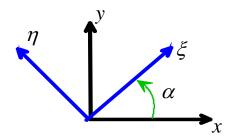


#### **EN224: Linear Elasticity**

## Homework 5: Complex variable methods for plane problems Due Friday April 22, 2005

**Division of Engineering Brown University** 



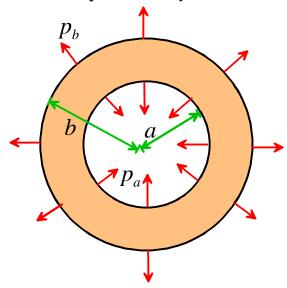
**1.** Let  $\Omega(z)$ ,  $\omega(z)$  be two complex potentials that generate stresses and displacements according to the usual formulation (no continuation)

$$\begin{split} &2\mu(u_x+iu_y) = \left(3-4\nu\right)\Omega(z)-z\overline{\Omega'(z)}-\overline{\omega(z)}\\ &\sigma_{xx}+\sigma_{yy} = 2\left\{\Omega'(z)+\overline{\Omega'(z)}\right\}\\ &\sigma_{xx}-\sigma_{yy}+2i\sigma_{xy} = -2\left\{z\overline{\Omega''(z)}+\overline{\omega'(z)}\right\} \end{split}$$

Show that the displacement and stress components in the  $(\xi,\eta)$  basis shown in the figure can be calculated as

$$\begin{split} &2\mu(u_{\xi}+iu_{\eta}) = \left[\left(3-4\nu\right)\Omega(z)-z\overline{\Omega'(z)}-\overline{\omega(z)}\right]e^{-i\alpha}\\ &\sigma_{\xi\xi}+\sigma_{\eta\eta}=2\left\{\Omega'(z)+\overline{\Omega'(z)}\right\}\\ &\sigma_{\xi\xi}-\sigma_{\eta\eta}+2i\sigma_{\xi\eta}=-2\left\{z\overline{\Omega''(z)}+\overline{\omega'(z)}\right\}e^{-2i\alpha} \end{split}$$

2. Complex variable solution to a pressurized cylinder.



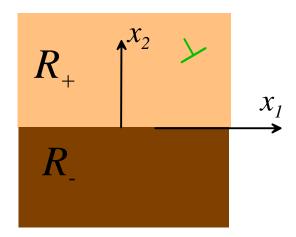
2.1 Using the results of the preceding problem, show that the complex potentials that generate stress and displacement fields in a pressurized cylinder (see above) must satisfy

$$\Omega'(z) + \overline{\Omega'(z)} - \left(z\overline{\Omega''(z)} + \overline{\omega'(z)}\right)e^{-2i\theta} = p_a \qquad z = ae^{i\theta}$$

$$\Omega'(z) + \overline{\Omega'(z)} - \left(z\overline{\Omega''(z)} + \overline{\omega'(z)}\right)e^{-2i\theta} = p_b \qquad z = be^{i\theta}$$

2.2 By expanding  $\Omega(z)$  and  $\omega(z)$  as Laurent series, find the potentials that solve this problem. To simplify the algebra, note that  $e^{-2i\theta} = r^2/z^2$   $\overline{z} = r^2/z$  on  $z = re^{i\theta}$ , and assume that the solution can be generated from terms in the series that after substitution in the boundary conditions, are independent of z.

### 3. Dislocation near a rigid interface



3.1 Begin by finding an analytic continuation that automatically satisfies D=0 on  $z=\overline{z}$ . To do this, start with the standard complex variable formulation

$$\begin{split} &2\mu(u_1+iu_2) = \left(3-4\nu\right)\Omega(z) - z\overline{\Omega'(z)} - \overline{\omega(z)} \\ &\sigma_{11} + \sigma_{22} = 2\left\{\Omega'(z) + \overline{\Omega'(z)}\right\} \\ &\sigma_{11} - \sigma_{22} + 2i\sigma_{12} = -2\left\{z\overline{\Omega''(z)} + \overline{\omega'(z)}\right\} \end{split}$$

Express the boundary condition in terms of  $\Omega$  and  $\omega$  defined in R+. Next, express the boundary condition in terms of potentials  $\overline{\Omega(\overline{z})}$   $\overline{\omega(\overline{z})}$  which are analytic in R-. Use the result to show that

$$(3-4\nu)\lim_{z\to L+}\overline{\Omega(z)} = \lim_{z\to L-} \left\{ \overline{z}\Omega'(\overline{z}) + \omega(\overline{z}) \right\}$$

where L denotes the real axis. Hence, conclude that this implies that

$$\Omega(z) = \begin{cases} \Omega(z) & z \in R + \\ \left(z\overline{\Omega'(\overline{z})} + \overline{\omega(\overline{z})}\right)/(3 - 4\nu) & z \in R - \end{cases}$$

is analytic in R. Use this to calculate an expression for  $\omega(z)$   $z \in R+$ , and hence show that a solution with D=0 on  $z=\overline{z}$  can be generated by finding a single potential  $\Omega(z)$  that is analytic in R, and calculating displacements and stresses from

$$\begin{split} &2\mu(u_1+iu_2)=\left(3-4\nu\right)\left\{\Omega(z)-\Omega(\overline{z})\right\}+(\overline{z}-z)\overline{\Omega'(z)}\\ &\sigma_{11}+\sigma_{22}=2\left\{\Omega'(z)+\overline{\Omega'(z)}\right\}\\ &\sigma_{11}-\sigma_{22}+2i\sigma_{12}=-2\left\{\left(3-4\nu\right)\Omega'(\overline{z})-\overline{\Omega'(z)}+(z-\overline{z})\overline{\Omega''(z)}\right\} \end{split}$$

3.2 Let 
$$\Omega_0(z) = i \frac{b_1 + ib_2}{8\pi(1-v)} \log(z-z_0)$$
  $\omega_0(z) = -i \frac{b_1 - ib_2}{8\pi(1-v)} \log(z-z_0)$  generate the

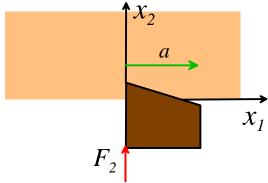
solution for a dislocation at position  $z_0$  in an infinite solid. Deduce that, to satisfy D=0 on L, we must superpose a second potential  $\Omega(z)$  satisfying

$$(3-4\nu) \left( \lim_{z \to L^{+}} \Omega(z) - \lim_{z \to L^{-}} \Omega(z) \right) = -(3-4\nu)\Omega_{0}(x_{1}) + x_{1}\overline{\Omega_{0}'(x_{1})} + \overline{\omega_{0}(x_{1})}$$

and calculating stresses and displacements from this potential using the formulation in 3.1.

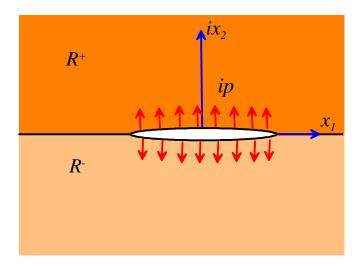
3.3 Using (3.1) as a guide, write down the potential  $\Omega(z)$  in terms of  $\Omega_0(z)$  and  $\omega_0(z)$ 

#### 4. Stress induced by indentation with a rigid wedge



Suppose that an elastic half-space is indented by a rigid frictionless wedge, with profile  $f(x_1) = \varepsilon x_1$ . Calculate the potential  $\Omega'(z)$  that generates the stress field in the solid in terms of a and  $F_2$ . Hence, determine the contact pressure distribution and the slope of the surface for  $x_1 > 0$ . Use the conditions that the contact pressure cannot be tensile, and the two solids cannot overlap to deduce the relationship between contact width a and the force applied to the punch.

# 5. An alternative solution for the pressurized crack



There is often more than one choice of analytic continuation for a particular boundary value problem. To illustrate this, in this problem we will devise an alternative procedure to solve the pressurized crack problem that was discussed in class.

Consider a crack that is subjected to equal and opposite tractions  $t_1 + it_2 = ip$  on its faces. Symmetry conditions imply that  $u_2 = 0$  or  $\sigma_{12} = 0$  on  $z = \overline{z}$  outside the crack. Moreover, it is evidently sufficient to find a solution in the upper half-plane, since the solution in the lower half-plane follows by symmetry.

5.1 Starting with the standard complex variable formulation

$$\begin{split} &2\mu(u_1+iu_2) = \left(3-4v\right)\Omega(z) - z\overline{\Omega'(z)} - \overline{\omega(z)} \\ &\sigma_{11} + \sigma_{22} = 2\left\{\Omega'(z) + \overline{\Omega'(z)}\right\} \\ &\sigma_{22} - i\sigma_{12} = \Omega'(z) + \overline{\Omega'(z)} + z\overline{\Omega''(z)} + \overline{\omega'(z)} \end{split}$$

show that setting  $\omega(z) = \Omega(z) - z\Omega'(z)$  will automatically satisfy  $\sigma_{12} = 0$  on  $z = \overline{z}$ .

5.2 With this choice of  $\omega(z)$ , show that the condition that  $u_2 = 0$  on  $z = \overline{z}$  implies that

$$\lim_{z \to L+} \Omega(z) = \lim_{z \to L-} \overline{\Omega(\overline{z})}$$

showing that

$$\theta(z) = \begin{cases} \Omega(z) & z \in R + \\ \overline{\Omega(\overline{z})} & z \in R - \end{cases}$$

is continuous outside the crack, and analytic in the whole plane. Deduce that  $\overline{\Omega(z)} = \theta(\overline{z})$ .

- 5.3 Hence show that traction boundary condition on the crack faces leads to a Hilbert problem for  $\theta(z)$
- 5.4 Write down the general solution to the Hilbert problem.
- 5.5 Hence, find an expression for the stress intensity factors induced at the right hand crack tip by a pair of equal and opposite point forces acting on the crack faces.

