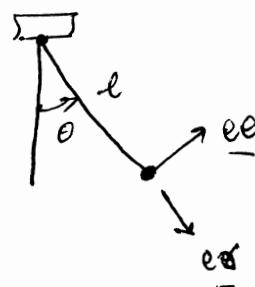
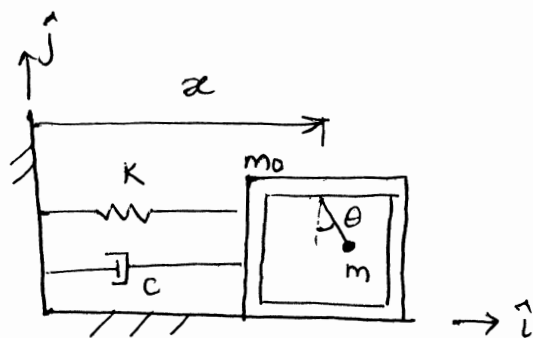


Problem 1

Hw6



For mass m_0 velocity $\underline{v} = \dot{x} \underline{\hat{i}}$

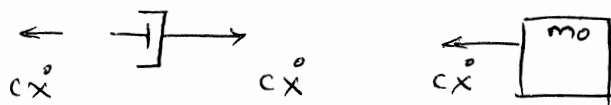
For mass m velocity $\underline{v} = \dot{x} \underline{\hat{i}} + l \dot{\theta} \underline{\hat{e}}_{\theta}$
 $= (\dot{x} + l \dot{\theta} \cos \theta) \underline{\hat{i}} + l \dot{\theta} \sin \theta \underline{\hat{j}}$

$\Rightarrow$ Kinetic Energy $= T = \frac{1}{2} m_0 \dot{x}^2 + \frac{1}{2} m (\dot{x}^2 + l^2 \dot{\theta}^2 + 2 l \dot{\theta} \dot{x} \cos \theta)$

Potential Energy $= V = \frac{1}{2} k (x - l_0)^2 + m g l (1 - \cos \theta)$

Lagrangian $L = T - V$

Non-conservative force:



$$Q_x^{NC} \delta x = -c \dot{x} \delta x \Rightarrow Q_x^{NC} = -c \dot{x}, \quad Q_{\theta}^{NC} = 0$$

$$\text{Then, } \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{x}} \right) - \frac{\partial L}{\partial x} = Q_x^{NC} \quad \{$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{\theta}} \right) - \frac{\partial L}{\partial \theta} = Q_{\theta}^{NC} \quad \text{gives}$$

Equations of motion:

$$1) (m + m_0) \ddot{x} + m l (\ddot{\theta} \cos \theta - \dot{\theta}^2 \sin \theta) + k(x - l_0) + c \dot{x} = 0$$

$$2) l \ddot{\theta} + \ddot{x} \cos \theta + g \sin \theta = 0$$

Problem 2 Greenwood 6-26.

Kinetic Energy:

$$T = \frac{1}{2} 3m(\dot{x}^2 + \dot{y}^2) + \frac{1}{2} m\ell^2 \dot{\theta}^2 \quad (1.1)$$

Non-holonomic Constraint:

$$\frac{\ell}{2\sqrt{3}} \dot{\theta} - \dot{x} \sin \theta + \dot{y} \cos \theta = 0 \quad (1.2)$$

Lagrange's Equations:

$$3m\ddot{x} = -\lambda \sin \theta \quad (1.3)$$

$$3m\ddot{y} = \lambda \cos \theta \quad (1.4)$$

$$m\ell^2 \ddot{\theta} = +\frac{1}{2\sqrt{3}} \ell \lambda \quad (1.5)$$

Eliminate λ :

$$\ddot{x} \cos \theta + \ddot{y} \sin \theta = 0 \quad (1.6)$$

$$\ell \ddot{\theta} + \frac{\sqrt{3}}{2} \ddot{x} \sin \theta - \frac{\sqrt{3}}{2} \ddot{y} \cos \theta = 0 \quad (1.7)$$

Introduce basis $\mathbf{e}_r, \mathbf{e}_\theta$:

$$\mathbf{e}_r = \mathbf{i} \cos \theta + \mathbf{j} \sin \theta \quad (1.8)$$

$$\mathbf{e}_\theta = -\mathbf{i} \sin \theta + \mathbf{j} \cos \theta$$

Note that for $\bar{\mathbf{v}}_r = \bar{v}_r \mathbf{e}_r + \bar{v}_\theta \mathbf{e}_\theta$, Eqn. (1.6) corresponds to:

$$\mathbf{e}_r \cdot \bar{\mathbf{a}} = 0 \Rightarrow \dot{\bar{v}}_r - \dot{\theta} \bar{v}_\theta = 0 \quad (1.9)$$

and Eqn. (1.5) corresponds to:

$$\frac{\ell}{2\sqrt{3}} \dot{\theta} + \bar{v}_\theta = 0. \quad (1.10)$$

Eliminating $\bar{v}_\theta$ from Eqns. one obtains

$$\dot{\bar{v}}_r + \frac{\ell}{2\sqrt{3}} (\dot{\theta})^2 = 0 \quad (1.11)$$

whereas differentiation of (1.5) gives

$$\frac{5\ell}{2\sqrt{3}} \ddot{\theta} - \dot{\theta} \bar{v}_r = 0. \quad (1.12)$$

Eliminating $\bar{v}_r$ between (1.11) and (1.12) gives

$$5\dot{\theta} \ddot{\theta} - 5(\ddot{\theta})^2 + (\dot{\theta})^4 = 0. \quad (1.13)$$

Eqn. (1.13) is a nonlinear ordinary differential equation for $\dot{\theta}(t)$. Because this equation is homogeneous (of degree 4) in the time derivatives, this equation does not determine the time dependence of $\dot{\theta}(t)$. Instead it suggests looking for a solution $\dot{\theta} = \dot{\theta}(\theta)$. Because exponentials can cancel out as a common factor we look for a solution of the form

$$\dot{\theta} = \exp[f(\theta)] \quad (1.14)$$

for which

$$\begin{aligned} \ddot{\theta} &= \{\exp[f(\theta)]\}^2 f'(\theta) \\ \ddot{\theta} &= \{\exp[f(\theta)]\}^3 [f''(\theta) + 2(f'(\theta))^2]. \end{aligned} \quad (1.15)$$

Substitution of (1.14) and (1.15) into (1.13) and cancellation of the common exponential term leads to the following differential equation for $f(\theta)$.

$$5f''(\theta) + 5[f'(\theta)]^2 + 1 = 0 \quad (1.16)$$

Equation (1.16) has the solutions

$$f'(\theta) = \pm \frac{i}{\sqrt{5}}, \quad i \equiv \sqrt{-1}. \quad (1.17)$$

Integration of (1.17) gives

$$f(\theta) = \pm i\theta + c \quad (1.18)$$

where c is a constant. Substitution of (1.18) into (1.14) gives a solution of the form

$$\dot{\theta} = C \left[\cos(\theta/\sqrt{5}) \pm i \sin(\theta/\sqrt{5}) \right] \quad (1.19)$$

where $C \equiv \exp[c]$ is a constant to be determined from the initial condition for $\dot{\theta}(0)$, which, from the initial conditions $\theta(0) = 0$, $\dot{y}(0) = v_0$, and the non-holonomic constraint (1.5) is given by

$$C = -\frac{2\sqrt{3}v_0}{\ell}. \quad (1.20)$$

Substitution of (1.20) into (1.19) and selecting only the real part of (1.19), which is also a solution of (1.13), the solution of (1.13) becomes

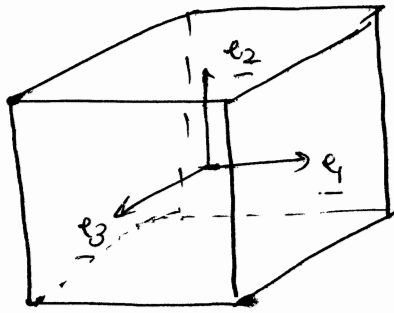
$$\dot{\theta} = -\frac{2\sqrt{3}v_0}{\ell} \cos(\theta/\sqrt{5}) \quad (1.21)$$

for $[\theta/\sqrt{5}] \geq -\pi/2$; once θ becomes equal to $-\sqrt{5}\pi/2$ it becomes constant and stays at this value indefinitely. Substitution of (1.21) into (1.4) gives

$$\lambda = \frac{12\sqrt{3}mv_0^2}{\ell\sqrt{5}} \sin(2\theta/\sqrt{5}) \quad (1.22)$$

for $[\theta/\sqrt{5}] \geq -\pi/2$; afterwards, $\lambda = 0$, $\bar{v}_\theta = 0$, $\bar{v}_r = -\sqrt{5}v_0$.

Problem 3



$$\underline{e}_1' = \frac{1}{\sqrt{3}} (\underline{e}_1 + \underline{e}_2 + \underline{e}_3)$$

$$\underline{e}_2' = \frac{1}{\sqrt{6}} (\underline{e}_1 + \underline{e}_2 - 2\underline{e}_3)$$

$$\underline{e}_3' = \frac{1}{\sqrt{2}} (-\underline{e}_1 + \underline{e}_2)$$

$$\bar{\mathbf{I}} \Rightarrow I_{xx} = I_{yy} = I_{zz} = \frac{ma^2}{6} \quad (\text{for cube of side } a)$$

$$[\bar{\mathbf{I}}] = \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix} \cdot \frac{ma^2}{6}$$

$$[\alpha] = \begin{bmatrix} 1/\sqrt{3} & 1/\sqrt{3} & 1/\sqrt{3} \\ 1/\sqrt{6} & 1/\sqrt{6} & -2/\sqrt{6} \\ -1/\sqrt{2} & 1/\sqrt{2} & 0 \end{bmatrix}$$

$$[\bar{\mathbf{I}}'] = [\alpha][\bar{\mathbf{I}}][\alpha]^T = [\bar{\mathbf{I}}] = \frac{ma^2}{6} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

(Since $\bar{\mathbf{I}}$ is a diagonal matrix)