

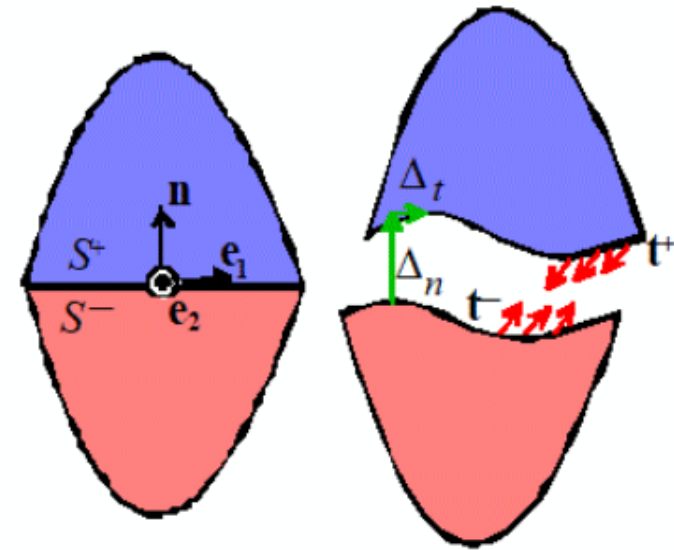
11) Interfaces and Contact

11.1 "Cohesive Zone" models of fracture

Goal: model crack nucleation and growth along a weak interface

Approach: Introduce planes Γ where solid can separate

Introduce additional constitutive equations that specify tractions acting on material planes as functions of separation



Cohesive Zone Laws

Kinematics & Kinetics

Let $\Delta_n = (\underline{u}^+ - \underline{u}^-) \cdot \underline{n}$

$$\Delta_\alpha = (\underline{u}^+ - \underline{u}^-) \cdot \underline{e}^\alpha$$

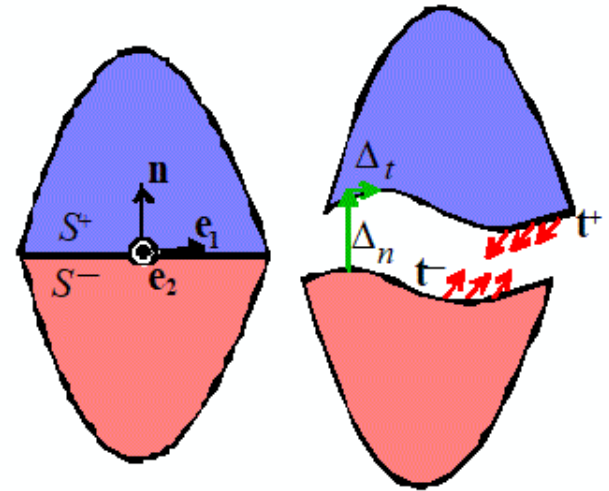
also $T_n = \underline{t}^- \cdot \underline{n}$

$$T_\alpha = \underline{t}^- \cdot \underline{e}^\alpha$$

We usually assume Δ_n , Δ_α are small

For large relative displacements use contact elements

Constitutive laws must give $\{T_n, T_\alpha\} = f(\Delta_n, \Delta_\alpha)$



Examples: Reversible cohesive zone law

Let $\Phi(\Delta_n, \Delta_\alpha)$ be work done per unit area to separate planes

$$T_n = \frac{\partial \Phi}{\partial \Delta_n} \quad T_\alpha = \frac{\partial \Phi}{\partial \Delta_\alpha}$$

Example: Xu-Needleman

$$\Phi(\Delta_n, \Delta_t) = \phi_n - \phi_n \left(1 + \frac{\Delta_n}{\delta_n} \right) \exp\left(-\frac{\Delta_n}{\delta_n}\right) \exp\left(-\frac{\beta^2 \Delta_t^2}{\delta_n^2}\right)$$

ϕ_n = interface energy

δ_n, β : material props controlling compliance

Hence

$$T_n = \sigma_{\max} \frac{\Delta_n}{\delta_n} \exp\left(1 - \frac{\Delta_n}{\delta_n}\right)$$

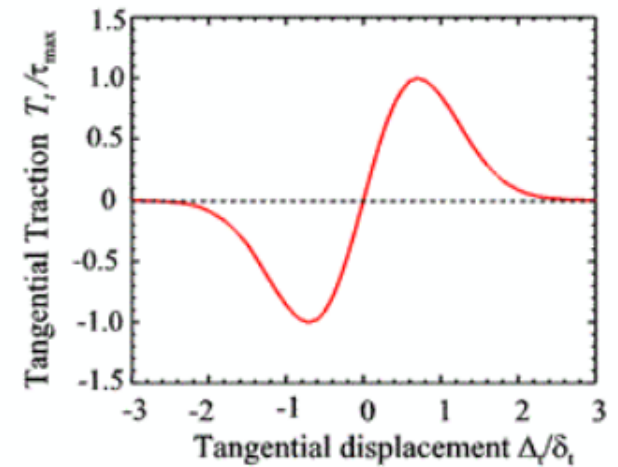
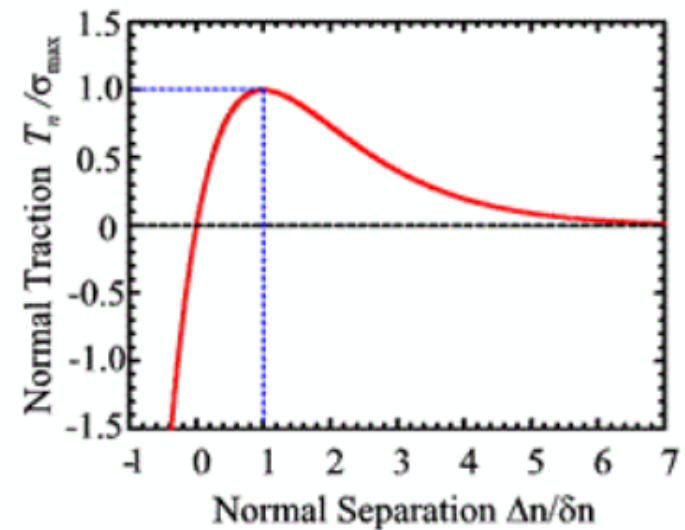
$$T_\alpha = 2\sigma_{\max} \left(\frac{\beta^2 \Delta_\alpha}{\delta_n}\right) \left(1 + \frac{\Delta_n}{\delta_n}\right) \exp\left(1 - \frac{\Delta_n}{\delta_n}\right) \exp\left(-\frac{\beta^2 \Delta_t^2}{\delta_n^2}\right)$$

$$\sigma_{\max} = \frac{\sigma_n}{\delta_n \exp(1)}$$

$$\tau_{\max} = \beta^2 \sigma_{\max}$$

Notes: (1) Interface is reversible:
Heals if interfaces contact
after separation

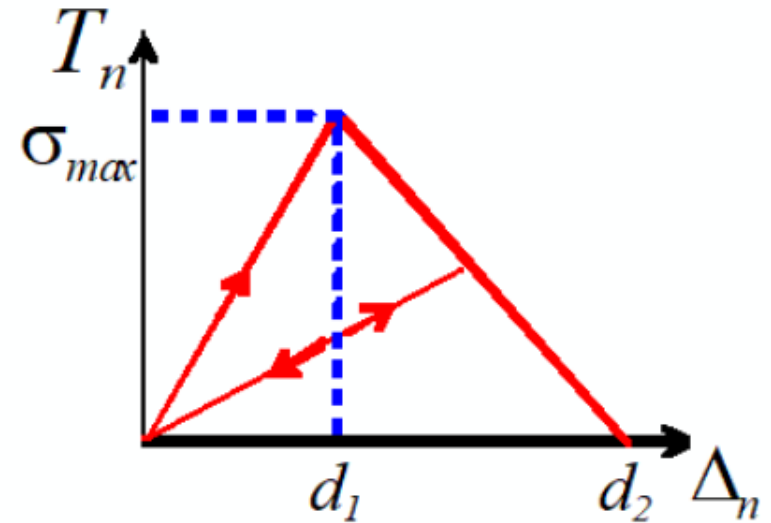
(2) Interface can behave strangely under
large pressure



Irreversible Cohesive zone laws

Assumptions:

- (1) No inter-penetration $\Delta_n \geq 0$
- (2) Irreversible behavior in tension
 - (i) Undamaged interface has stiffness $k_0 = \frac{\sigma_{max}}{d_1}$
 - (2) If separation exceeds d_1 interface stiffness drops



Damage is quantified by an internal variable $0 < D < 1$

Undamaged solid has $D = 0$
 Total failure $D = 1$

To quantify magnitude of separation and traction

$$\lambda = \sqrt{\Delta_n^2 + \beta^2 (\Delta_1^2 + \Delta_2^2)}$$

$$\tau = \sqrt{T_n^2 + (T_1^2 + T_2^2) / \beta^2}$$

Assume that interface potential $\Phi(\lambda, D)$ has form

$$\Phi = \frac{1}{2} k_0 (1-D) \lambda^2$$

Hence $T_n = \frac{\partial \Phi}{\partial \Delta_n} = k_0 (1-D) \Delta_n$

$$T_\alpha = \frac{\partial \Phi}{\partial \Delta_\alpha} = \beta^2 k_0 (1-D) \Delta_\alpha$$

We also need an evolution equation for D

We have that $\underbrace{k_0(1-D)}_{= \sigma_{\max}/d_1} \lambda = \sigma_{\max} \left(1 - \frac{\lambda - d_1}{d_2 - d_1} \right)$

Take time derivative

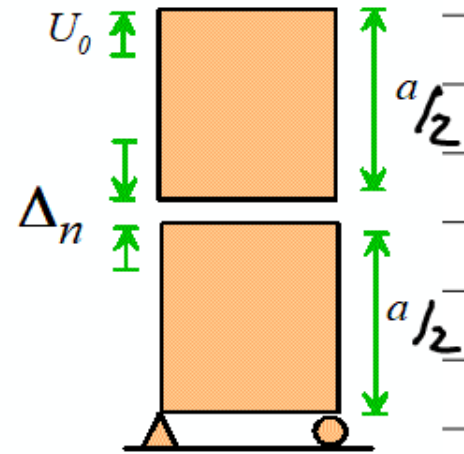
$$\frac{dD}{dt} = \begin{cases} 0 & \lambda < \frac{d_1 d_2}{(1-D)d_2 + Dd_1} \text{ or } d\lambda/dt < 0 \text{ or } D=1 \\ \left((1-D) + \frac{d_1}{d_2 - d_1} \right) \frac{1}{\lambda} \frac{d\lambda}{dt} & \text{Otherwise} \end{cases}$$

History dependent - D is a state variable

Viscous Regularization

Problem: Crack nucleation can be an unstable process

Example: Xu-Needleman interface between elastic blocks in uniaxial tension



$$\frac{\sigma}{\sigma_{\max}} = \frac{\Delta_n}{\delta_n} \exp\left(1 - \frac{\Delta_n}{\delta_n}\right)$$

$$U_0 = \frac{\sigma}{E} a + \Delta_n \Rightarrow \frac{U_0 E}{\sigma_{\max} a} = \frac{\sigma}{\sigma_{\max}} + \frac{\Delta_n}{\delta_n} \underbrace{\frac{E \delta_n}{\sigma_{\max} a}}_{\sim 1}$$

Stress - displacement relation

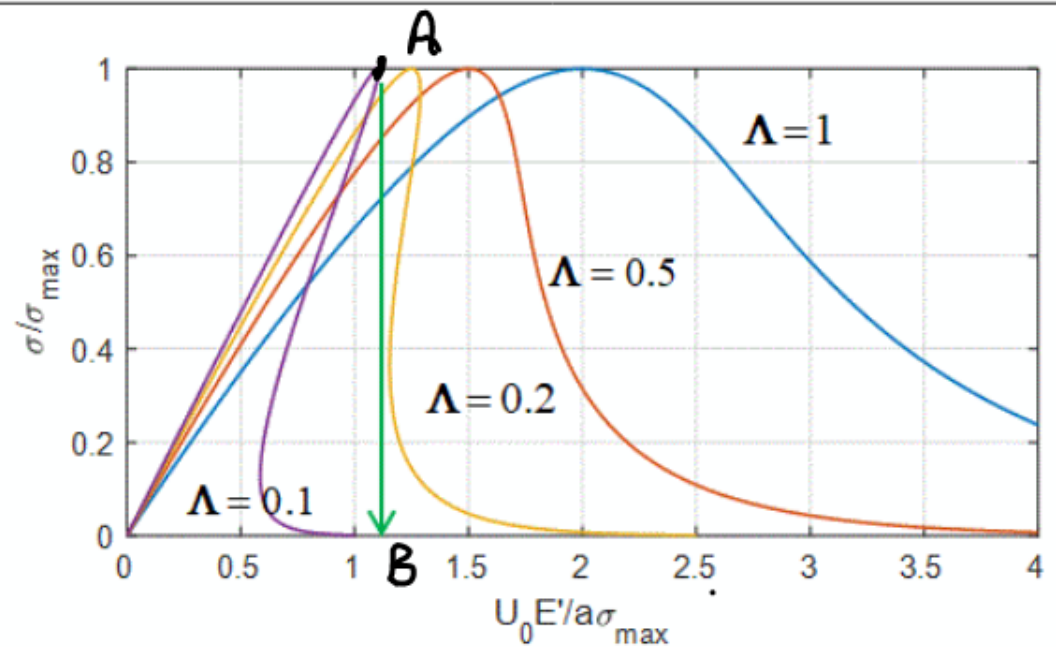
For large Λ , we see stable separation as U_0 increases

For small Λ we see a snap from A-B

For a displacement controlled static analysis there is no equilibrium path from A-B

Newton iterations will fail at A

We can fix this by adding some viscosity to interface



For example we could use

$$\left. \begin{aligned} T_n &= \frac{\partial \Phi}{\partial \Delta_n} + \eta \frac{d\Delta_n}{dt} \\ T_\alpha &= \frac{\partial \Phi}{\partial \Delta_\alpha} + \eta \frac{d\Delta_\alpha}{dt} \end{aligned} \right\} \eta \text{ is a small viscosity}$$

Choice of η dictated by computational cost & accuracy

FE implementation

Assume small strain linear elasticity

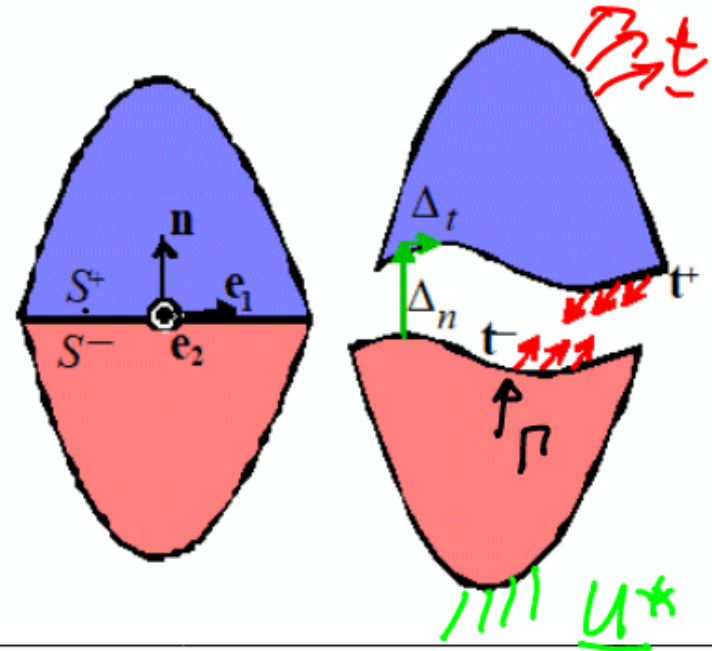
Governing Equations

$$\frac{\partial \sigma_{ij}}{\partial x_j} = 0 \quad \sigma_{ij} = C_{ijkl} \frac{\partial u_k}{\partial x_l}$$

Boundary Conditions : $\sigma_{ij} n_j = t_i^*$ on S_2
 $u_i = u_i^*$ on S_1

Interface Relations $n_i \sigma_{ij} n_j = T_n$ on Γ_-
 $= -T_n$ on Γ_+

$$e_i^\alpha \sigma_{ij} n_j = T_\alpha \text{ on } \Gamma_- \\ = -T_\alpha \text{ on } \Gamma_+$$



Constitutive eqs for cohesive zone $(T_n, T_\alpha) = f(\Delta_n, \Delta_\alpha)$

We need a new weak form for equilibria

Let η_i be a kinematically admissible displacement field

Note η_i may be discontinuous across Γ

Let η_i^+ , η_i^- denote η_i on two sides of Γ

Weak form of equilibrium

$$\int_{\mathcal{R}} \frac{\partial \sigma_{ij}}{\partial x_j} \eta_i \, dV = 0 \quad \Rightarrow \quad \int_{\mathcal{R}} \left\{ \frac{\partial}{\partial x_j} (\sigma_{ij} \eta_i) - \sigma_{ij} \frac{\partial \eta_i}{\partial x_j} \right\} dV = 0$$

Apply divergence theorem

$$\Rightarrow \int_{\Gamma^-} \sigma_{ij}^- n_j \eta_i^- dA - \int_{\Gamma_+} \sigma_{ij} n_j \eta_i^+ dA + \int_S \sigma_{ij} n_j \eta_i dA - \int_{\mathcal{R}} \sigma_{ij} \frac{d\eta_i}{dx_j} dV = 0 \quad \forall \text{ admiss } \eta_i^-$$

Note on Γ $\sigma_{ij}^- n_j = T_n n_i + T_\alpha e_i^\alpha$

$$\begin{aligned} \Rightarrow \int_{\mathcal{R}} \sigma_{ij} \frac{\partial \eta_i}{\partial x_j} dV + \int_{\Gamma} T_n (\eta_i^+ - \eta_i^-) n_i dA \\ + \int_{\Gamma} T_\alpha (\eta_i^+ - \eta_i^-) e_i^\alpha dA - \int_{S_2} t_i^* \eta_i dA = 0 \quad \forall \text{ admiss } \eta_i \end{aligned}$$

This looks like usual PVW but with some additional terms on interface

Introduce usual interpolations inside volume

$$u_i = N^a u_i^a \quad \eta_i = N^a \eta_i^a$$

Account for displacement discontinuity using cohesive zone elements

Generate meshes inside two solids joined by interface with coincident nodes on Γ

