

## Review

Tensor: Linear mapping of one vector onto another  $\mathbf{v} = \mathbf{S}\mathbf{u}$

Tensor components (orthonormal basis)  $S_{ij} = \mathbf{m}_i \cdot \mathbf{S}\mathbf{m}_j$   $v_i = S_{ij}v_j$

Dyadic Product  $\mathbf{S} = \mathbf{a} \otimes \mathbf{b}$  satisfies  $\mathbf{S}\mathbf{u} = (\mathbf{a} \otimes \mathbf{b})\mathbf{u} = \mathbf{a}(\mathbf{b} \cdot \mathbf{u})$   $S_{ij} = a_i b_j$

Dyadic expansion  $\mathbf{S} = S_{ij} \mathbf{m}_i \otimes \mathbf{m}_j$

Change of Basis  $\mathbf{S} = S_{ij}^{(\mathbf{m})} \mathbf{m}_i \otimes \mathbf{m}_j = S_{ij}^{(\mathbf{e})} \mathbf{e}_i \otimes \mathbf{e}_j$   $S_{kl}^{(\mathbf{m})} = Q_{ki} S_{ij}^{(\mathbf{e})} Q_{lj}$

$$Q_{ij} = \mathbf{m}_i \cdot \mathbf{e}_j$$

### Basic Tensor Operations

Addition:  $\mathbf{U} = \mathbf{S} + \mathbf{T}$   $U_{ij} = S_{ij} + T_{ij}$

Tensor-vector products:  $\mathbf{v} = \mathbf{S}\mathbf{u}$   $v_i = S_{ij}u_j$   $\mathbf{v} = \mathbf{u}\mathbf{S}$   $v_i = u_j S_{ji}$

Product of two tensors:  $\mathbf{U} = \mathbf{T}\mathbf{S}$   $U_{ij} = T_{ik} S_{kj}$  NB:  $(\mathbf{T}\mathbf{S}) \neq (\mathbf{S}\mathbf{T})$

Transpose:  $\mathbf{S}^T$   $\mathbf{u} \cdot \mathbf{S}^T = \mathbf{S} \cdot \mathbf{u}$   $S_{ij}^T = S_{ji}$

Note:  $(\mathbf{S}\mathbf{T})^T = \mathbf{T}^T \mathbf{S}^T$  (show with index notation)

Trace:  $\text{trace}(\mathbf{S}) = S_{11} + S_{22} + S_{33}$   $\text{trace}(\mathbf{S}) = S_{kk}$

## Contracted Products

Inner Product  $T : S \equiv T_{ij} S_{ij}$

Outer Product  $T \cdot S \equiv T_{ij} S_{ji} \equiv T : S^T$

## Determinant

$$\lambda = \det \mathbf{A} \equiv \lambda = \frac{1}{6} \epsilon_{ijk} \epsilon_{lmn} A_{li} A_{mj} A_{nk} = \epsilon_{ijk} A_{i1} A_{j2} A_{k3}$$

$$\Leftrightarrow \epsilon_{lmn} \lambda = \epsilon_{ijk} A_{li} A_{mj} A_{nk} = \epsilon_{ijk} A_{il} A_{jm} A_{kn}$$

## Useful Results

$$\det(\mathbf{A}\mathbf{B}) = \det(\mathbf{A}) \det(\mathbf{B})$$

$$\det(\mathbf{A}^T) = \det(\mathbf{A})$$

A tensor is nonsingular  $\Leftrightarrow \det(S) \neq 0$

## Inverse of a nonsingular tensor

$$S^{-1} S = S S^{-1} = I \quad S_{ij}^{-1} S_{jk} = \delta_{ik}$$

$$S_{ij}^{-1} = \frac{1}{\det(S)} \epsilon_{jpr} \epsilon_{ikl} S_{pk} S_{ql}$$

A tensor is orthogonal if  $S^{-1} = S^T$   
 is "proper" orthogonal if  $S^{-1} = S^T$   
 and  $\det(S) > 0$

### Example

Let  $J = \det(S)$ . Show that  $\frac{\partial J}{\partial S_{mn}} = J S_{nm}^{-1}$ . (this is a very useful result – we often need to differentiate

the determinant of a tensor when working with constitutive equations for materials, for example)

Note  $\frac{\partial S_{ij}}{\partial S_{mn}} = \delta_{im} \delta_{jn}$

Recall  $J = \frac{1}{6} \epsilon_{pqr} \epsilon_{ijk} S_{pi} S_{qj} S_{rk}$

$$\frac{\partial J}{\partial S_{mn}} = \frac{1}{6} \epsilon_{pqr} \epsilon_{ijk} \left\{ \delta_{pm} \delta_{in} S_{qj} S_{rk} + S_{pi} \delta_{qm} \delta_{jn} S_{rk} + S_{pi} S_{qj} \delta_{rm} \delta_{kn} \right\}$$

Recall  $\epsilon_{ijk} = \epsilon_{jki}$

$$\frac{\partial J}{\partial S_{mn}} = \frac{1}{2} \epsilon_{mqr} \epsilon_{njk} S_{qj} S_{rk} = \det(S) S_{nm}^{-1}$$

## Invariants of a tensor

Definition: A scalar function of components that is independent of basis

Example:  $\text{trace}(S) = S_{kk}$

$$S_{kk}^{(m)} = Q_{ki} S_{ij}^e Q_{kj}$$

$$= Q_{ki} Q_{kj} S_{ij}^e = \delta_{ij} S_{ij}^e = S_{ii}^e$$

3 independent invariants

Standard set :  $I_1 = \text{trace}(S)$

$$I_2 = \frac{1}{2} (\text{trace}(S)^2 - S \cdot S)$$

$$= \frac{1}{2} (S_{ii} S_{jj} - S_{ij} S_{ji})$$

$$I_3 = \det(S)$$

## Eigenvalues & Eigenvectors $\{\lambda, \underline{u}\}$

$$S \underline{u} = \lambda \underline{u}$$

To calculate  $\lambda$ :  $(S - \lambda I) \underline{u} = \underline{0}$

$$\Rightarrow \det(S - \lambda I) = 0$$

$$\lambda^3 - I_1 \lambda^2 + I_2 \lambda - I_3 = 0 \quad \text{Characteristic equation}$$

### Observations:

(1) At least 1 real eigenvalue

(2) A symmetric tensor has real eigenvalues and orthogonal eigenvectors

Proof: Let  $\{\lambda_u, \underline{u}\}$ ,  $\{\lambda_v, \underline{v}\}$

$$\begin{aligned} \underline{v} \cdot S \underline{u} &= \lambda_u \underline{v} \cdot \underline{u} = \underline{u} \cdot S^T \underline{v} = \underline{u} \cdot S \underline{v} \\ &= \lambda_v \underline{v} \cdot \underline{u} \end{aligned}$$

$$(\lambda_u - \lambda_v) \underline{u} \cdot \underline{v} = 0 \Rightarrow \underline{u} \cdot \underline{v} = 0$$

Also Let  $(\lambda, \underline{u})$ ,  $(\bar{\lambda}, \bar{\underline{u}})$  are a conjugate pair

$$\begin{aligned} \bar{\underline{u}} \cdot S \underline{u} &= \lambda \underline{u} \cdot \bar{\underline{u}} = \underline{u} \cdot S^T \bar{\underline{u}} = \underline{u} \cdot S \bar{\underline{u}} \\ &= \bar{\lambda} \bar{\underline{u}} \cdot \underline{u} \Rightarrow \lambda = \bar{\lambda} \end{aligned}$$

(3) Orthogonal tensors have eigenvalues  $1$ ,  $e^{\pm i\theta}$  for some  $\theta$

$$\text{Let } R^T R = I$$

$$\begin{aligned}
 \text{Proof: } \quad \underline{R} \underline{\bar{u}} \cdot \underline{R} \underline{u} &= \bar{\lambda} \lambda \underline{\bar{u}} \cdot \underline{u} \\
 &= \underline{\bar{u}} \cdot \underline{R}^T \underline{R} \underline{u} \\
 &= \underline{\bar{u}} \cdot \underline{u} \Rightarrow \bar{\lambda} \lambda = 1
 \end{aligned}$$

### Cayley - Hamilton Theorem:

Every tensor satisfies its own characteristic eq

$$S^3 - I_1 S^2 + I_2 S - I_3 = 0 \quad *$$

Proof: Let  $T(\alpha)$  satisfy  $T(S - \alpha I) = \det(S - \alpha I) I$

$$T(S - \alpha I) = (-\alpha^3 + I_1 \alpha^2 - I_2 \alpha + I_3) I$$

$$\text{Guess } T = I \alpha^2 + T_2 \alpha + T_3$$



$$\begin{aligned} \Rightarrow T(S - \alpha I) &= -\alpha^3 I + (S - T_2)\alpha^2 - (T_3 - T_2 S)\alpha + T_3 S \\ &= -\alpha^3 I + I_1 I \alpha^2 - I_2 I \alpha + I_3 I \end{aligned}$$

$$S - T_2 = I_1 I \quad T_3 - T_2 S = I_2 I \quad T_3 S = I_3 I$$

$\Rightarrow$  Subst back into \*

$$S^3 - (S - T_2)S^2 + (T_3 - T_2 S)S - T_3 S = 0!$$