### Impulse-momentum for a system of particles

 $\mathbf{R}_{ij}$  Force exerted on particle i by particle j

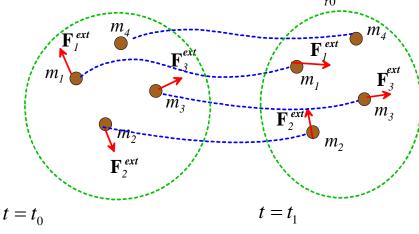
 $\mathbf{F}_{i}^{ext}$  External force on particle i

 $\mathbf{v}_i$  Velocity of particle i



Total External Force  $\mathbf{F}^{TOT}(t)$ 

Total External Impulse  $\mathbf{I}^{TOT} = \int_{0}^{1} \mathbf{F}^{TOT}(t) dt$ 



Total momentum  $\mathbf{p}_0^{TOT}$ 

Total momentum  $\mathbf{p}_1^{TOT}$ 

#### Impulse-momentum for the system:

$$\mathbf{F}^{TOT} = \frac{d\mathbf{p}^{TOT}}{dt}$$

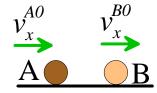
$$\mathbf{I}^{TOT} = \mathbf{p}_1^{TOT} - \mathbf{p}_0^{TOT}$$

#### **Special case – zero external impulse:**

$$\mathbf{p}_1^{TOT} = \mathbf{p}_0^{TOT}$$

(Linear momentum conserved)

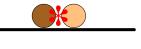
### Collisions



Momentum

$$m_A v_x^{A1} + m_B v_x^{B1} = m_A v_x^{A0} + m_B v_x^{B0}$$

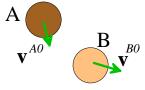
Restitution formula 
$$v^{B1} - v^{A1} = -e(v^{B0} - v^{A0})$$



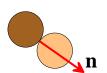
$$v_x^{BI}$$

$$v^{B1} = v^{B0} - \frac{m_A}{m_A + m_B} (1 + e) \left( v^{B0} - v^{A0} \right)$$

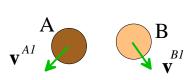
$$v^{A1} = v^{A0} + \frac{m_B}{m_A + m_B} (1 + e) \left( v^{B0} - v^{A0} \right)$$



By Momentum 
$$m_B \mathbf{v}^{B1} + m_A \mathbf{v}^{A1} = m_B \mathbf{v}^{B0} + m_A \mathbf{v}^{A0}$$



Restitution formula 
$$(\mathbf{v}^{B1} - \mathbf{v}^{A1}) = (\mathbf{v}^{B0} - \mathbf{v}^{A0}) - (1+e)[(\mathbf{v}^{B0} - \mathbf{v}^{A0}) \cdot \mathbf{n}]\mathbf{n}$$



$$\mathbf{v}^{A1} = \mathbf{v}^{A0} + \frac{m_B}{m_B + m_A} (1 + e) \left[ \left( \mathbf{v}^{B0} - \mathbf{v}^{A0} \right) \cdot \mathbf{n} \right] \mathbf{n}$$

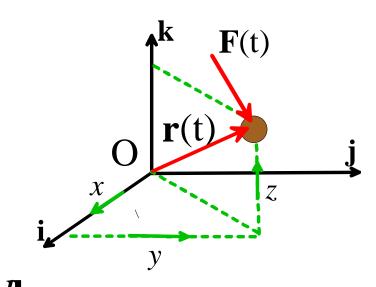
$$\mathbf{v}^{B1} = \mathbf{v}^{B0} - \frac{m_A}{m_B + m_A} (1 + e) \left[ \left( \mathbf{v}^{B0} - \mathbf{v}^{A0} \right) \cdot \mathbf{n} \right] \mathbf{n}$$

# **Angular Impulse-Momentum Equations for a Particle**

$$\mathbf{A} = \int_{t_0}^{t_1} \mathbf{r}(t) \times \mathbf{F}(t) dt$$

**Angular Momentum** 

$$\mathbf{h} = \mathbf{r} \times \mathbf{p} = \mathbf{r} \times m\mathbf{v}$$



Impulse-Momentum relations

$$\mathbf{r} \times \mathbf{F} = \frac{d\mathbf{h}}{dt}$$

$$\mathbf{A} = \mathbf{h}_1 - \mathbf{h}_0$$

$$\mathbf{A} = \mathbf{0} \Rightarrow \mathbf{h}_1 = \mathbf{h}_0$$

Angular momentum conserved

Useful for central force problems (when forces on a particle always act through a single point, eg planetary gravity)

# Free Vibrations – concept checklist

#### You should be able to:

- 1. Understand simple harmonic motion (amplitude, period, frequency, phase)
- 2. Identify # DOF (and hence # vibration modes) for a system
- 3. Understand (qualitatively) meaning of 'natural frequency' and 'Vibration mode' of a system
- 4. Calculate natural frequency of a 1DOF system (linear and nonlinear)
- 5. Write the EOM for simple spring-mass-damper systems by inspection
- 6. Understand natural frequency, damped natural frequency, and 'Damping factor' for a dissipative 1DOF vibrating system
- 7. Know formulas for nat freq, damped nat freq and 'damping factor' for spring-mass system in terms of k,m,c
- 8. Understand underdamped, critically damped, and overdamped motion of a dissipative 1DOF vibrating system
- 9. Be able to determine damping factor and natural frequency from a measured free vibration response
- 10. Be able to predict motion of a freely vibrating 1DOF system given its initial velocity and position, and apply this to design-type problems

# Vibrations and simple harmonic motion

### **Typical vibration response**

Period, frequency, angular frequency amplitude

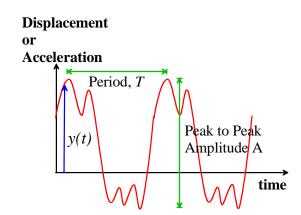
### **Simple Harmonic Motion**

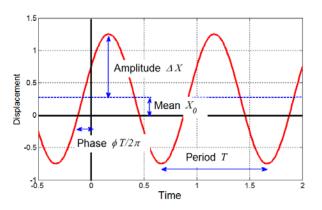
$$x(t) = X_0 + \Delta X \sin(\omega t + \phi)$$

$$v(t) = \Delta V \cos(\omega t + \phi)$$

$$a(t) = -\Delta A \sin(\omega t + \phi)$$

$$\Delta V = \omega \Delta X \qquad \Delta A = \omega \Delta V$$



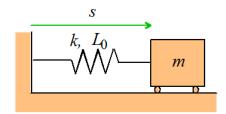


# Vibration of 1DOF conservative systems

#### **Harmonic Oscillator**

Derive EOM (F=ma) 
$$\frac{m}{k} \frac{d^2s}{dt^2} + s = L_0$$

**Canonical Vibration Problem:** The spring mass system is released with velocity  $v_0$  from position  $s_0$  at time t=0 . Find s(t).



#### Compare with 'standard' differential equation

Equation 
$$\frac{1}{\omega_n^2} \frac{d^2x}{dt^2} + x = C \quad \text{Initial Conditions} \quad x = x_0 \quad \frac{dx}{dt} = v_0 \quad t = 0$$
Solution 
$$\begin{aligned} x &= C + X_0 \sin(\omega_n t + \phi) \\ X_0 &= \sqrt{(x_0 - C)^2 + v_0^2 / \omega_n^2} \quad \phi = \tan^{-1} \left( \frac{(x_0 - C)\omega_n}{v_0} \right) \end{aligned}$$
Or 
$$x(t) = C + (x_0 - C)\cos\omega_n t + \frac{v_0}{\omega} \sin\omega_n t$$

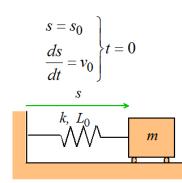
$$x = s C = L_0 x_0 = s_0$$

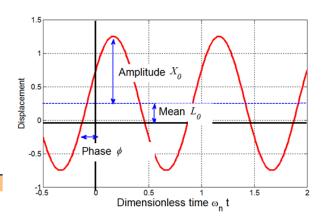
$$\frac{1}{\omega_n^2} = \frac{m}{k}$$

#### Solution

$$s(t) = L_0 + \sqrt{(s_0 - L_0)^2 + v_0^2 / \omega_n^2} \sin(\omega_n t + \phi)$$

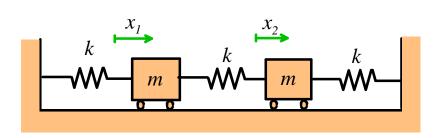
Natural Frequency 
$$\omega_n = \sqrt{\frac{k}{m}}$$

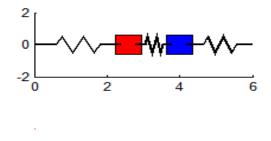


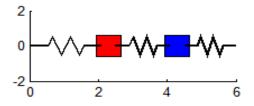


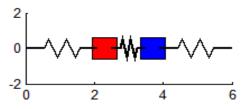
# Vibration modes and natural frequencies

- •Vibration modes: special initial deflections that cause entire system to vibrate harmonically
- •Natural Frequencies are the corresponding vibration frequencies







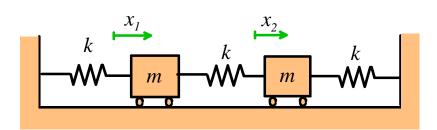


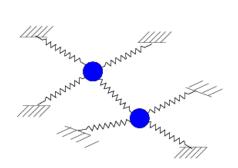
# Number of DOF (and vibration modes)

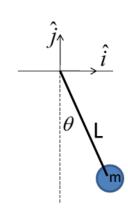
In 2D: # DOF =  $2^*\#$  particles +  $3^*\#$  rigid bodies - # constraints In 3D: # DOF =  $3^*\#$  particles +  $6^*\#$  rigid bodies - # constraints

Expected # vibration modes = # DOF - # rigid body modes

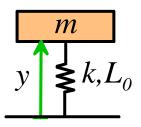
A 'rigid body mode' is steady rotation or translation of the entire system at constant speed. The maximum number of 'rigid body' modes (in 3D) is 6; in 2D it is 3. Usually only things like a vehicle or a molecule, which can move around freely, have rigid body modes.







# Calculating nat freqs for 1DOF systems – the basics

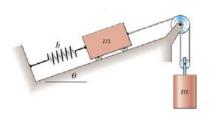


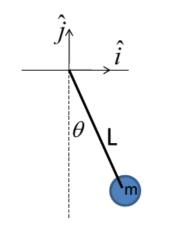
EOM for small vibration of any 1DOF undamped system has form

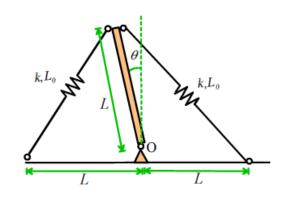
$$\frac{1}{\omega_n^2} \frac{d^2 y}{dt^2} + y = C$$

 $\omega_n$  is the natural frequency

- 1. Get EOM (F=ma or energy)
- 2. Linearize (sometimes)
- 3. Arrange in standard form
- 4. Read off nat freq.







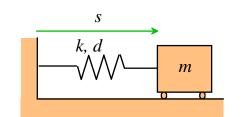
### Tricks for calculating natural frequencies of 1DOF undamped systems

Using energy conservation to find EOM

$$KE + PE = \frac{1}{2}m\left(\frac{ds}{dt}\right)^{2} + \frac{1}{2}k(s - L_{0})^{2} = const$$

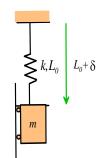
$$\Rightarrow \frac{d}{dt}(KE + PE) = m\left(\frac{ds}{dt}\right)\frac{d^{2}s}{dt^{2}} + k(s - L_{0})\frac{ds}{dt} = 0$$

$$\Rightarrow m\frac{d^{2}s}{dt^{2}} + ks = kL_{0}$$



Nat freq is related to static deflection

$$\omega_n = \sqrt{\frac{g}{\delta}}$$



### **Linearizing EOM**

Sometimes EOM has form

$$\frac{d^2y}{dt^2} + f(y) = C$$

 $\frac{\hat{j}}{\theta} \downarrow \frac{\hat{i}}{L}$ 

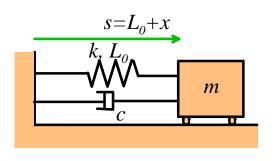
We cant solve this in general... Instead, assume y is small

$$m\frac{d^{2}y}{dt^{2}} + f(0) + \frac{df}{dy}\bigg|_{y=0} y + \dots = C$$

$$m\frac{d^{2}y}{dt^{2}} + \frac{df}{dy}\bigg|_{y=0} y = C - f(0)$$

There are short-cuts to doing the Taylor expansion

# Writing down EOM for spring-mass-damper systems

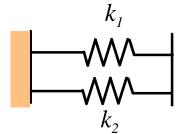


Commit this to memory! (or be able to derive it...)

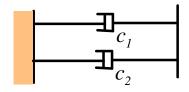
$$\mathbf{F} = m\mathbf{a} \Rightarrow \frac{d^2x}{dt^2} + \frac{c}{m}\frac{dx}{dt} + \frac{k}{m}x = 0$$

$$\frac{d^2x}{dt^2} + 2\zeta\omega_n \frac{dx}{dt} + \omega_n^2 x = 0 \qquad \omega_n = \sqrt{\frac{k}{m}} \qquad \zeta = \frac{c}{2\sqrt{km}}$$

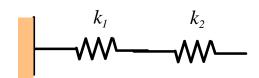
x(t) is the 'dynamic variable' (deflection from static equilibrium)



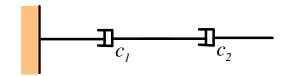
Parallel: stiffness  $k = k_1 + k_2$ 



Parallel: coefficient  $c = c_1 + c_2$ 



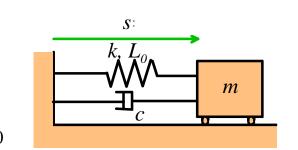
Series: stiffness  $\frac{1}{k} = \frac{1}{k_1} + \frac{1}{k_2}$ 



Parallel: coefficient  $\frac{1}{c} = \frac{1}{c_1} + \frac{1}{c_2}$ 

### Canonical damped vibration problem

**EOM** 
$$m\frac{d^2s}{dt^2} + c\frac{ds}{dt} + ks = kL_0$$
 with  $s = s_0$   $\frac{ds}{dt} = v_0$   $t = 0$ 



Standard Form 
$$\frac{1}{\omega_n^2} \frac{d^2x}{dt^2} + \frac{2\varsigma}{\omega_n} \frac{dx}{dt} + x = C$$
  $x = x_0$   $\frac{dx}{dt} = v_0$   $t = 0$ 

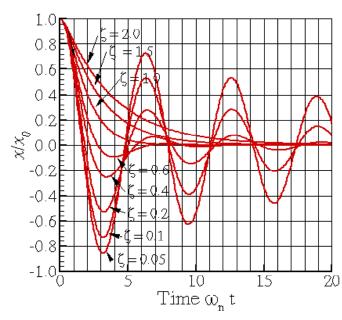
$$s \equiv x$$
  $\omega_n = \sqrt{\frac{k}{m}}$   $\zeta = \frac{c}{2\sqrt{km}}$   $C = L_0$   $x_0 \equiv s_0$ 

$$\omega_d = \omega_n \sqrt{1 - \zeta^2}$$

Overdamped  $\varsigma > 1$ 

Critically Damped  $\varsigma = 1$ 

Underdamped  $\varsigma < 1$ 

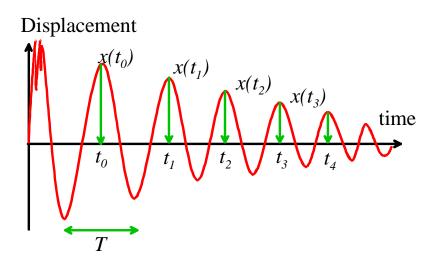


Overdamped 
$$\zeta > 1$$
  $x(t) = C + \exp(-\varsigma \omega_n t) \left\{ \frac{v_0 + (\varsigma \omega_n + \omega_d)(x_0 - C)}{2\omega_d} \exp(\omega_d t) - \frac{v_0 + (\varsigma \omega_n - \omega_d)(x_0 - C)}{2\omega_d} \exp(-\omega_d t) \right\}$ 

Critically Damped 
$$\varsigma = 1$$
  $x(t) = C + \{(x_0 - C) + [v_0 + \omega_n(x_0 - C)]t\} \exp(-\omega_n t)$ 

Underdamped 
$$\zeta < 1$$
  $x(t) = C + \exp(-\varsigma \omega_n t) \left\{ (x_0 - C) \cos \omega_d t + \frac{v_0 + \varsigma \omega_n (x_0 - C)}{\omega_d} \sin \omega_d t \right\}$ 

# Calculating natural frequency and damping factor from a measured vibration response



Measure log decrement: 
$$\delta = \frac{1}{n} \log \left( \frac{x(t_0)}{x(t_n)} \right)$$

Measure period: T

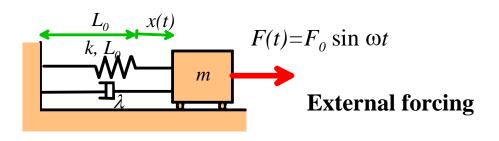
$$\varsigma = \frac{\delta}{\sqrt{4\pi^2 + \delta^2}} \qquad \omega_n = \frac{\sqrt{4\pi^2 + \delta^2}}{T}$$

# Forced Vibrations – concept checklist

#### You should be able to:

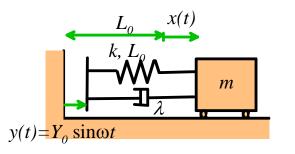
- Be able to derive equations of motion for spring-mass systems subjected to external forcing (several types) and solve EOM using complex vars, or by comparing to solution tables
- 2. Understand (qualitatively) meaning of 'transient' and 'steady-state' response of a forced vibration system (see Java simulation on web)
- Understand the meaning of 'Amplitude' and 'phase' of steady-state response of a forced vibration system
- 4. Understand amplitude-v-frequency formulas (or graphs), resonance, high and low frequency response for 3 systems
- 5. Determine the amplitude of steady-state vibration of forced spring-mass systems.
- 6. Deduce damping coefficient and natural frequency from measured forced response of a vibrating system
- 7. Use forced vibration concepts to design engineering systems

### **EOM** for forced vibrating systems



$$\frac{1}{\omega_n^2} \frac{d^2 x}{dt^2} + \frac{2\varsigma}{\omega_n} \frac{dx}{dt} + x = KF_0 \sin \omega t$$

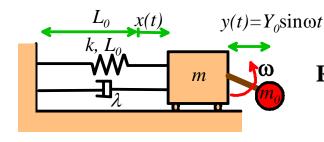
$$\omega_n = \sqrt{\frac{k}{m}}, \quad \varsigma = \frac{\lambda}{2\sqrt{km}}, \quad K = \frac{1}{k}$$



### **Base Excitation**

$$\frac{1}{\omega_n^2} \frac{d^2 x}{dt^2} + \frac{2\varsigma}{\omega_n} \frac{dx}{dt} + x = K \left( y + \frac{2\varsigma}{\omega_n} \frac{dy}{dt} \right)$$

$$\omega_n = \sqrt{\frac{k}{m}}, \quad \varsigma = \frac{\lambda}{2\sqrt{km}}, \quad K = 1$$



**Rotor Excitation** 
$$\frac{1}{\omega_n^2} \frac{d^2 x}{dt^2} + \frac{2\varsigma}{\omega_n} \frac{dx}{dt} + x = -\frac{K}{\omega_n^2} \frac{d^2 y}{dt^2} = K \frac{Y_0 \omega^2}{\omega_n^2} \sin \omega t$$

$$\omega_n = \sqrt{\frac{k}{M}}$$
  $\varsigma = \frac{\lambda}{2\sqrt{kM}}$   $K = \frac{m_0}{M}$   $M = m + m_0$ 

### **Steady-state and Transient solution to EOM**

Equation 
$$\frac{1}{\omega_n^2} \frac{d^2x}{dt^2} + \frac{2\varsigma}{\omega_n} \frac{dx}{dt} + x = C + KF_0 \sin(\omega t)$$
 Initial Conditions  $x = x_0$   $\frac{dx}{dt} = v_0$   $t = 0$ 

Full Solution 
$$x(t) = C + x_h(t) + x_p(t)$$

Steady state part (particular integral)  $x_p(t) = X_0 \sin(\omega t + \phi)$ 

$$X_{0} = \frac{KF_{0}}{\left\{ \left( 1 - \omega^{2} / \omega_{n}^{2} \right)^{2} + \left( 2\varsigma\omega / \omega_{n} \right)^{2} \right\}^{1/2}} \qquad \phi = \tan^{-1} \frac{-2\varsigma\omega / \omega_{n}}{1 - \omega^{2} / \omega_{n}^{2}}$$

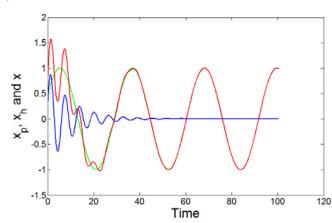
Transient part (complementary integral)

Critically Damped 
$$\zeta = 1$$
  $x_h(t) = C + \left\{ x_0^h + \left[ v_0^h + \omega_n x_0^h \right] t \right\} \exp(-\omega_n t)$ 

Underdamped 
$$\zeta < 1 \qquad x_h(t) = C + \exp(-\varsigma \omega_n t) \left\{ x_0^h \cos \omega_d t + \frac{v_0^h + \varsigma \omega_n x_0^h}{\omega_d} \sin \omega_d t \right\}$$

$$\omega_d = \omega_n \sqrt{\left|1 - \varsigma^2\right|}$$

$$x_0^h = x_0 - C - x_p(0) = x_0 - C - X_0 \sin \phi$$
  $v_0^h = v_0 - \frac{dx_p}{dt}\Big|_{t=0} = v_0 - X_0 \omega \cos \phi$ 



### Canonical externally forced system (steady state solution)

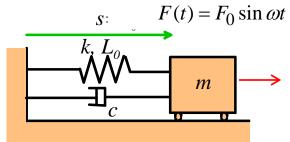
Steady state solution to 
$$\frac{1}{\omega_n^2} \frac{d^2x}{dt^2} + \frac{2\varsigma}{\omega_n} \frac{dx}{dt} + x = C + KF_0 \sin(\omega t) \qquad \omega = 2\pi / T$$

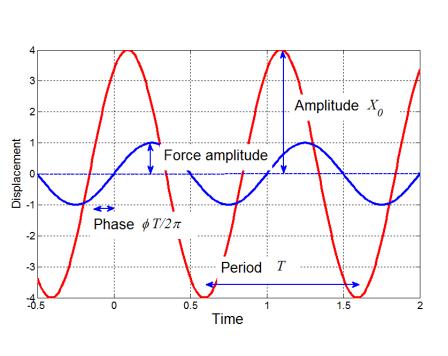
$$\omega_n = \sqrt{\frac{k}{m}} \zeta = \frac{c}{2\sqrt{km}} K = \frac{1}{k}$$

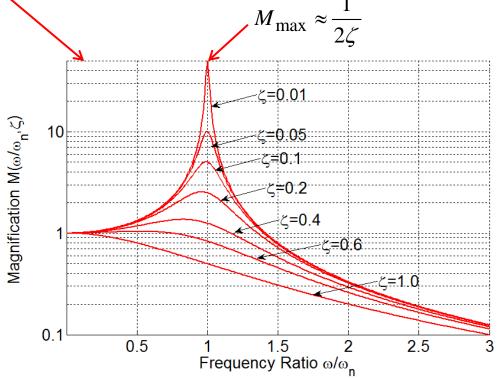
$$x_p(t) = X_0 \sin(\omega t + \phi)$$

$$X_{0} = KF_{0}M(\omega/\omega_{n}, \zeta) \qquad M = \frac{1}{\left\{ \left( 1 - \omega^{2}/\omega_{n}^{2} \right)^{2} + \left( 2\varsigma\omega/\omega_{n} \right)^{2} \right\}^{1/2}} \qquad \phi = \tan^{-1} \frac{-2\varsigma\omega/\omega_{n}}{1 - \omega^{2}/\omega_{n}^{2}}$$

$$\phi = \tan^{-1} \frac{-2\varsigma\omega / \omega_n}{1 - \omega^2 / \omega_n^2}$$







#### Canonical base excited system (steady state solution)

Steady state solution to 
$$\frac{1}{\omega_n^2} \frac{d^2x}{dt^2} + \frac{2\varsigma}{\omega_n} \frac{dx}{dt} + x = C + K \left( y + \frac{2\varsigma}{\omega_n} \frac{dy}{dt} \right)$$

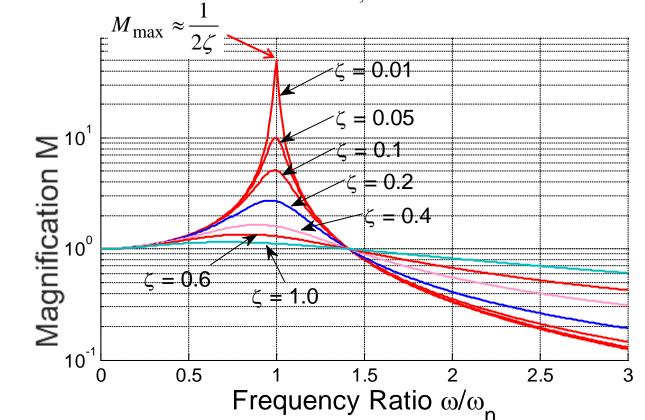
$$\omega_n = \sqrt{\frac{k}{m}} \zeta = \frac{c}{2\sqrt{km}} K = 1$$

$$x_p(t) = X_0 \sin(\omega t + \phi)$$

$$y(t) = Y_0 \sin \omega t$$

$$\omega_{n} = \sqrt{\frac{k}{m}} \zeta = \frac{c}{2\sqrt{km}} K = 1 \qquad x_{p}(t) = X_{0} \sin(\omega t + \phi)$$

$$X_{0} = KY_{0}M(\omega, \omega_{n}, \zeta) \qquad M = \frac{\left\{1 + \left(2\zeta\omega/\omega_{n}\right)^{2}\right\}^{1/2}}{\left\{\left(1 - \omega^{2}/\omega_{n}^{2}\right)^{2} + \left(2\zeta\omega/\omega_{n}\right)^{2}\right\}^{1/2}} \qquad \phi = \tan^{-1}\frac{-2\zeta\omega^{3}/\omega_{n}^{3}}{1 - (1 - 4\zeta^{2})\omega^{2}/\omega_{n}^{2}}$$



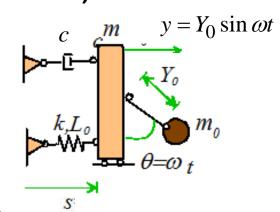
#### Canonical rotor excited system (steady state solution)

Steady state solution to 
$$\frac{1}{\omega_n^2} \frac{d^2x}{dt^2} + \frac{2\varsigma}{\omega_n} \frac{dx}{dt} + x = C - \frac{K}{\omega_n^2} \frac{d^2y}{dt^2}$$

$$\omega_n = \sqrt{\frac{k}{m + m_0}} \zeta = \frac{c}{2\sqrt{k(m + m_0)}} K = \frac{m_0}{m + m_0} x_p(t) = X_0 \sin(\omega t + \phi)$$

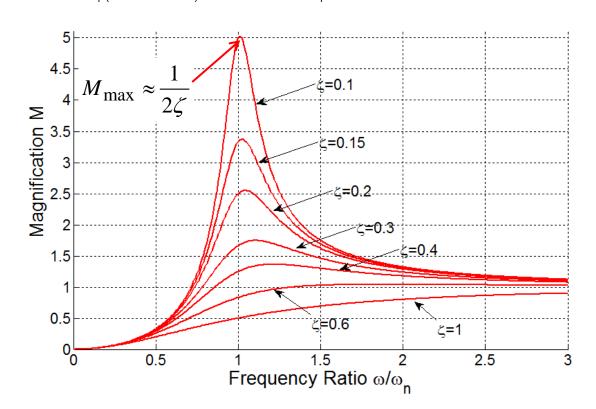
$$K = \frac{m_0}{m + m_0}$$

$$x_p(t) = X_0 \sin(\omega t + \phi)$$



$$X_0 = KY_0 M(\omega, \omega_n, \zeta)$$

$$X_0 = KY_0 M(\omega, \omega_n, \zeta) \qquad M = \frac{\omega^2 / \omega_n^2}{\left\{ \left( 1 - \omega^2 / \omega_n^2 \right)^2 + \left( 2\varsigma\omega / \omega_n \right)^2 \right\}^{1/2}} \qquad \phi = \tan^{-1} \frac{-2\varsigma\omega / \omega_n}{1 - \omega^2 / \omega_n^2}$$



### **Dynamics of Rigid Bodies- concept checklist**

- 1. Understand angular velocity and acceleration vectors; be able to integrate / differentiate angular velocities / accelerations for planar motion.
- 2. Understand formulas relating velocity/acceleration of two points on a rigid body
- 3. Understand constraints at joints and contacts between rigid bodies
- 4. Be able to relate velocities, accelerations, or angular velocities/accelerations of two members in a system of links or rigid bodies
- 5. Be able to analyze motion in systems of gears
- 6. Understand formulas relating velocity/angular velocity and acceleration/angular acceleration of a rolling wheel
- 7. Be able to calculate mass moments of inertia of simple shapes; use parallel axis theorem to shift axis of inertia or calculate mass moments of inertia for a set of rigid bodies connected together
- 8. Understand  $\sum \mathbf{M}_G = I_G \boldsymbol{\alpha}$  for planar motion of a rigid body
- 9. Understand and know when you can use  $\sum \mathbf{M}_0 = I_0 \boldsymbol{\alpha}$
- 10. Be able to calculate accelerations / forces in a system of planar rigid bodies subjected to forces using dynamics equations and kinematics equations
- 11. Understand power/work/potential energy of a rigid body; use energy methods to analyze motion in a system of rigid bodies
- 12. Understand angular momentum of a rigid body; use angular momentum to analyze motion of rigid bodies

# Describing rotational motion of a rigid body

### **Angular velocity vector:**

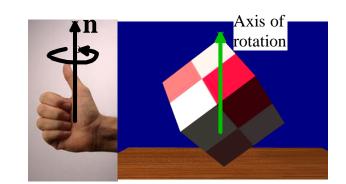
- Direction parallel to rotation axis (RH screw rule)
- Magnitude angle (radians) turned per sec



$$\mathbf{\omega} = \frac{d\theta}{dt}\mathbf{n} = \omega\mathbf{n}$$

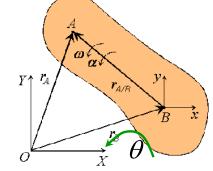
 $\mathbf{\omega} = \frac{d\theta}{dt}\mathbf{n} = \omega\mathbf{n}$  Angular acceleration vector:  $\mathbf{\alpha} = \frac{d\omega}{dt}\mathbf{n}$ 

For planar motion: 
$$\omega = \frac{d\theta}{dt}$$
  $\alpha = \frac{d\omega}{dt} = \frac{d^2\theta}{dt^2}$   $\omega = \frac{d\theta}{dt} \mathbf{k}$   $\omega = \frac{d^2\theta}{dt^2} \mathbf{k}$ 



# Pure Moments (torques): M = Mn

A pure moment is a generalized force that induces rotational motion without translation of center of mass



A motor shaft is an example of an object that exerts a moment – the shaft is parallel to the direction of the moment n

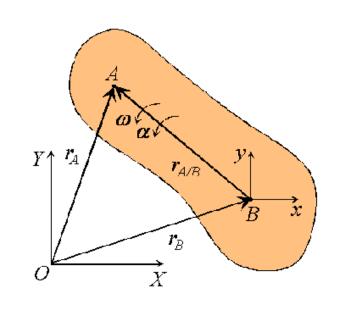
# Rigid body kinematics

Velocities of two points on a rigid body are related by

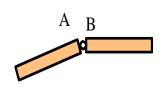
$$\mathbf{v}_A = \mathbf{v}_B + \mathbf{\omega} \times \mathbf{r}_{A/B}$$

Accelerations of two points on a rigid body are related by

$$\mathbf{a}_{A} = \mathbf{a}_{B} + \boldsymbol{\alpha} \times \mathbf{r}_{A/B} + \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}_{A/B})$$

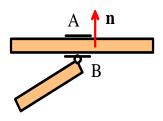


# Continuity conditions



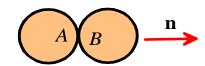
$$\mathbf{v}_A = \mathbf{v}_B$$

$$\mathbf{a}_A = \mathbf{a}_B$$



$$\mathbf{v}_A \cdot \mathbf{n} = \mathbf{v}_B \cdot \mathbf{n}$$

$$\mathbf{a}_{A}\cdot\mathbf{n}=\mathbf{a}_{R}\cdot\mathbf{n}$$



No slip 
$$\mathbf{v}_A = \mathbf{v}_B$$
 Tangential accels equal

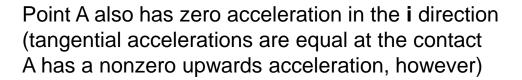
Slip 
$$\mathbf{v}_A \cdot \mathbf{n} = \mathbf{v}_B \cdot \mathbf{n}$$
 Accels arbitrary

### Kinematics of a Rolling Wheel

Wheel has angular velocity  $\omega = \omega \mathbf{k}$ 

Wheel has angular acceleration  $\, oldsymbol{lpha} = lpha {f k} \,$ 

Wheel rolls without slip
This means that velocity of A is zero
(wheel has same velocity as the ground, see animation)



The rigid body formula tells us that

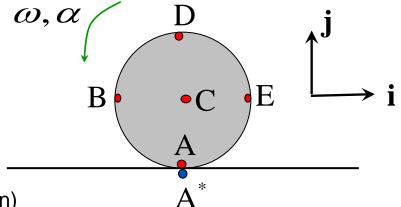
$$\mathbf{v}_C = \mathbf{v}_A + \boldsymbol{\omega} \times \mathbf{r}_{C/A}$$

$$= \mathbf{0} + \omega \mathbf{k} \times R \mathbf{j}$$

$$\mathbf{v}_C = -\omega R \mathbf{i}$$

Then differentiate wrt time to see  $\mathbf{a}_{C}=-\alpha R\mathbf{i}$ 

To find velocity or accel at A, B, D, E use the standard rigid body formulas....



### **Dynamics of rigid bodies**

Preliminary definitions: mass moments of inertia used in planar motion (general 3D more complex)

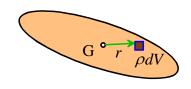
Mass density  $\rho$ 

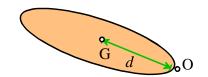
Total Mass: 
$$M = \int_{V} \rho dV$$

Total Mass: 
$$M = \int_{V} \rho dV$$
 COM position:  $\mathbf{r}_{G} = \frac{1}{M} \int_{V} \rho \mathbf{r} dV$ 

Inertia about an axis through origin:  $I = \int_{V} \rho r^2 dV$ 

Parallel Axis Theorem:  $I_O = I_G + md^2$ 





#### **Equations of Motion**

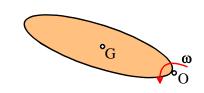
Translational motion  $\mathbf{F} = m\mathbf{a}_G$  (must use acceleration of COM)

$$G$$
 $G$ 
 $G$ 
 $M$ 

Rotational motion 
$$\sum \mathbf{M}_G = \sum_{Forces} \mathbf{r}_{F/G} \times \mathbf{F} + \sum_{Pure\ Moments} \mathbf{M} = I_G \boldsymbol{\alpha}$$

This rotational motion equation is valid **ONLY** for planar motion – 3D motion has another term

For rotation about a fixed axis only 
$$\sum \mathbf{M}_O = \sum_{Forces} \mathbf{r}_{F/O} \times \mathbf{F} + \sum_{Pure\ Moments} \mathbf{M} = I_O \alpha$$



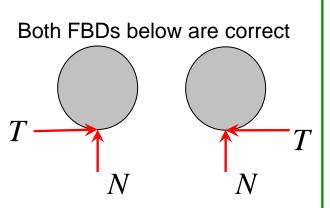
### Free body diagrams with friction

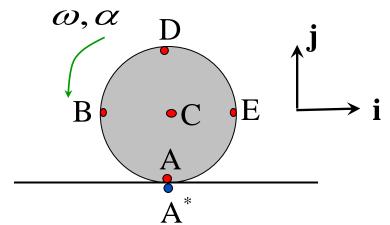
#### Rolling without slip

$$\mathbf{v}_{C} = -\omega R \mathbf{i}$$

$$\mathbf{a}_C = -\alpha R \mathbf{i}$$

$$|T| < \mu N$$



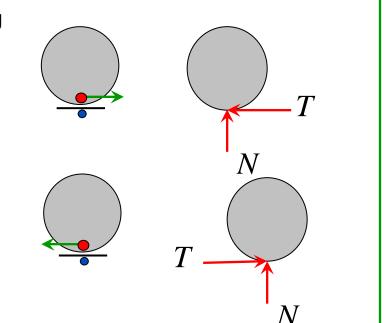


#### Rolling with sliding: Friction force must oppose sliding

$$v_{Cx} + \omega R > 0 \Rightarrow$$
 A moves to right wrt A\*

$$v_{Cx} + \omega R < 0 \Rightarrow$$
 A moves to left wrt A\*

$$T = \mu N$$



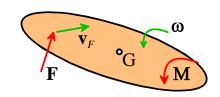
# Analyzing motion of systems of rigid bodies

- 1. Identify each particle/rigid body in the system
- 2. Draw a FBD for each particle / rigid body separately
- 3. Write down  $\mathbf{F} = m\mathbf{a}$  for each rigid body and particle
- 4. Write down  $\sum \mathbf{M}_G = I_G \boldsymbol{\alpha}$  for each rigid body (for rotation about a fixed point can also use  $\sum \mathbf{M}_O = I_O \boldsymbol{\alpha}$
- 5. Look for points in system where acceleration is known or related (eg contacts, joints, etc)
- 6. Use  $\mathbf{a}_G = \mathbf{a}_A + \boldsymbol{\alpha} \times \mathbf{r}_{G/A} + \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}_{G/A})$  to relate accelerations and angular accelerations of rigid bodies
- 7. Solve system of equations from 3, 4, 6 to calculate unknown reactions and accelerations / angular accelerations

### **Energy methods for rigid bodies**

Power (rate of work done) by forces and moments acting on a rigid body

$$P = \sum_{Forces} \mathbf{F} \cdot \mathbf{v}_F + \sum_{Pure\ Moments} \mathbf{M} \cdot \mathbf{\omega}$$

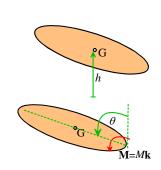


Total work done 
$$W = \int_{t_0}^{t_1} P(t) dt = \int_{t_0}^{t_1} \left( \sum_{Forces} \mathbf{F} \cdot \mathbf{v}_F + \sum_{Pure\ Moments} \mathbf{M} \cdot \mathbf{\omega} \right) dt$$

Gravitational potential energy of a rigid body – use position of COM V = mgh

Potential energy of a constant moment (planar motion only)  $V = -M\theta$ 

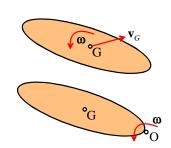
Potential energy of a torsional spring  $V = \frac{1}{2}\kappa\theta^2$ 



#### Kinetic energy of a rigid body

$$T = \frac{1}{2}m|\mathbf{v}_G|^2 + \frac{1}{2}I_G|\mathbf{\omega}|^2$$
 General – can always use this

$$T = \frac{1}{2}I_O \left|\omega\right|^2$$
 Rotation about a **fixed axis** only (use parallel axis theorem to find lo)

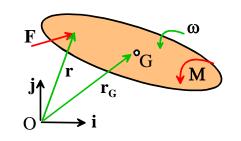


Power-KE relation 
$$P = \frac{dT}{dt}$$
 Work-KE relation  $W = T_1 - T_0$ 

Work- energy relation for a conservative system  $W^{ext} = T_1 + V_1 - (T_0 + V_0)$ If no external work is done on a conservative system  $T_1 + V_1 = (T_0 + V_0)$ 

# Angular momentum for rigid bodies

$$\mathbf{A}_{G} = \int_{t_{0}}^{t_{1}} \left( \sum_{Forces} (\mathbf{r} - \mathbf{r}_{G}) \times \mathbf{F} + \sum_{Pure\ Moments} \mathbf{M} \right) dt$$



Angular impulse about a fixed point 
$$\mathbf{A}_0 = \int_{t_0}^{t_1} \left( \sum_{Forces} \mathbf{r} \times \mathbf{F} + \sum_{Pure\ Moments} \mathbf{M} \right) dt$$

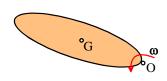
Angular momentum about COM  $\mathbf{h}_G = I_G \mathbf{\omega}$ 

$$\mathbf{h}_G = I_G \mathbf{\omega}$$

Angular momentum about a fixed point  $\mathbf{h}_o = \mathbf{r}_c \times m\mathbf{v}_c + I_c\mathbf{\omega}$ 

Special case: rotation about a fixed point  $\mathbf{h}_{\alpha} = I_{\alpha} \mathbf{\omega}$ 

$$\mathbf{h}_{o} = I_{o}\mathbf{\omega}$$



Impulse-momentum relations (COM) 
$$\sum_{Forces} (\mathbf{r} - \mathbf{r}_G) \times \mathbf{F} + \sum_{Pure\ Moments} \mathbf{M} = \frac{d\mathbf{h}_G}{dt} \qquad \mathbf{A}_G = \mathbf{h}_{G1} - \mathbf{h}_{G0}$$

$$\mathbf{A}_G = \mathbf{h}_{G1} - \mathbf{h}_{G0}$$

Momentum is conserved if  $A_c = 0$ 

Impulse-momentum relations (Fixed point)

$$\sum_{Forces} \mathbf{r} \times \mathbf{F} + \sum_{Pure\ Moments} \mathbf{M} = \frac{d\mathbf{h}_{O}}{dt}$$

$$\mathbf{A}_O = \mathbf{h}_{O1} - \mathbf{h}_{O0}$$

Momentum is conserved if  $A_0 = 0$