

HOLOGRAPHIC DOUBLE DIFFRACTIVE PRODUCTION OF HIGGS AND THE ADS GRAVITON (POMERON)

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(with Richard Brower and Marko Djuric)
Paris-2011

R.C. Brower, Marko Djurić, Chung-I Tan
“Holographic Double Diffractive Higgs Production”
in preparation

Brower, Polchinski, Strassler, Tan (BPST) The Pomeron and Gauge/String Duality (2006)

References:

- R. Brower, J. Polchinski, M. Strassler, and C-I Tan, “The Pomeron and Gauge/String Duality”, hep-th/0603115.
- R. Brower, M. Strassler, and C-I Tan, hep-th/0707.2408; hep-th/0710.4378.
- R. Brower, M. Djuric, and C-I Tan, arXiv:0812.0354.
- R. Brower, M. Djuric, I. Sarcevic and C-I Tan, “DIS and Gauge/String Duality”, arXiv:1007.2259
- R. Brower, M. Djuric, and C-I Tan, “Diffractive Higgs Production and Gauge/String Duality”, (in preparation)

Other related work, e.g.,

L. Cornalba, et al., (hep-th/0710.5480),

- Y. Hatta, E. Iancu, and A. H. Mueller, (hep-th/0710.2148),
- E. Levin, et al. (arXiv:0811.3586) and (arXiv:0902.3122).
- Many others, (e.g., Kovchegov’s talk, more recent work by Y. Hatta, et al., etc.)

Executive Summary:

Gauge/String Duality (AdS/CFT)  2-GLUONS \simeq GRAVITON

Goals:

- ◆ Generality of “Pomeron” in QCD beyond perturbation theory,

- ◆ Unification of Soft and Hard Physics
- ◆ New phenomenology based on “Large Pomeron intercept”, e.g., DIS at small-x: [DGLAP vs Pomeron \(talk by Marko Djuric\)](#)
- ◆ Diffractive Higgs Production

EXACT Weak/Strong EQUIVALENCE

N = 4 SYM is IIB string theory in AdS5 x S5

$$ds^2 = e^{2A(z)} [-dx^+ dx^- + dx_\perp dx_\perp + dz dz] + ds^2(Y^5)$$

for AdS_5 $e^{2A(z)} = \frac{R^2}{z^2}$

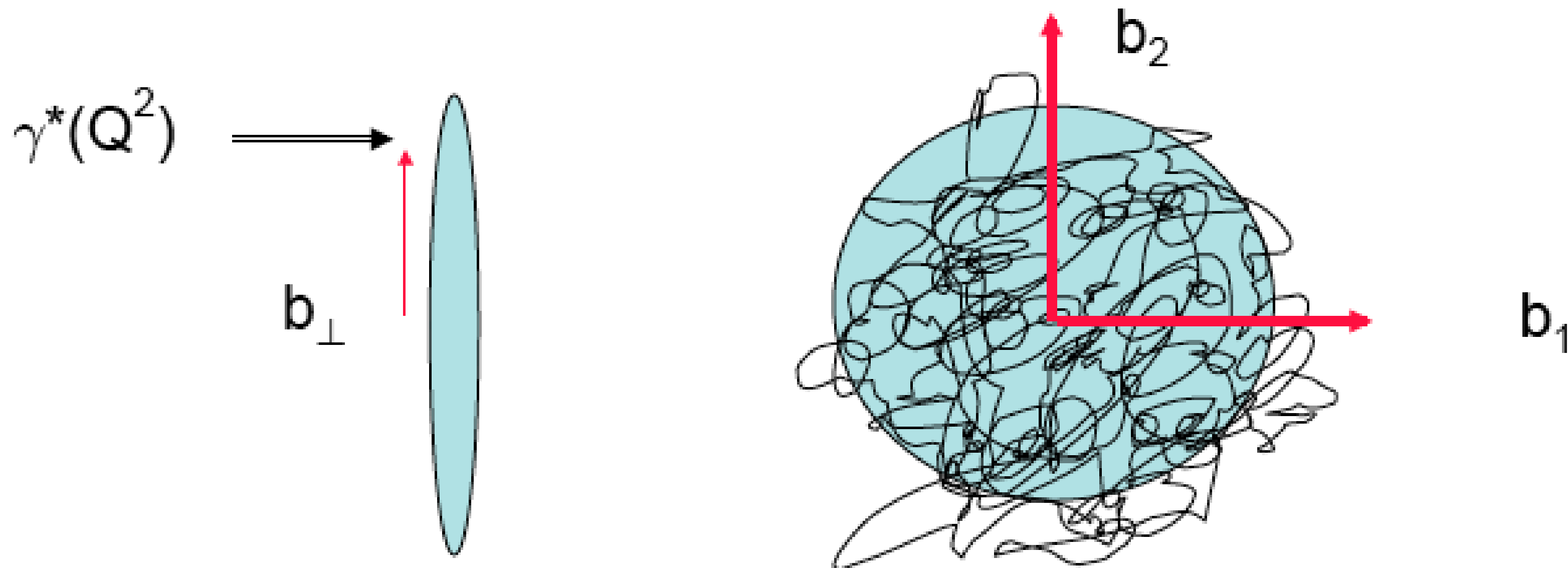
$R^2 d^2\Omega_5$

Overwhelming evidence for $\mathcal{N} = 4$ SYM
in the 'tHooft planar limit $N_c \rightarrow \infty$
at fixed $\lambda = g^2 N_c$

$$\langle e^{\int d^4x \phi_i(x) \mathcal{O}_i(x)} \rangle_{CFT} = \mathcal{Z}_{string} [\phi_i(x, z)|_{z \sim 0} \rightarrow \phi_i(x)]$$

QCD EMERGENCE OF 5-DIM ADS

“Fifth” co-ordinate is size z / z' of proj/target



5 kinematical Parameters:

2-d Longitudinal

$$p^{\pm} = p^0 \pm p^3 \simeq \exp[\pm \log(s/\Lambda_{qcd})]$$

2-d Transverse space:

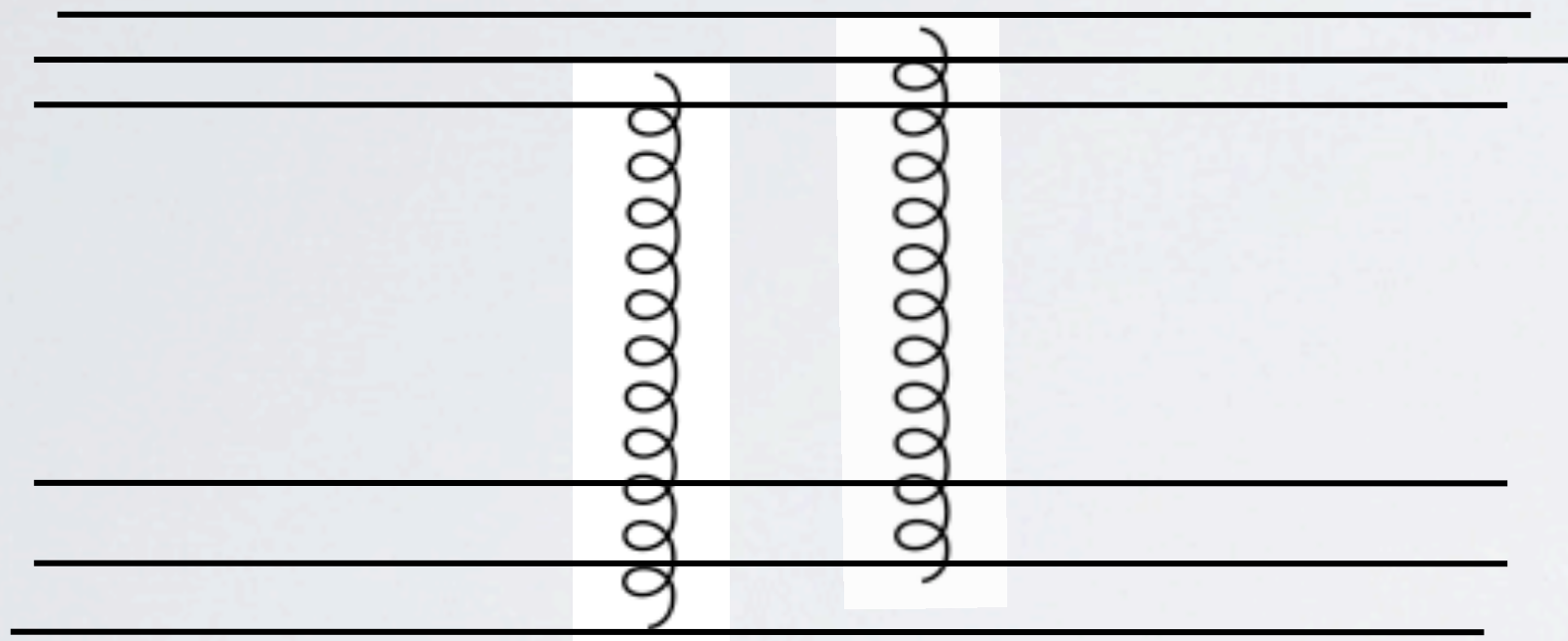
$$x'_{\perp} - x_{\perp} = b_{\perp}$$

1-d Resolution:

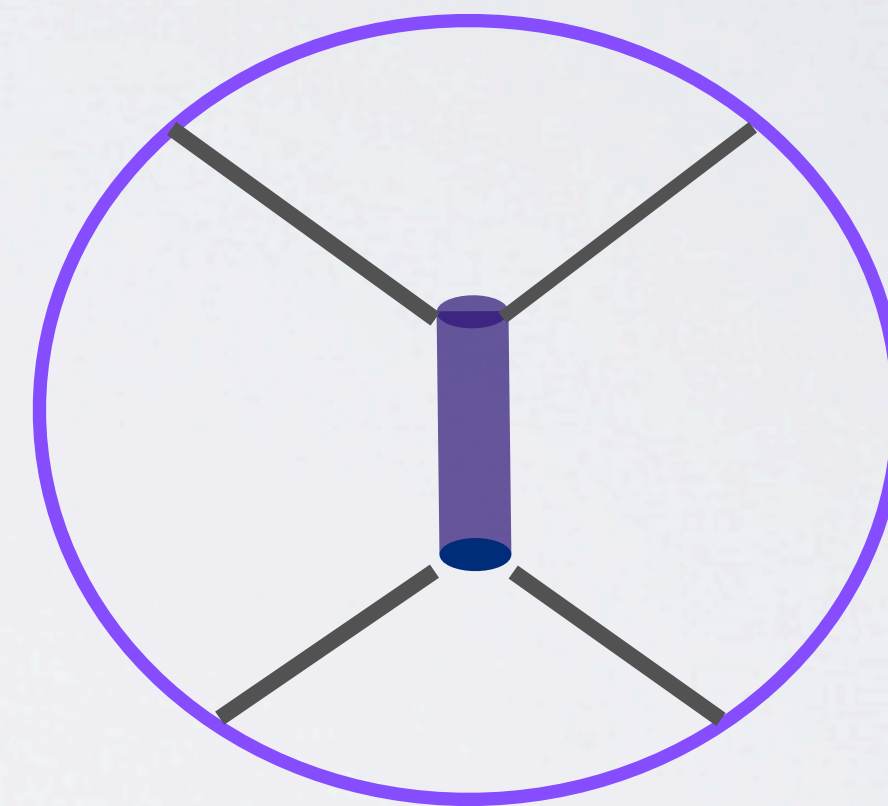
$$z = 1/Q \quad (\text{or } z' = 1/Q')$$

WHAT IS THE BARE POMERON ? LEADING 1/N TERM CYLINDER EXCHANGE

WEAK: TWO GLUON \Leftrightarrow STRONG: ADS GRAVITON



$$J_{cut} = 1 + 1 - 1 = 1$$



$$J = 2$$

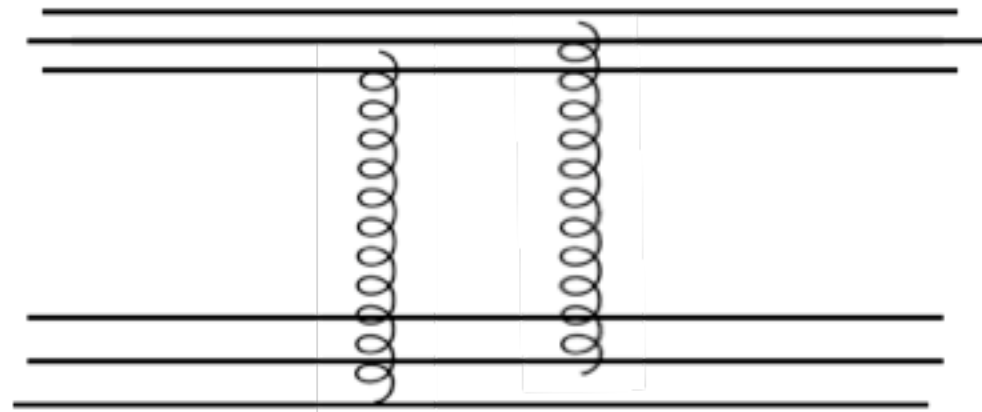
$$S = \frac{1}{2\kappa^2} \int d^4x dz \sqrt{-g(z)} \left(-\mathcal{R} + \frac{12}{R^2} + \frac{1}{2} g^{MN} \partial_M \phi \partial_N \phi \right)$$

F.E. Low. Phys. Rev. D 12 (1975), p. 163.
S. Nussinov. Phys. Rev. Lett. 34 (1975), p. 1286.

AdS Witten Diagram: Adv.
Theor. Math. Physics 2 (1998)253

Two gluon exchange (Low-Nussinov Pomeron!)

- $J_{\text{cut}} = 2(J-1) + 1 = 1$

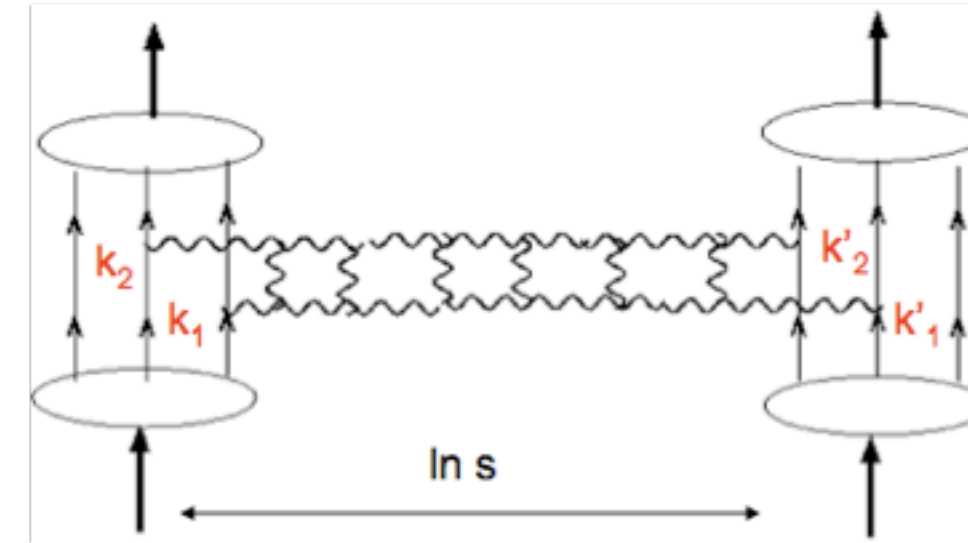


F.E. Low. Phys. Rev. D 12 (1975), p. 163.
S. Nussinov. Phys. Rev. Lett. 34 (1975), p. 1286.

BFKL: Balitsky & Lipatov; Fadin, Kuraev, Lipatov '75

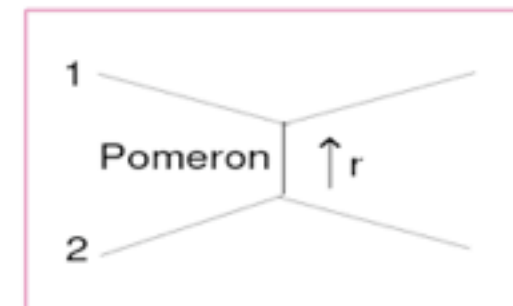
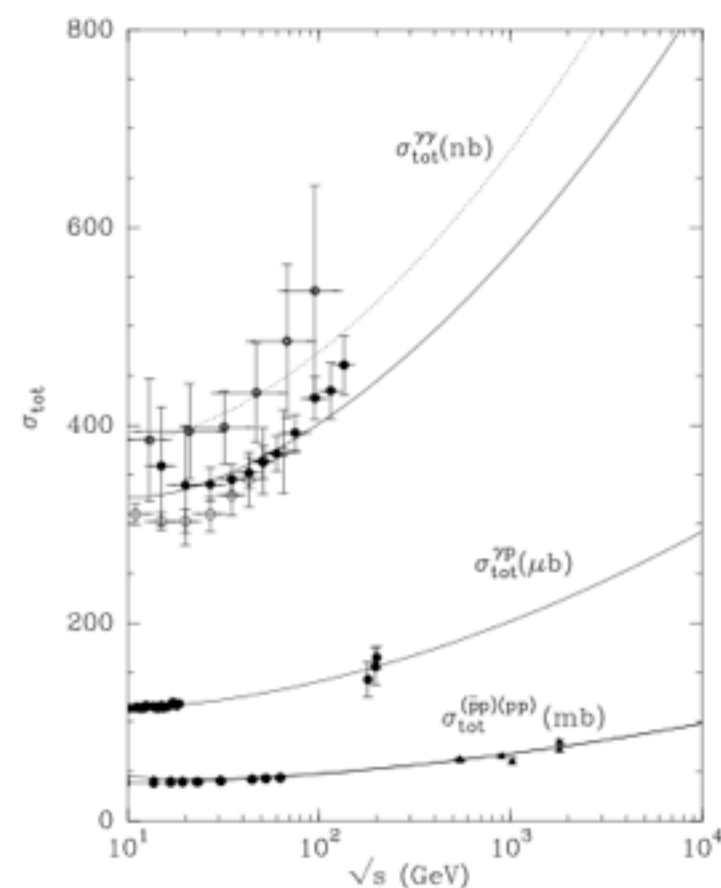
$$t = -(k_1 + k_2)^2 \rightarrow$$

$$\lambda = g^2 N_c \sim 0$$



- Sum diagrams 1st order in $g^2 N_c$ & all orders $(g^2 N_c \log s)^n$
- BKFL equation for 2 "reggized" gluon ladder is $L = 2$ $SL(2, C)$ spin chain to one loop order.
- Accidentally "planar" diagrams (e.g. $N_c = 1$) and conformal.

Total Cross Sections



$$\mathcal{A} \sim s^{J(t)} = s^{\alpha(0) + \alpha' t}$$

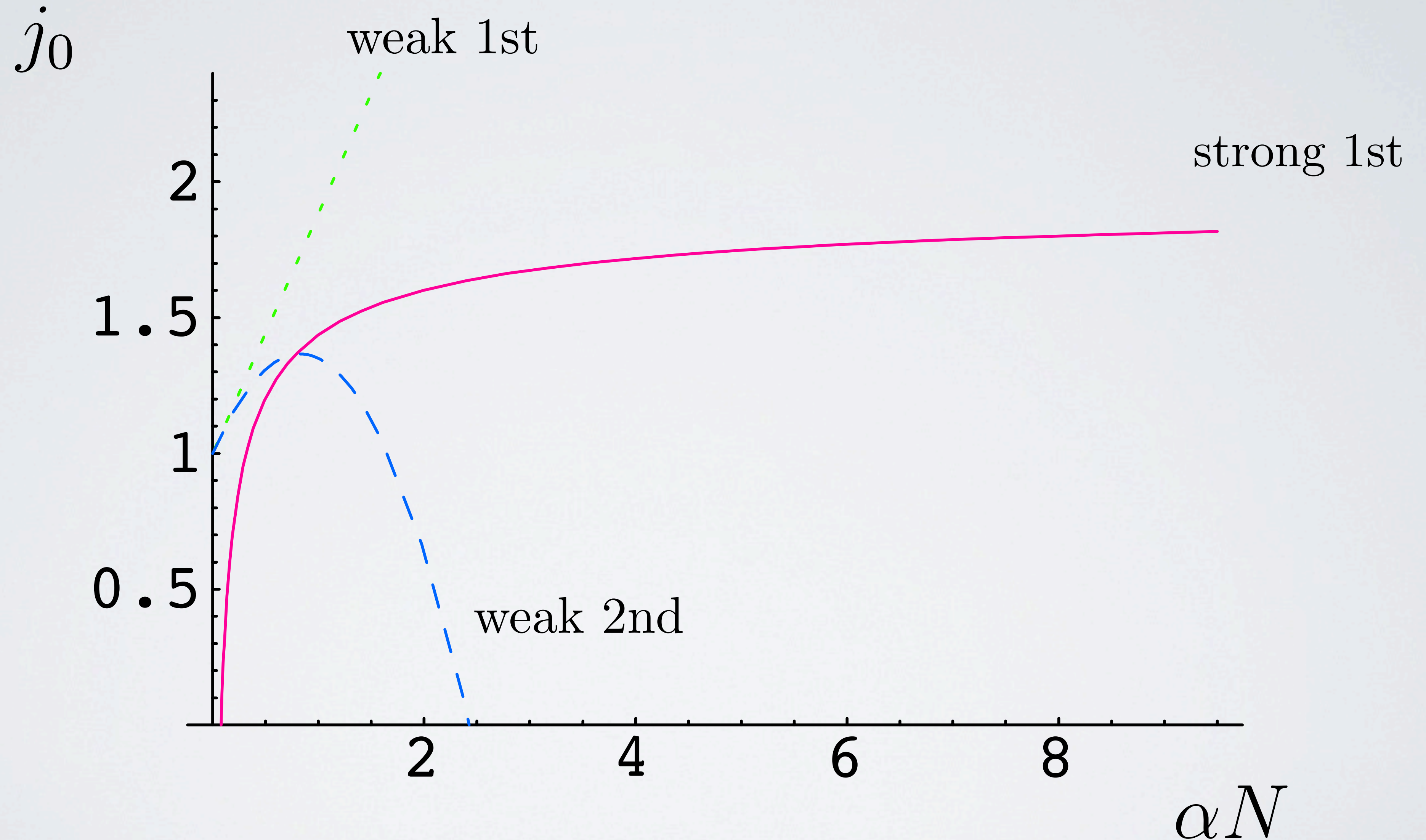
$$\sigma_{\text{total}} \sim \mathcal{A}(s, 0)/s \sim S^{J(0)-1} \sim s^{\alpha(0)-1}$$

$$\alpha(0) > 1$$

(IR) Pomeron as Closed String??

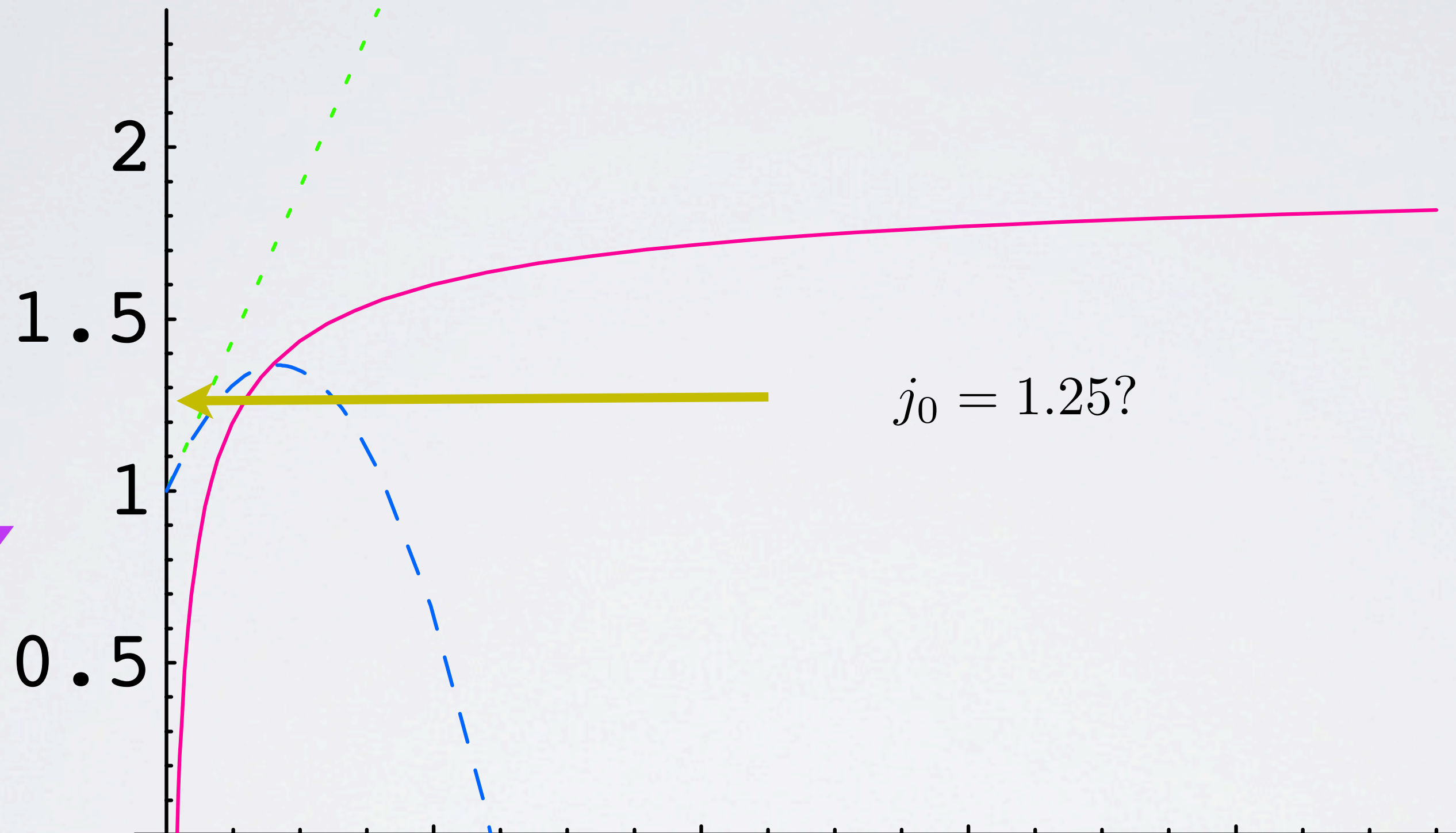
Using AdS/CFT to obtain a strong coupling, i.e., non-perturbative, treatment of high energy forward scattering

$\mathcal{N} = 4$ Strong vs Weak $g^2 N_c$



$\mathcal{N} = 4$ Strong vs Weak $g^2 N_c$

j_0



$j_0 = 1$

Two
Gluon

$j_0 = 1.25?$

$j_0 = 2$

Graviton

BFKL

QCD?

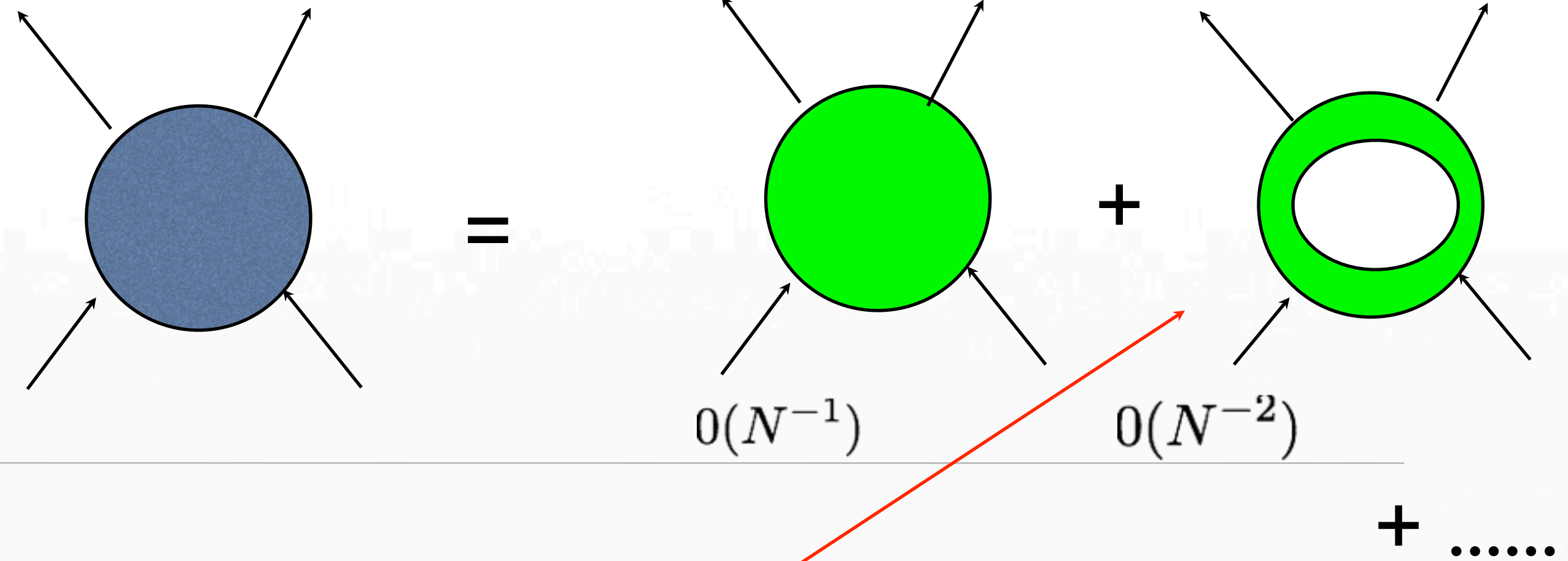
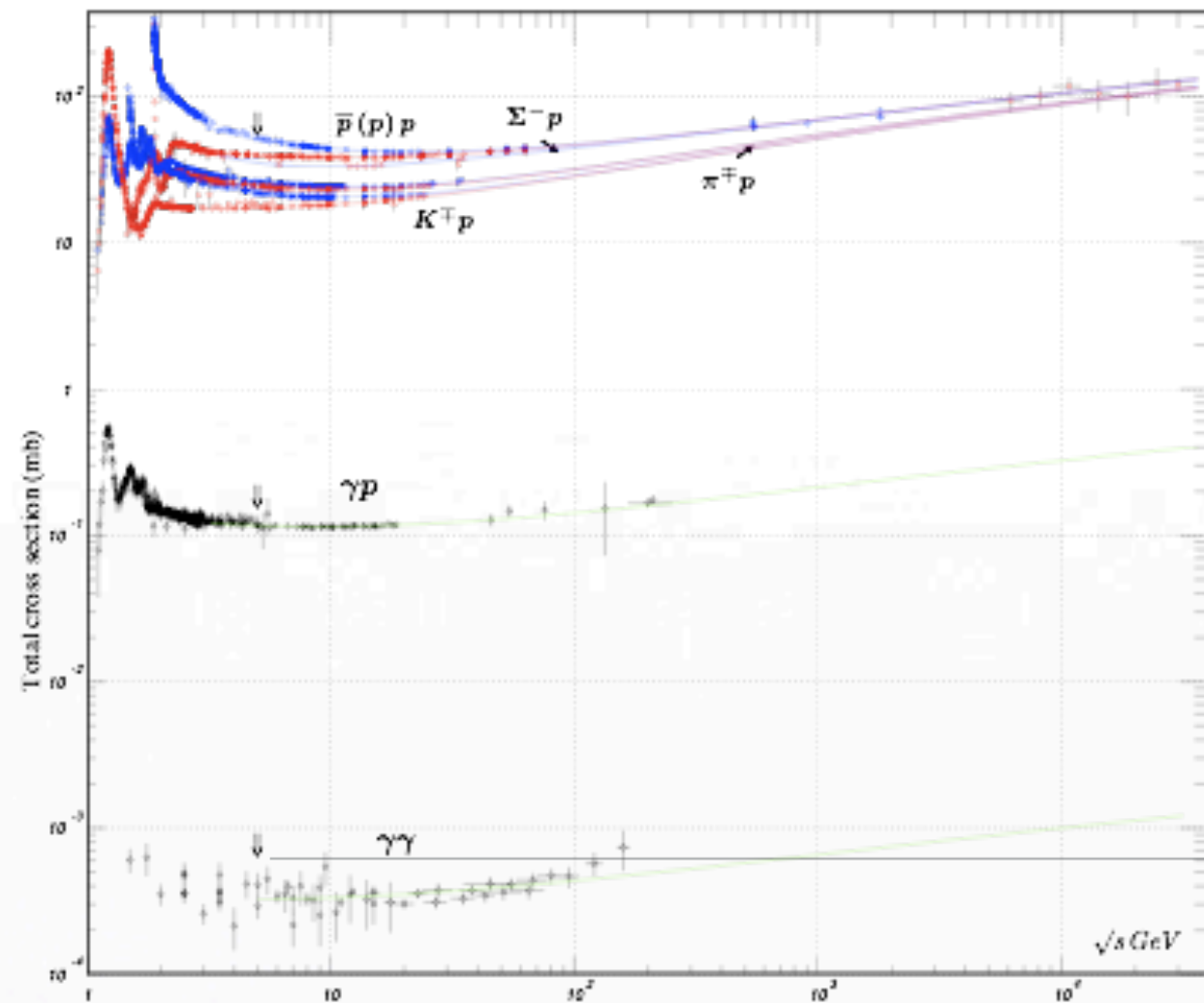
BPST

$$j_0 = 1 + \ln(2)g^2 N_c / \pi^2$$

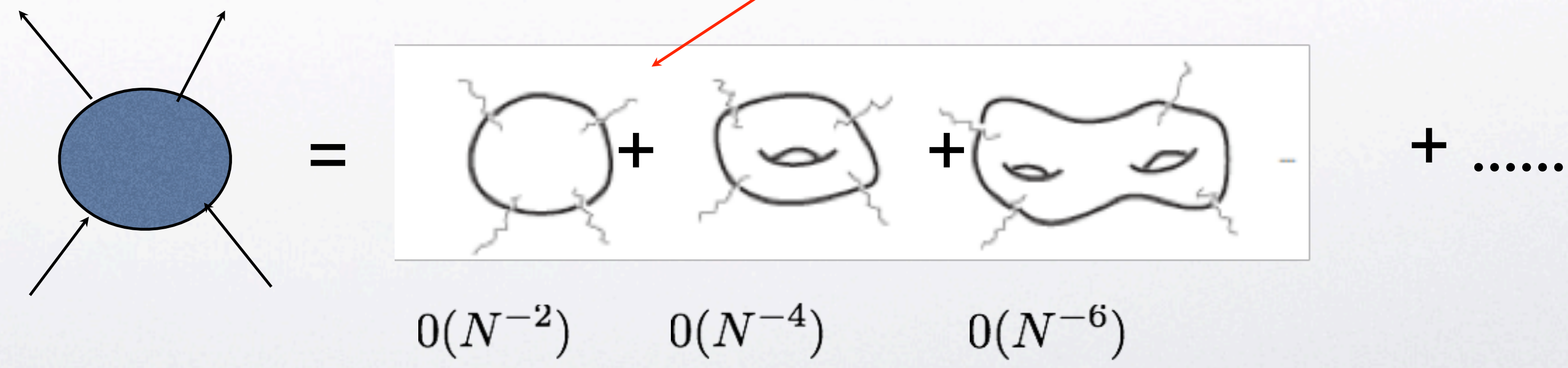
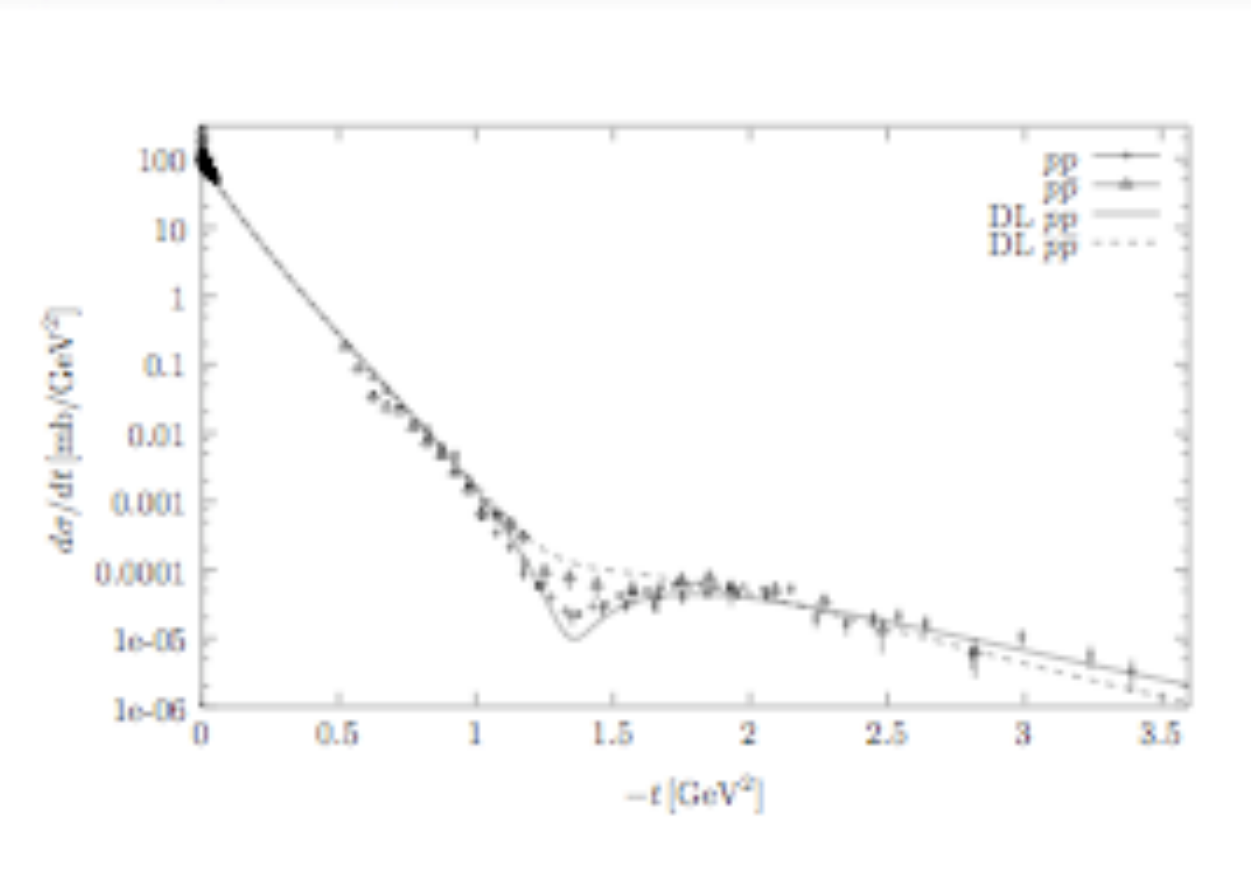
$$j_0 = 2 - 2/\sqrt{g^2 N_c}$$



Open-string Scattering



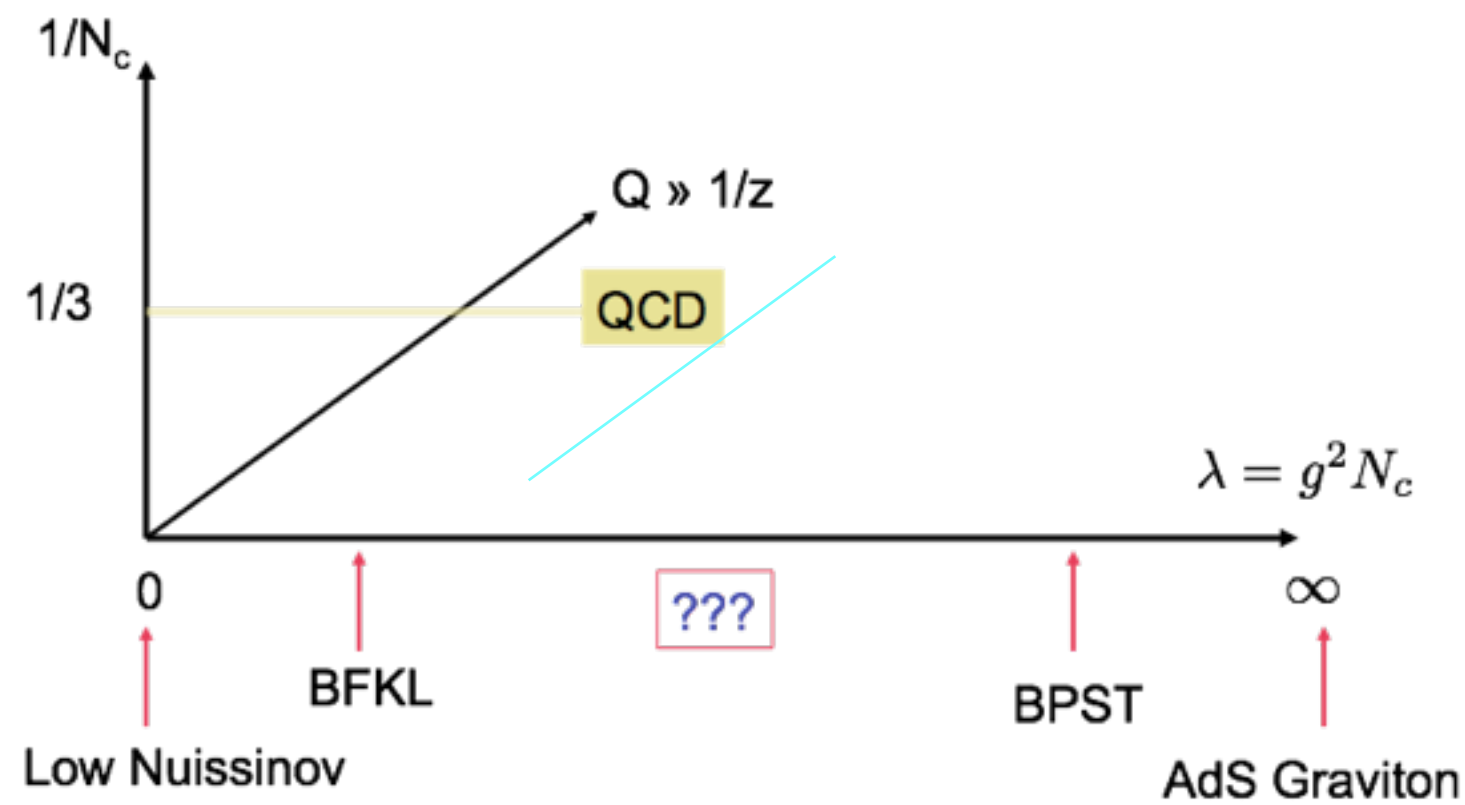
Closed-string Scattering



In this talk, will focus on “closed string” exchanges only. For DIS, applicable at $x \ll 1$.

Gauge/String Duality: QCD at Strong Coupling

Pomeron Parameter Space



- $C=+1$: Pomeron \Leftrightarrow Graviton:

$$\alpha_0^{(+)} = 2 - 2/\sqrt{\lambda} + O(1/\lambda)$$

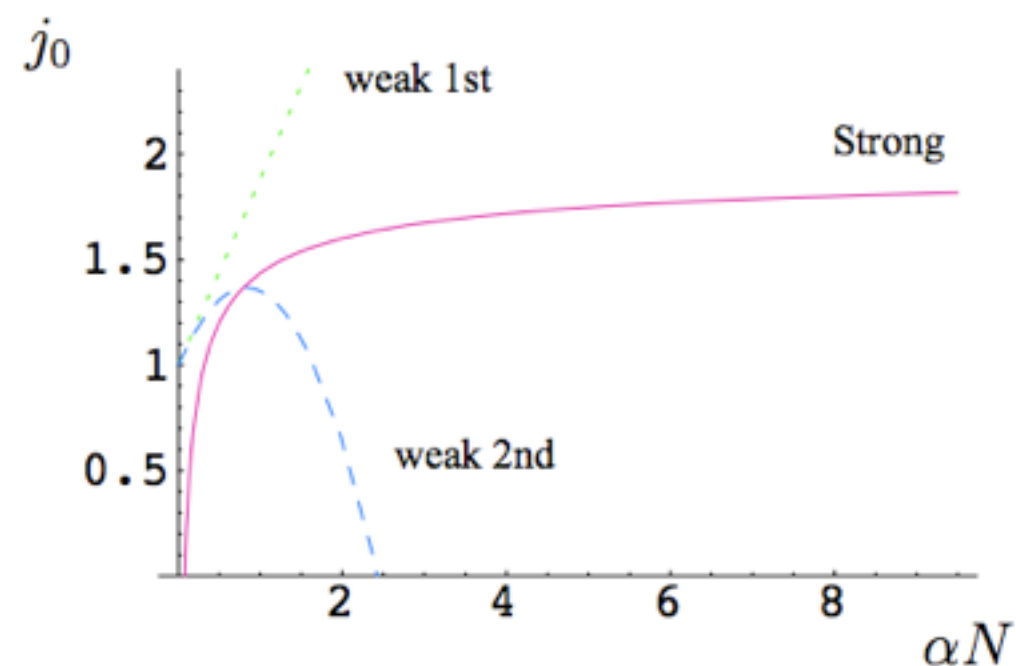
(*symmetric tensor* : $g_{\mu\nu}$)

- $C=-1$: Odderon \Leftrightarrow Kalb-Ramond

$$\alpha_0^{(-)} = 1 - m_{ads}^2/2\sqrt{\lambda} + O(1/\lambda)$$

(*anti - symmetric tensor* : $b_{\mu\nu}$)

$\mathcal{N} = 4$ Strong vs Weak BFKL



- Many New Questions:

Comparison of strong vs weak coupling kernel at t=0

Strong Coupling:

$$\mathcal{K}(r, r', s) = \frac{s^{j_0}}{\sqrt{4\pi\mathcal{D}\ln s}} e^{-(\ln r - \ln r')^2 / 4\mathcal{D}\ln s}$$

Diffusion in “warped co-ordinate”

$$j_0 = 2 - \frac{2}{\sqrt{g^2 N}} + O(1/g^2 N)$$

$$\mathcal{D} = \frac{1}{2\sqrt{g^2 N}} + O(1/g^2 N)$$

Weak Coupling:

$$K(s, k_{\perp}, k'_{\perp}) \approx \frac{s^{\alpha(0)-1}}{\sqrt{\pi\ln s}} e^{-[(\ln k'_{\perp} - \ln k_{\perp})^2 / 4\mathcal{D}\ln s]}$$

$$j_0 = 1 + \ln(2)g^2 N / \pi^2$$

$$\mathcal{D} = \frac{14\zeta(3)}{\pi} g^2 N / 4\pi^2.$$

ELASTIC, DIS, DOUBLE HIGGS: ADS BUILDING BLOCKS

$$A(s, t) = g_0^2 \int d^3\mathbf{b} d^3\mathbf{b}' e^{i\mathbf{q}_\perp \cdot (\mathbf{x} - \mathbf{x}')} \Phi_{13}(z) \mathcal{K}(s, \mathbf{x} - \mathbf{x}', z, z') \Phi_{24}(z')$$

$$\sigma_T(s) = \frac{1}{s} \text{Im} A(s, 0)$$

$$d^3\mathbf{b} \equiv dz d^2x_\perp \sqrt{-g(z)} \quad \text{where} \quad g(z) = \det[g_{nm}] = -e^{5A(z)}$$

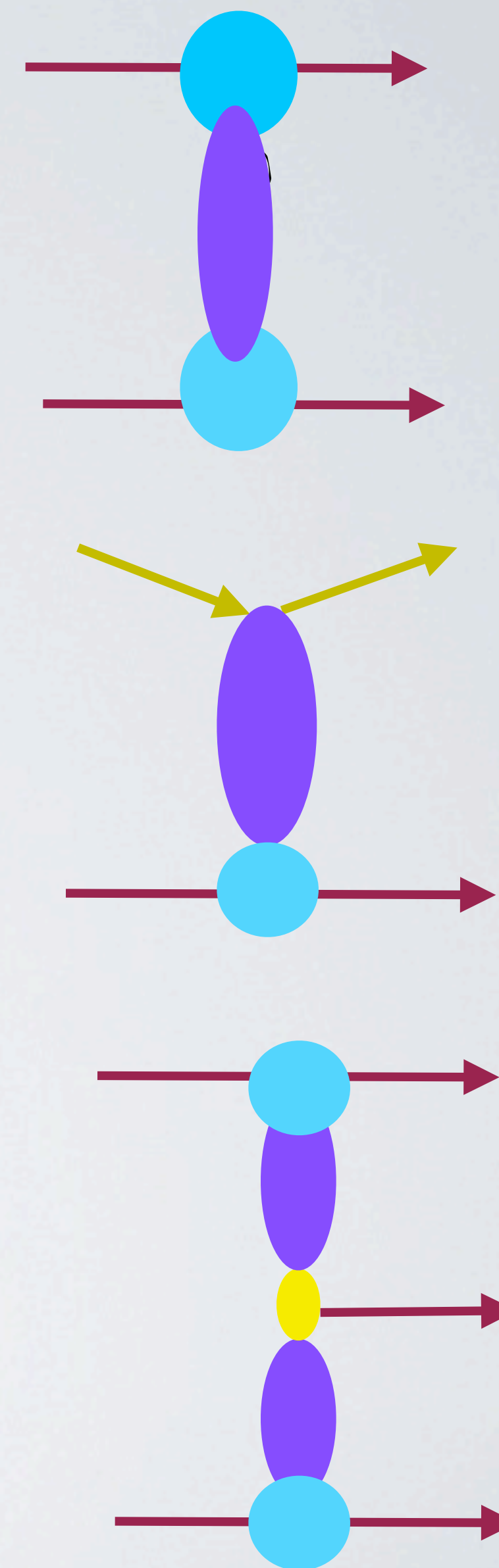
for $F_2(x, Q)$

$$\Phi_{13}(z) \rightarrow \Phi_{\gamma^* \gamma^*}(z, Q) = \frac{1}{z} [Qz]^4 (K_0^2(Qz) + K_1^2(Qz))$$

For Double Diffractive Higgs

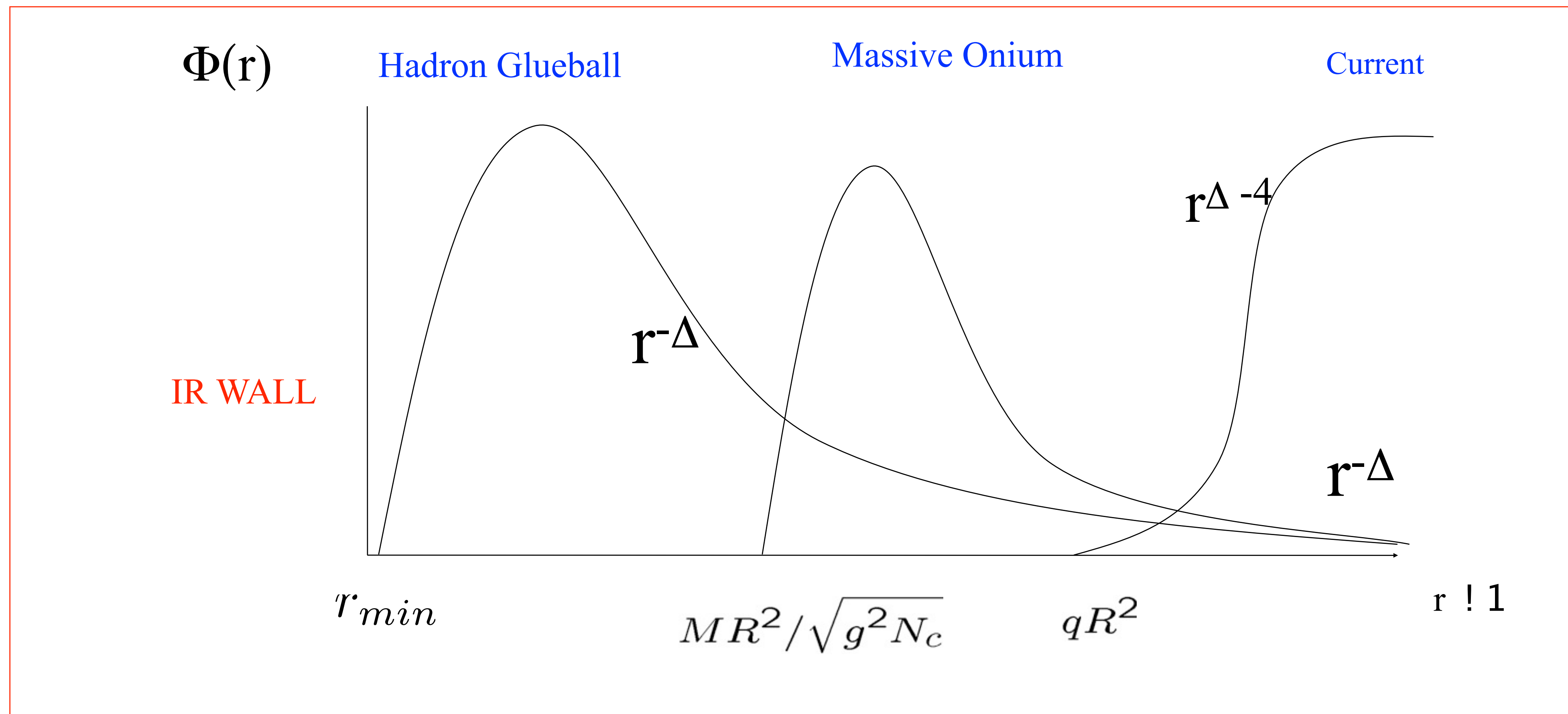
$$A(s_1, s_2, s, t_1, t_2) = \Phi_{13} * \mathcal{K}_1 * \mathcal{V}_H * \mathcal{K}_2 * \Phi_{24}$$

$$\mathcal{V}_H \rightarrow V_H \Phi_H = V_H (m_H z)^2 K_2(m_H z)$$



- Universality:

By choosing wave functions, Φ , can treat DIS, Higgs Production, Proton-Proton, etc., on equal footing.

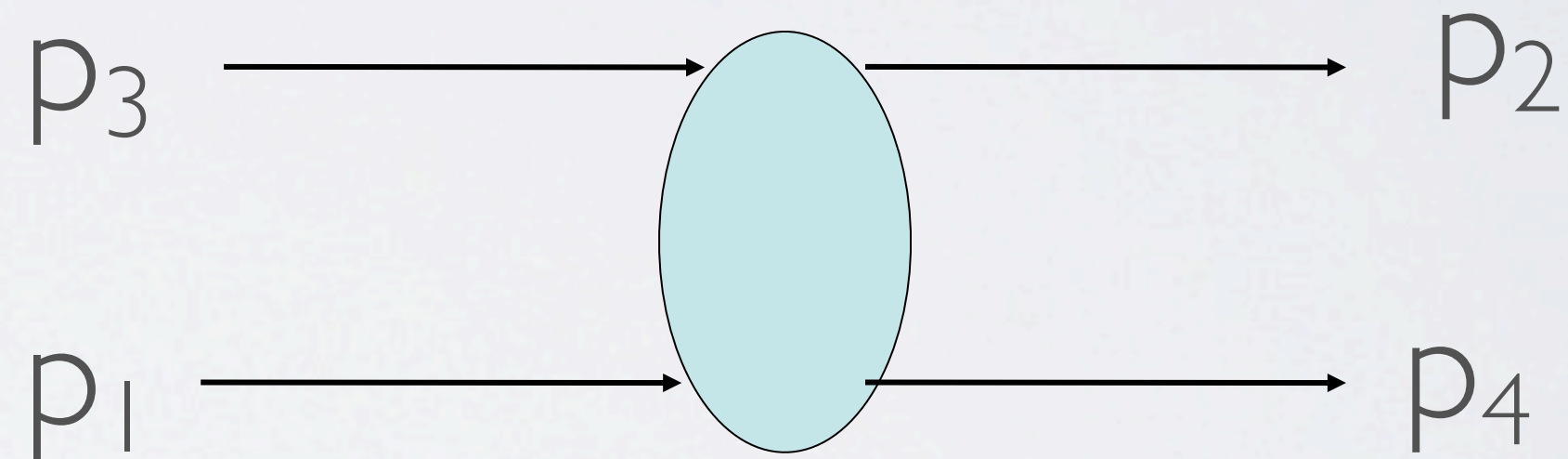


CLASSIC REGGE DIFFUSION IMPACT PARAMETER

$$A(s, t) \sim s^{\alpha' t + j_0} \rightarrow a(s, b_{\perp}) \sim s^{j_0} \exp[-b_{\perp}^2 / 4\alpha' \log(s)]$$

Diffusion in transverse (impact) space
rapidity = "time"

$$s = (p_1 + p_3)^2$$



$$t = (p_1 + p_2)^2$$

$$A(s, t = -q_{\perp}^2) = \int \frac{d^2 b_{\perp}}{4\pi^2} e^{-i q_{\perp} \cdot b_{\perp}} a(s, b_{\perp})$$

SL(2,C) ADS3 J-PLAN SCHRÖDINGER EQU.

$$M_{+-}\psi = [-\partial_u^2 + e^{-2u}\nabla_{\perp}] \psi = 2\sqrt{g^2 N_c} \partial_y \psi$$

Eigenvalues:

$$\nu^2$$

$$t = -q_{\perp}^2$$

$$E = 2 - j$$

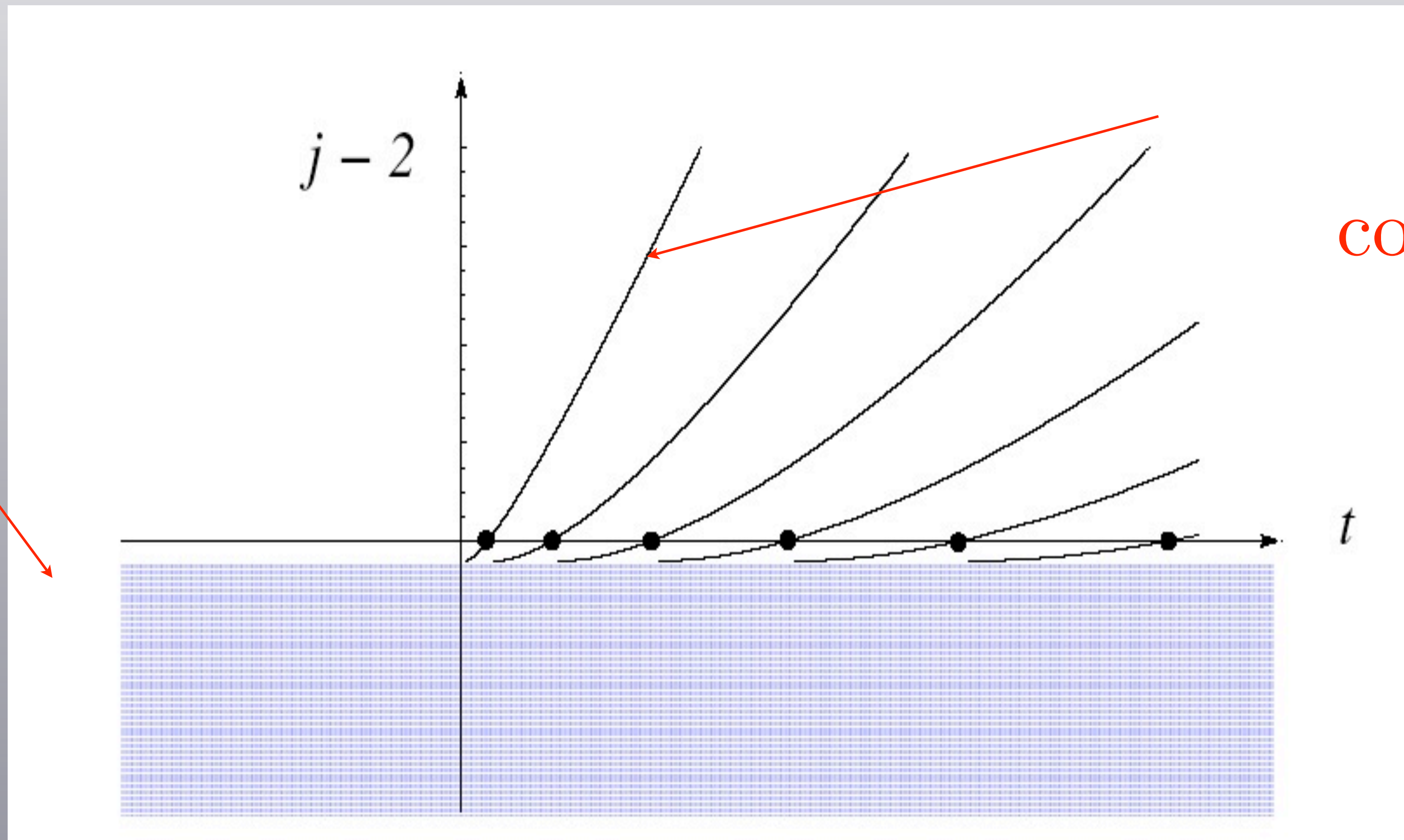
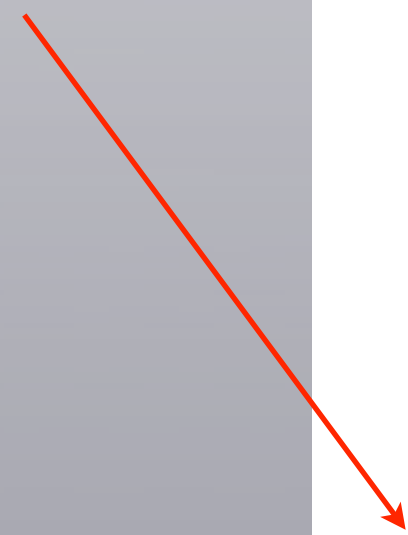
$$G_j(t, z, z') \sim \int q dq \frac{J_{(\Delta(j)-2)}(qz) J_{(\Delta(j)-2)}(qz')}{q^2 - t}$$

$$(\Delta(j) - 2)^2 = 2\sqrt{g^2 N_c}(j - j_0)$$

Unified Hard (conformal) and Soft (confining) Pomeron

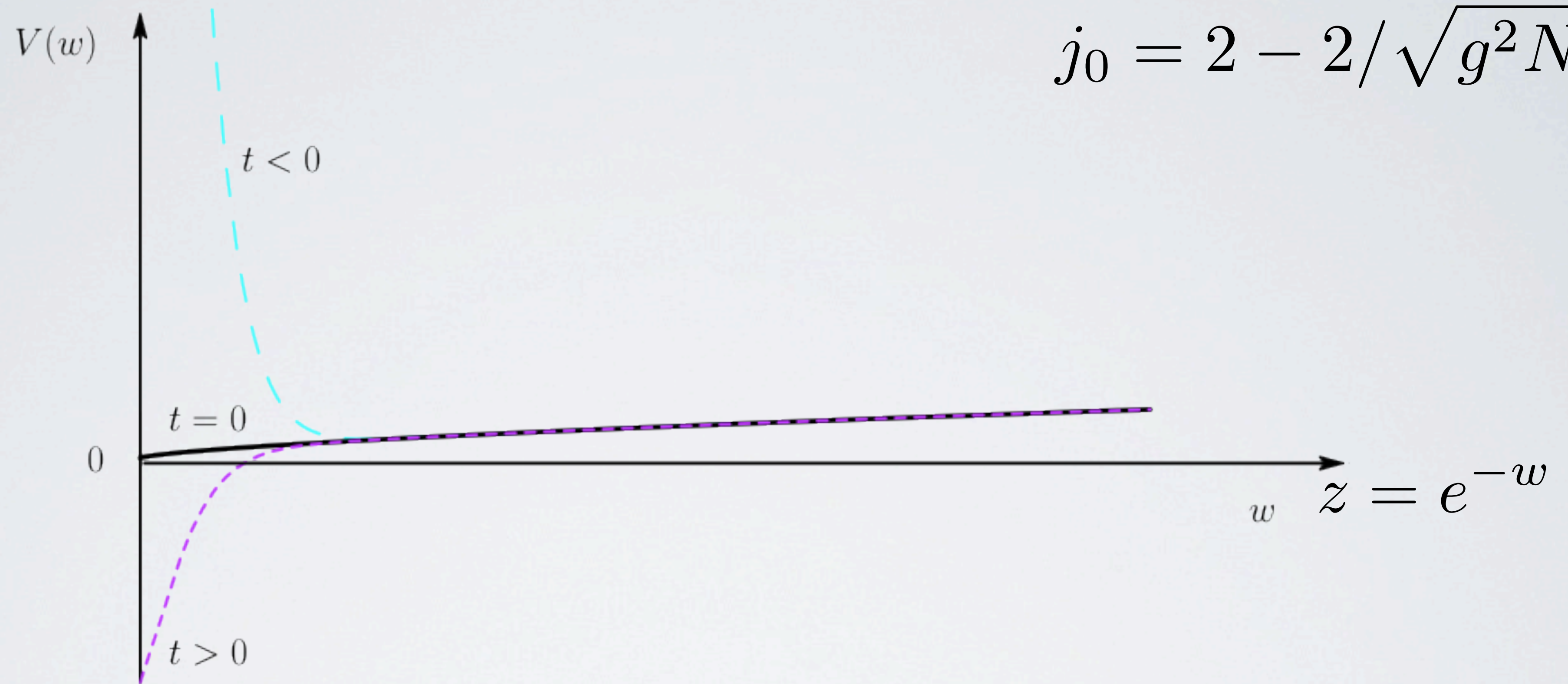
At finite λ , due to Confinement in AdS, *at* $t > 0$ asymptotical linear Regge trajectories

diffusion



confinement

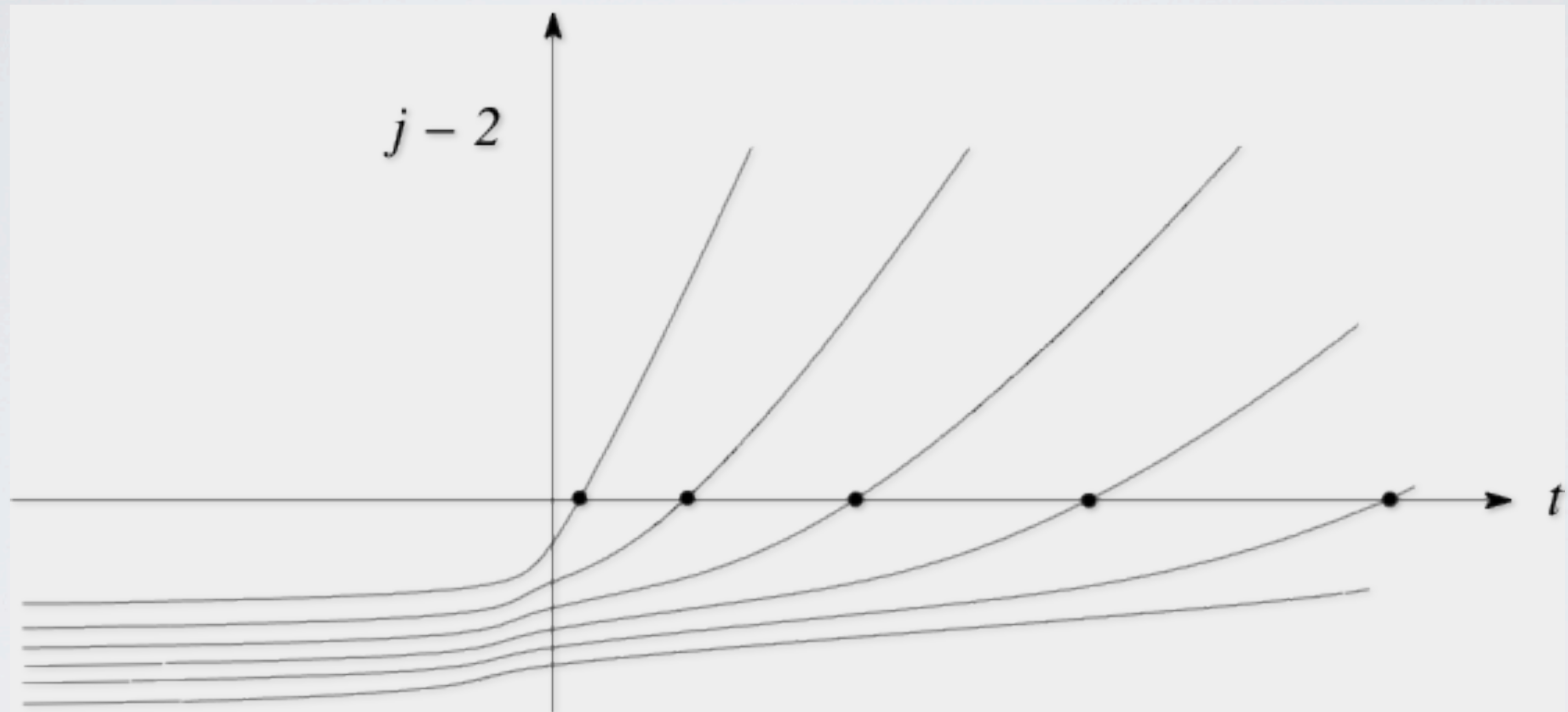
ADS J-PLANE SCHEODINGER EQ



$$j_0 = 2 - 2/\sqrt{g^2 N}$$

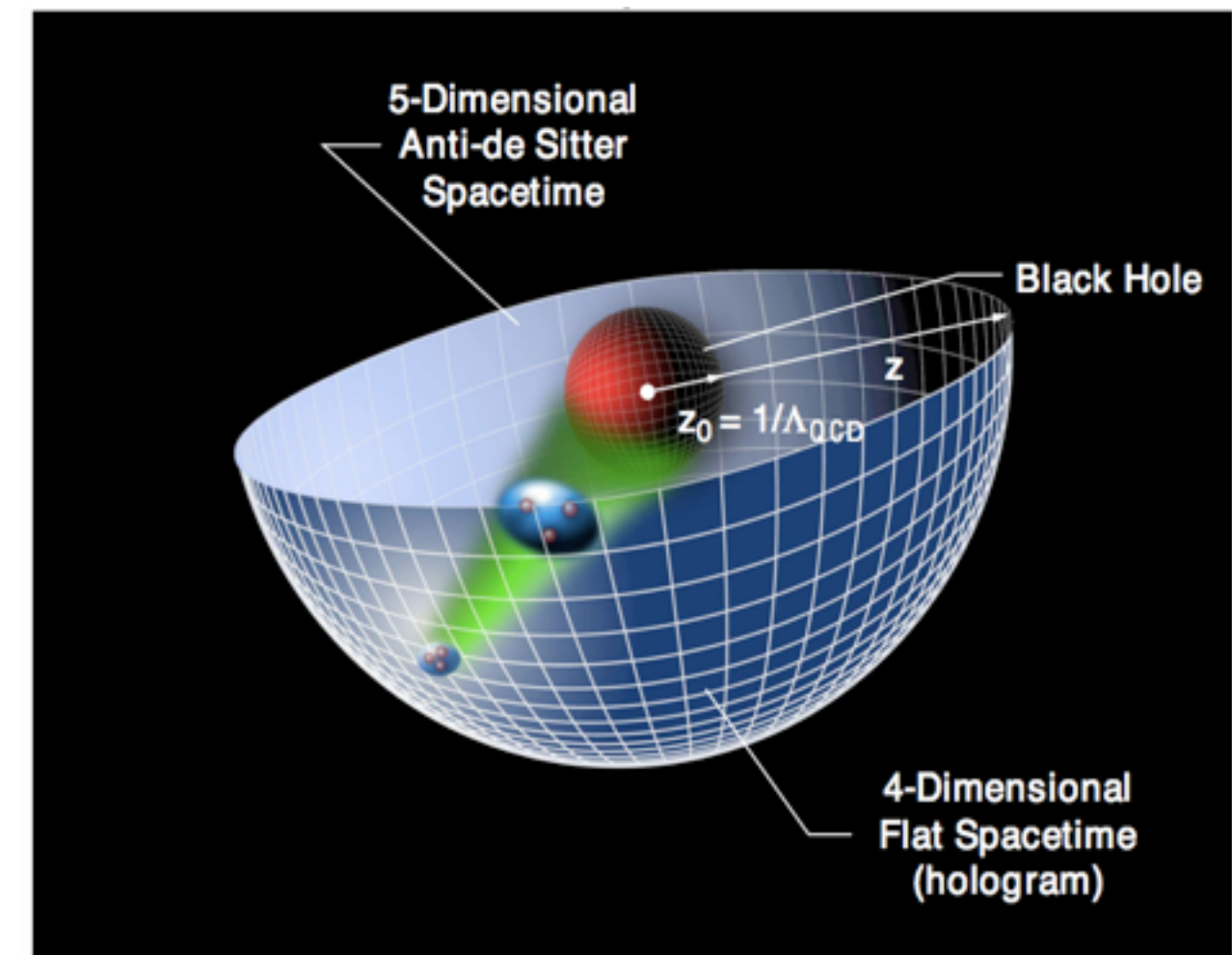
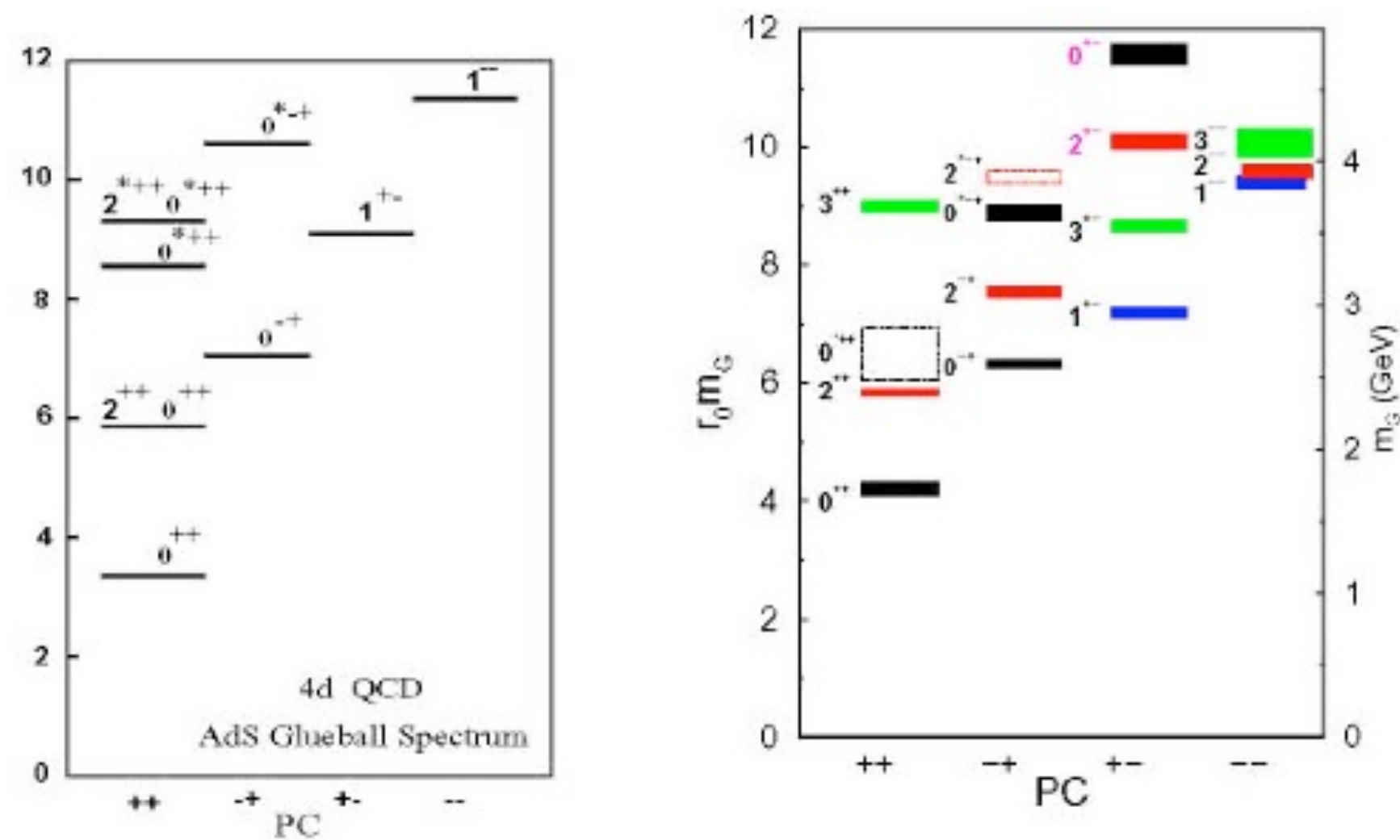
$$\left[-\partial_w^2 - te^{-2w} - cw \right] / 2 + \sqrt{g^2 N} (j - j_0) G_j(t, z, z') = \delta(w - w')$$

(STRONG) RUNNING



Confinement Deformation: Glueball Spectrum

$$(\lambda = \infty)$$



Four-Dimensional Mass:

$$E^2 = (p_1^2 + p_2^2 + p_3^2) + M^2$$

5-Dim Massless Mode:

$$0 = E^2 - (p_1^2 + p_2^2 + p_3^2 + p_r^2)$$

AdS/CFT \implies

In gauge theories with string-theoretical dual descriptions, the Pomeron emerges **unambiguously**.

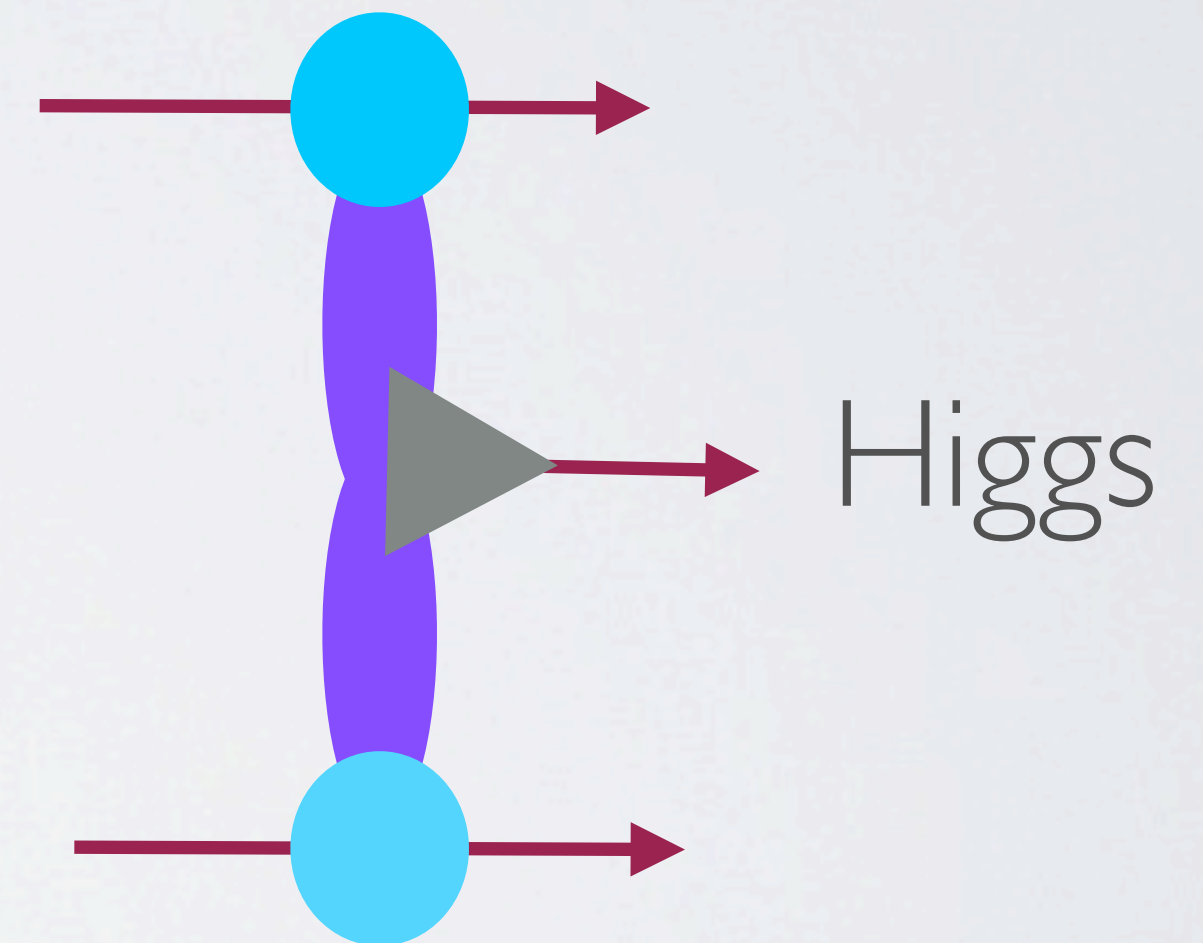
Pomeron can be associated with a Reggeized Graviton.

Both the IR (soft) Pomeron and the UV (BFKL) Pomeron are dealt in a unified single step.

R. Brower, J. Polchinski, M. Strassler, and C-I Tan, "The Pomeron and Gauge/String Duality", (hep-th/0603115.)

DOUBLE DIFFRACTIVE HIGGS PRODUCTION

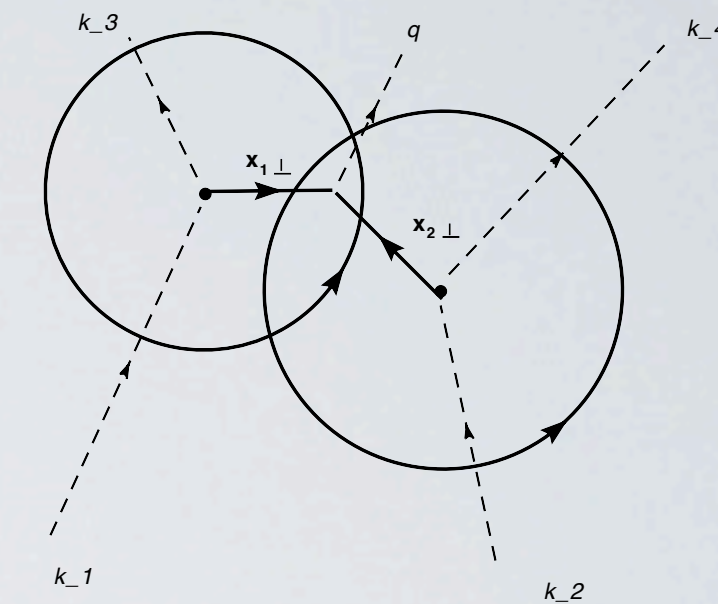
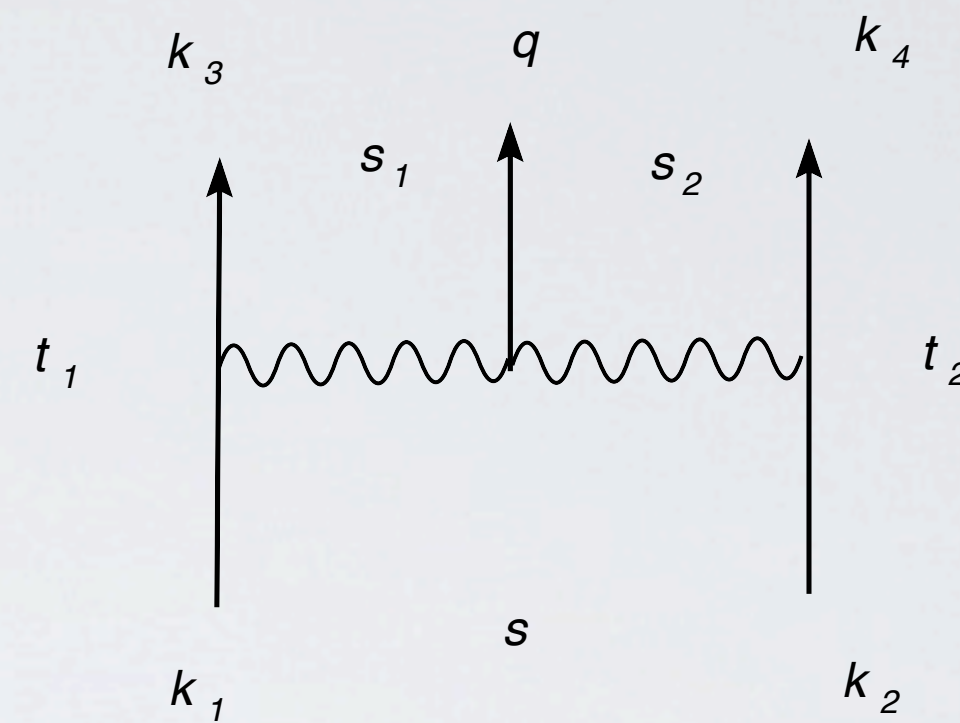
- Kinematics:
- Basic Building Block:
- New Feature:



KINEMATICS:

- Longitudinal Phase Space:

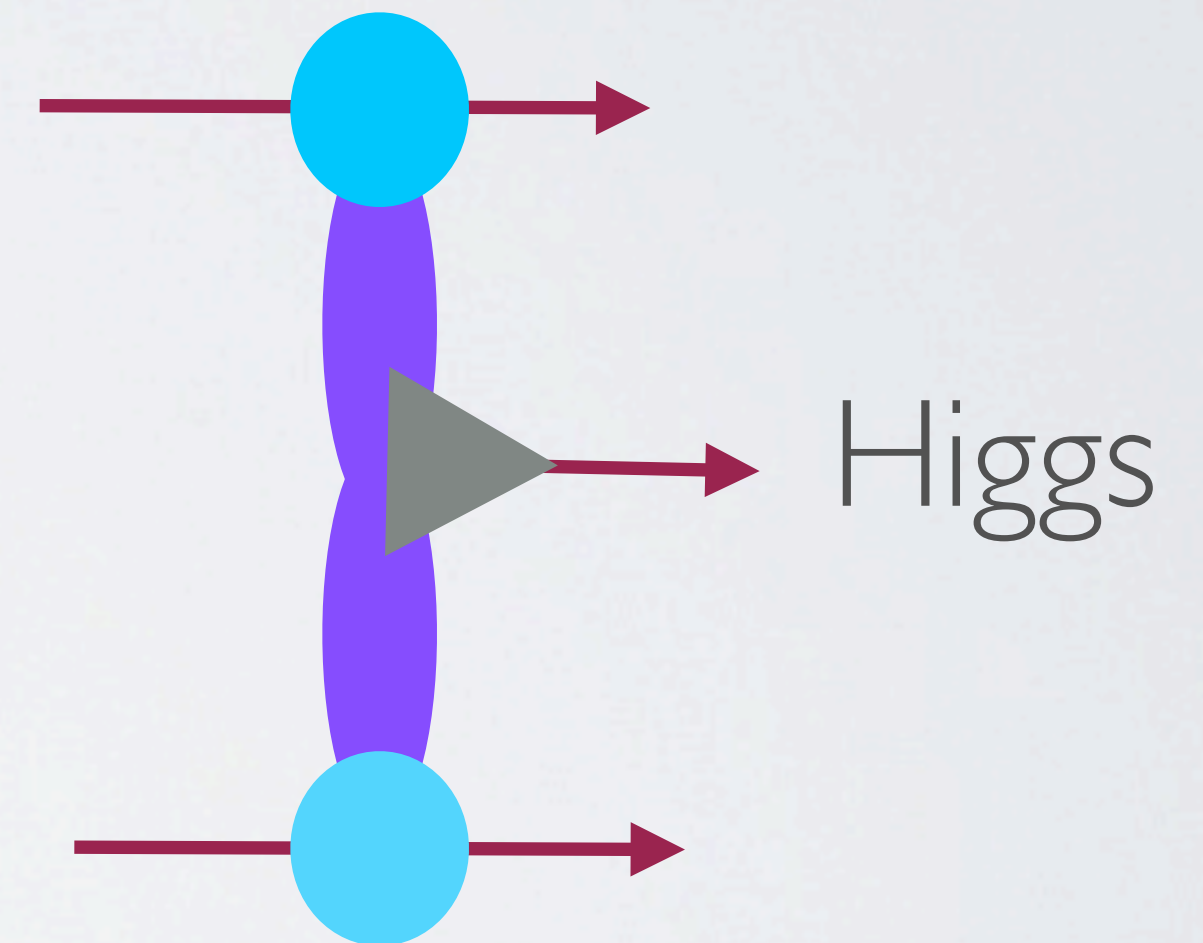
$$s \gg s_1, s_2 \gg t_1, t_2$$



- Transverse Correlation:

$$t_1 \simeq -q_{3,\perp}^2, \quad t_2 \simeq -q_{4,\perp}^2$$

$$\kappa = \frac{s_1 s_2}{s} \simeq m_H^2 + q_{H,\perp}^2$$



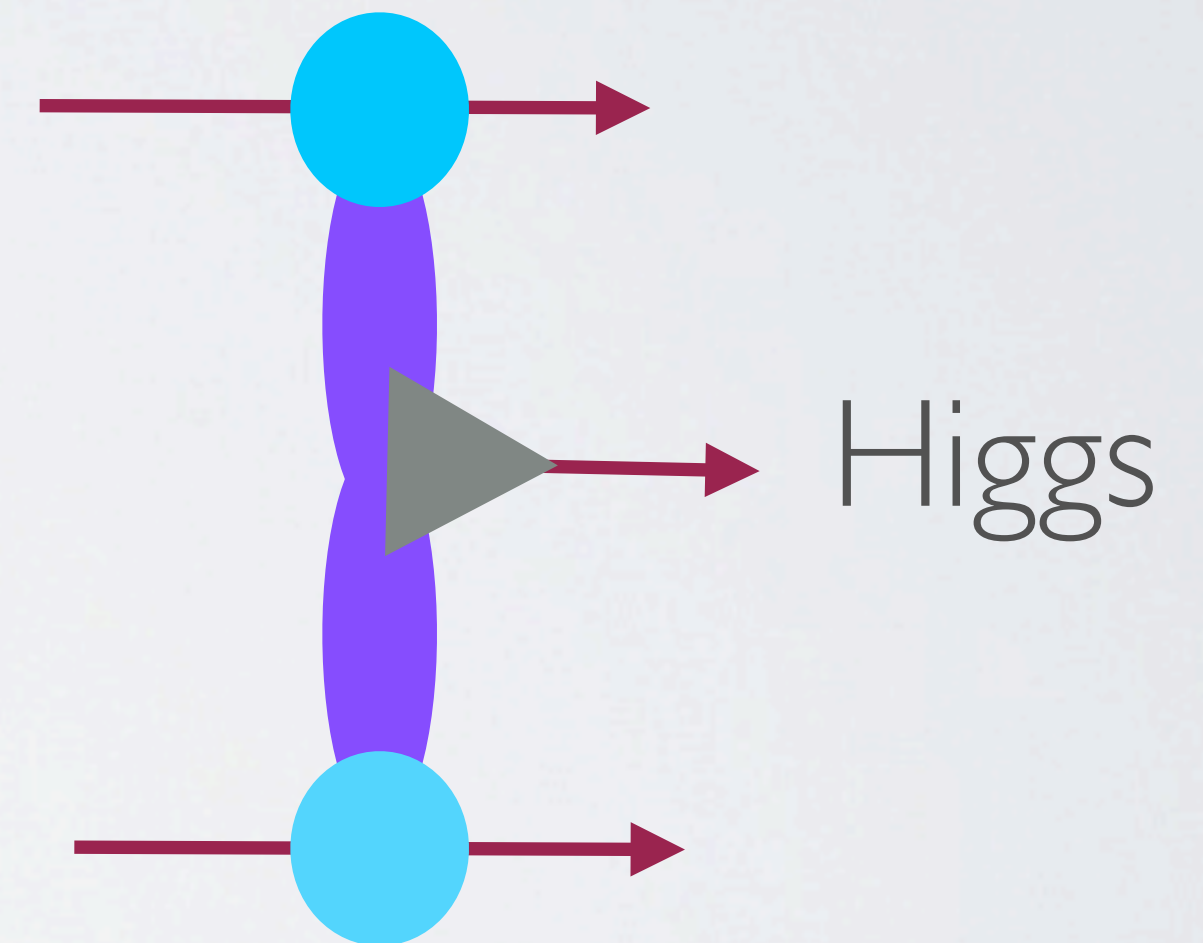
BASIC BUILDING BLOCK

- Elastic Vertex:
- Pomeron/Graviton Propagator:

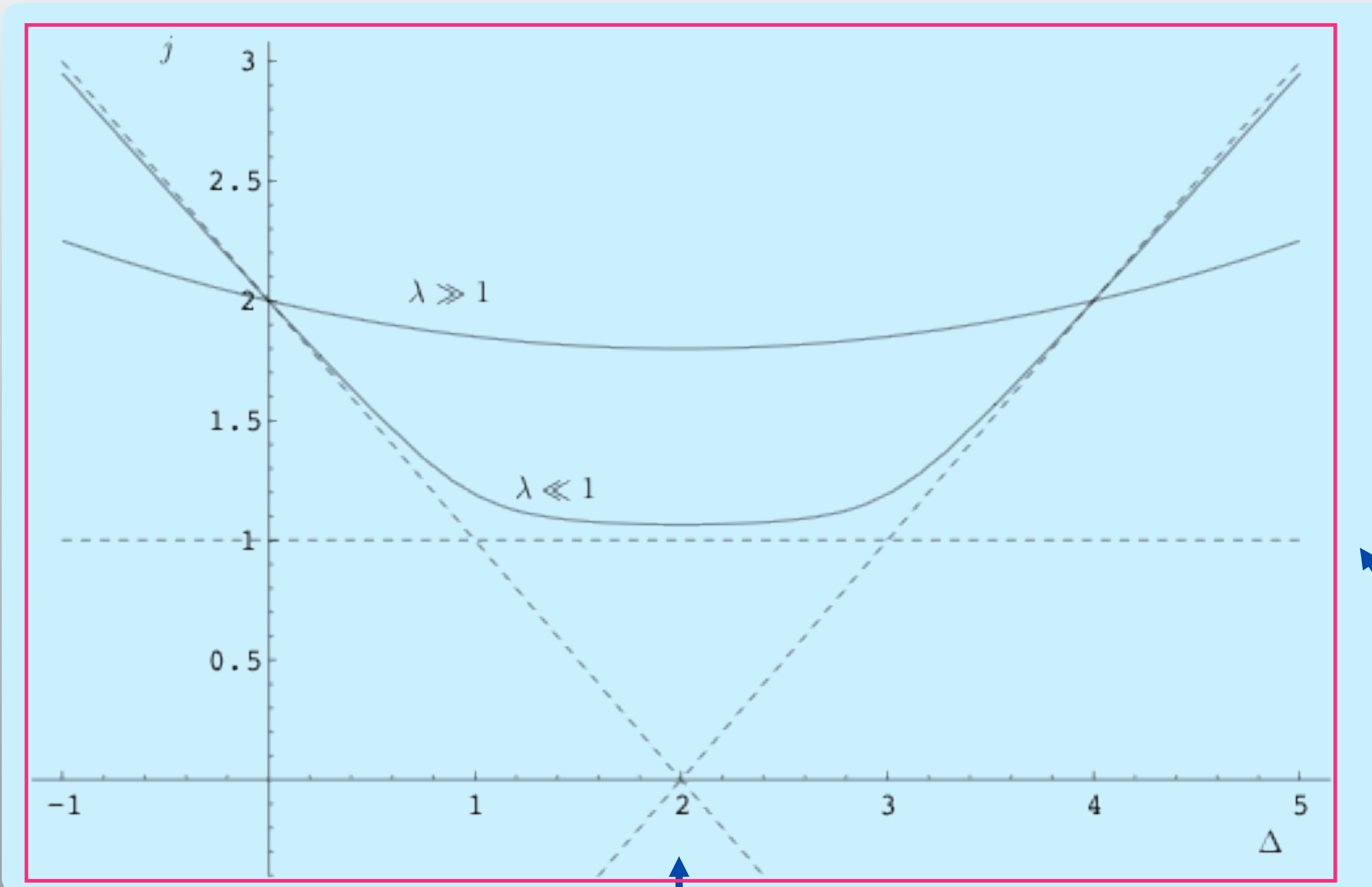
$$G_j(z, x^\perp, z', x'^\perp) = \frac{1}{4\pi z z'} \frac{e^{(2-\Delta(j))\xi}}{\sinh \xi},$$

$$\Delta(j) = 2 + \sqrt{2} \lambda^{1/4} \sqrt{(j-j_0)}$$

$$\mathcal{K}(s, b, z, z') = - \left(\frac{(z z')^2}{R^4} \right) \int \frac{dj}{2\pi i} \left(\frac{1 + e^{-i\pi j}}{\sin \pi j} \right) \widehat{s}^j G_j(z, x^\perp, z', x'^\perp; j)$$



$\mathcal{N} = 4$ SYM LEADING TWIST DELTA(J) VS J



$\lambda = 0$ DGLAP
(DIS moments)

$$\text{Tr}[F_{+\mu} D_+^{j-2} F_+^\mu]$$

$(0,2) T_{\mu\nu} \quad \gamma = 0$

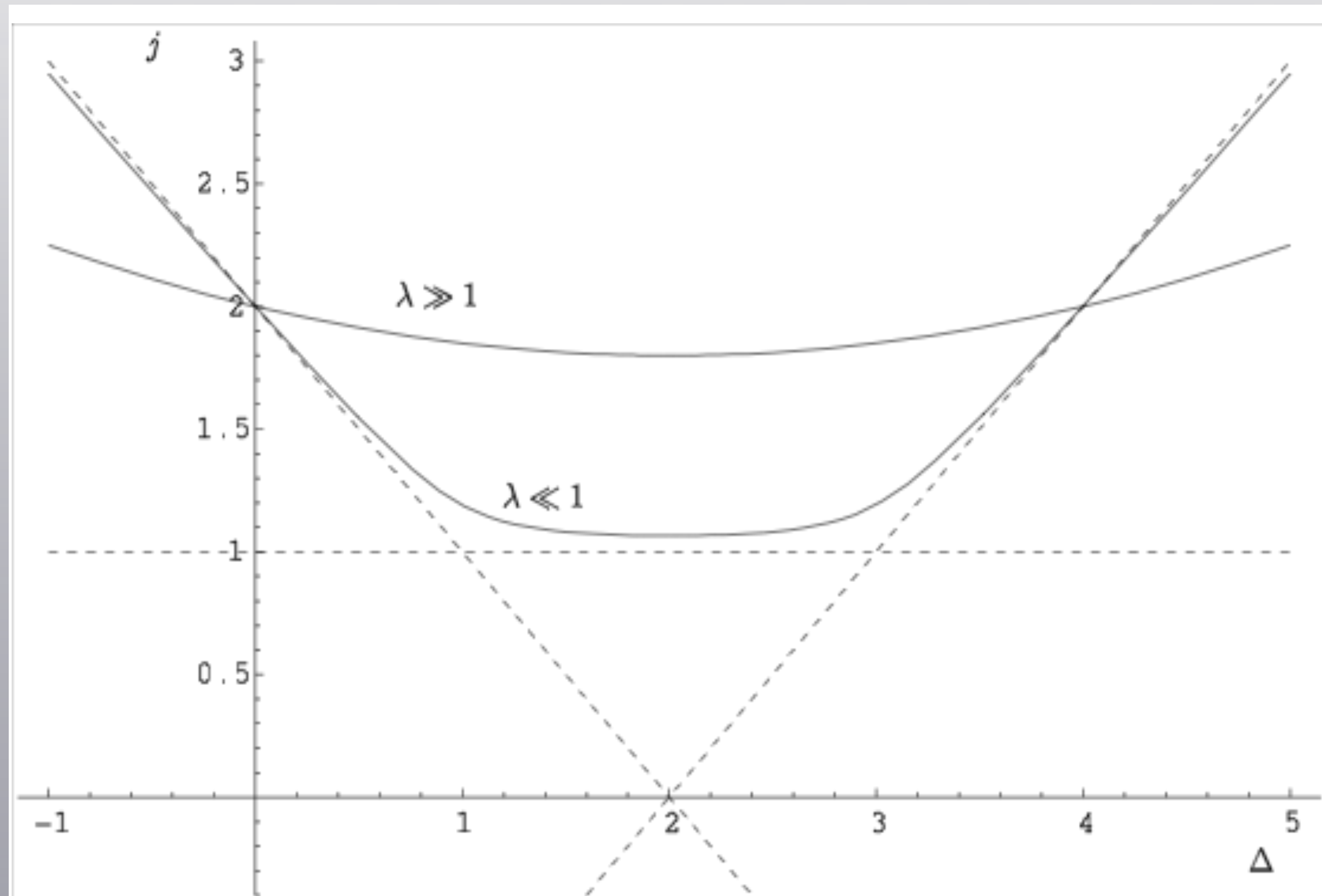
$\lambda = 0$, BFKL

$$\lambda = g^2 N = 0$$

$j = j_0 @ \text{min } \Delta$

MOMENTS AND ANOMALOUS DIMENSION

$$M_n(Q^2) = \int_0^1 dx x^{n-2} F_2(x, Q^2) \rightarrow Q^{4-\Delta_n}$$



$$\gamma_2 = 0$$

$$\Delta(j) = 2 + \sqrt{2} \sqrt{\sqrt{g^2 N_c} (j - j_0)}$$

$$\gamma_n = 2 \sqrt{1 + \sqrt{g^2 N} (n - 2)/2} - n$$

Simultaneous compatible large Q^2 and small x evolutions!

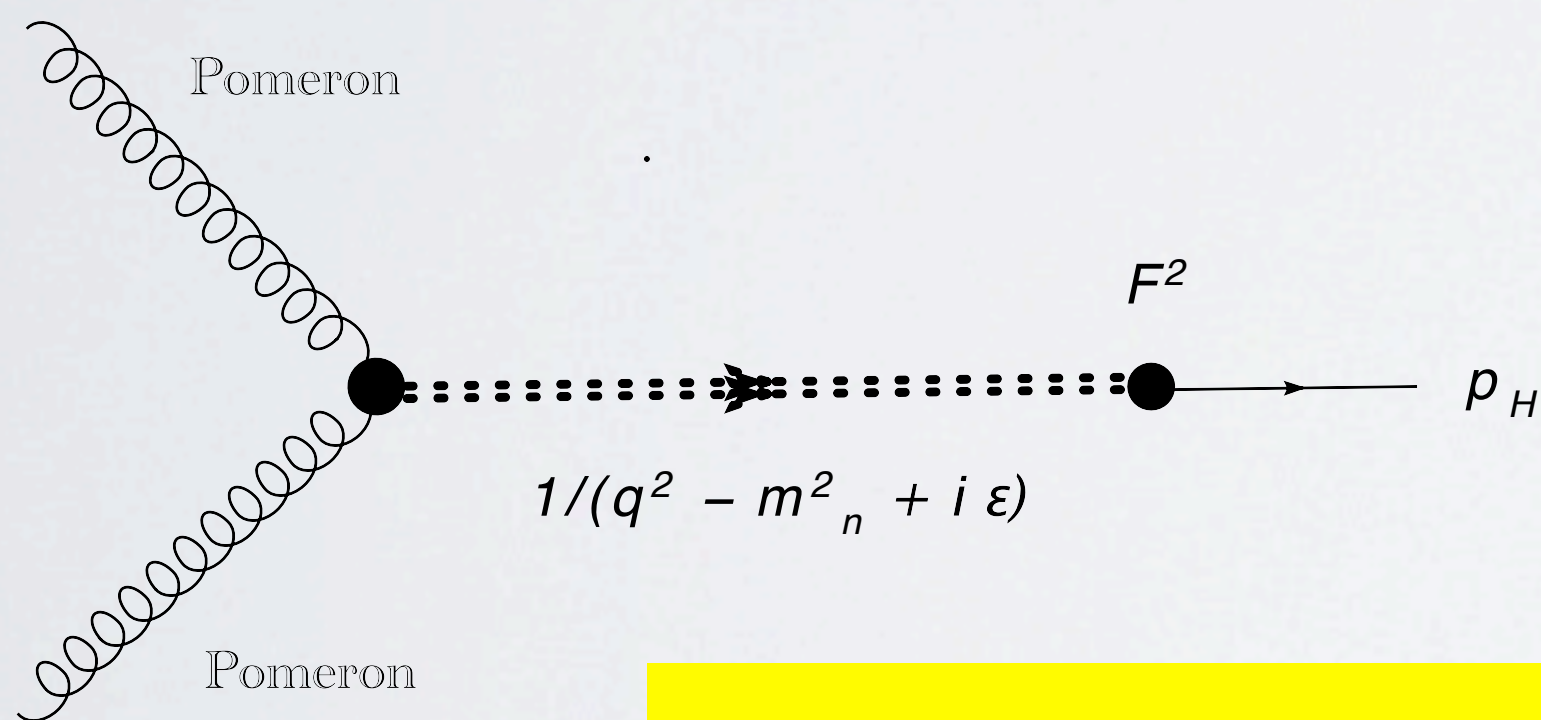
Energy-Momentum Conservation built-in automatically.

NEW FEATURE: POMERON-POMERON FUSION VERTEX

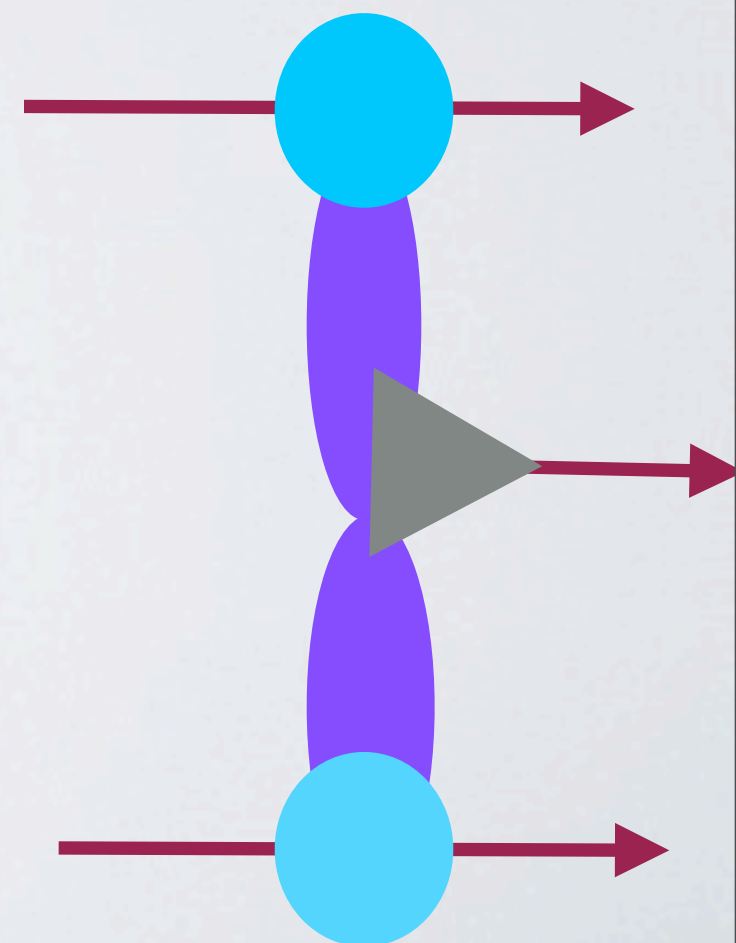
• Higgs coupled to Heavy quarks: $\mathcal{L} = -\frac{g}{2M_W} m_t \bar{t}(x)t(x)\phi_H(x)$.

• Integrating out heavy quarks leads to coupling to F^2 .

$$\mathcal{L} = \frac{\alpha_s g}{24\pi M_W} F_{\mu\nu}^a F^{a\mu\nu} \phi_H = L(m_H^2) F_{\mu\nu}^a F^{a\mu\nu} \phi_H$$



$$\mathcal{V}_H = V_{PP\phi} * K(m_H^2, z) * L(m_H^2)$$



CENTRAL VERTEX AND SCALAR INVARIANCE

- In scale invariance theory with exact AdS5-background, graviton-graviton-dilaton vertex vanishes identically
- To have non-vanishing expectation values for $\langle F^2 \rangle, \langle F^2 T_{\mu\nu} \rangle, \langle F^2 T_{\mu\nu} T_{\mu'\nu'} \rangle$, scale invariance must be broken.

- With Confinement deformation, will have non-vanishing

$$\langle F^2 \rangle, \langle F^2 T_{\mu\nu} \rangle, \langle F^2 T_{\mu\nu} T_{\mu'\nu'} \rangle \neq 0,$$

- Will also have non-vanishing graviton-graviton-dilaton vertex: $V_{PP\phi} \neq 0$

MODEL FOR CONFINEMENT DILATON GRAVITY

$$S = M_P^2 \int d^5x \sqrt{g} \left(-\mathcal{R} - V(\phi) + \frac{1}{2} G^{MN} \partial_M \phi \partial_N \phi - \lambda(\phi) T(z) \right)$$

$$G_{mn} = g_{mn} + h_{mn}$$

$$\phi = \phi_{cl} + \varphi$$

• Constant Background:

$$\phi_{cl} = \text{constant}$$

$$V_{PP\phi} = 0$$

• Non-trivial background:

$$\phi_{cl} \neq \text{constant}$$

$$V_{PP\phi} \neq 0$$

$$S_{int} = \frac{M_P^2}{4} \int dz d^4x \sqrt{-g} h^{nm} h_{mn} [V'(\phi_{cl})\varphi - g^{zz} \partial_z \phi_{cl}(z) \partial_z \varphi] .$$

• Asymptotic Freedom -- “coupling” running with “scale” in UV:

PHENOMENOLOGICAL ESTIMATES FOR DIFFRACTIVE HIGGS PRODUCTION

- Normalizing Pomeron-Pomeron-Higgs coupling by trace-anomaly by going on-shell

$$\gamma_{GGH}(q^2 = 0) = \frac{2M_G^2}{3vb} = 2^{1/4} G_F^{1/2} \frac{2M_G^2}{27}$$

- Use Strong Coupling Pomeron/Graviton Kernel to continue back to scattering region where $t < 0$.
- Use phenomenological parametrization for diffraction peaks.
- Estimate double-Pomeron Higgs production:

$$\frac{d\sigma}{dy_H} \simeq (1/\pi) \times m_1^{-4} \times |\gamma_{GGH}(0)|^2 \times \frac{\sigma(s)}{\sigma(m_H^2)} \times R_{el}^2(m_H \sqrt{s})$$

III. Beyond Pomeron

- Sum over all Pomeron graph (string perturbative, $1/N^2$)
- Eikonal summation in AdS_3
- Constraints from Conformal Invariance, Unitarity, Analyticity, Confinement, Universality, etc.
- Froissart Bound?
- Other “non-perturbative”

Review of High Energy Scattering in String Theory

DIS in AdS

It is convenient to introduce the “eikonal” $\chi(s, b, z, z')$ since we will show shortly that it is possible to express the amplitude in an eikonal sum as

$$A(s, t) = 2is \int d^2b e^{i\vec{q}\cdot\vec{b}} \int dz dz' P_{13}(z) P_{24}(z') \{1 - e^{i\chi(s, b, z, z')}\},$$

To first order in the eikonal,

$$A_4(s, t) \simeq \int d^2b e^{-i\mathbf{b}\mathbf{q}_\perp} \int dz dz' P_{13}(z) P_{24}(z') (2s\chi(s, b, z, z')),$$

where

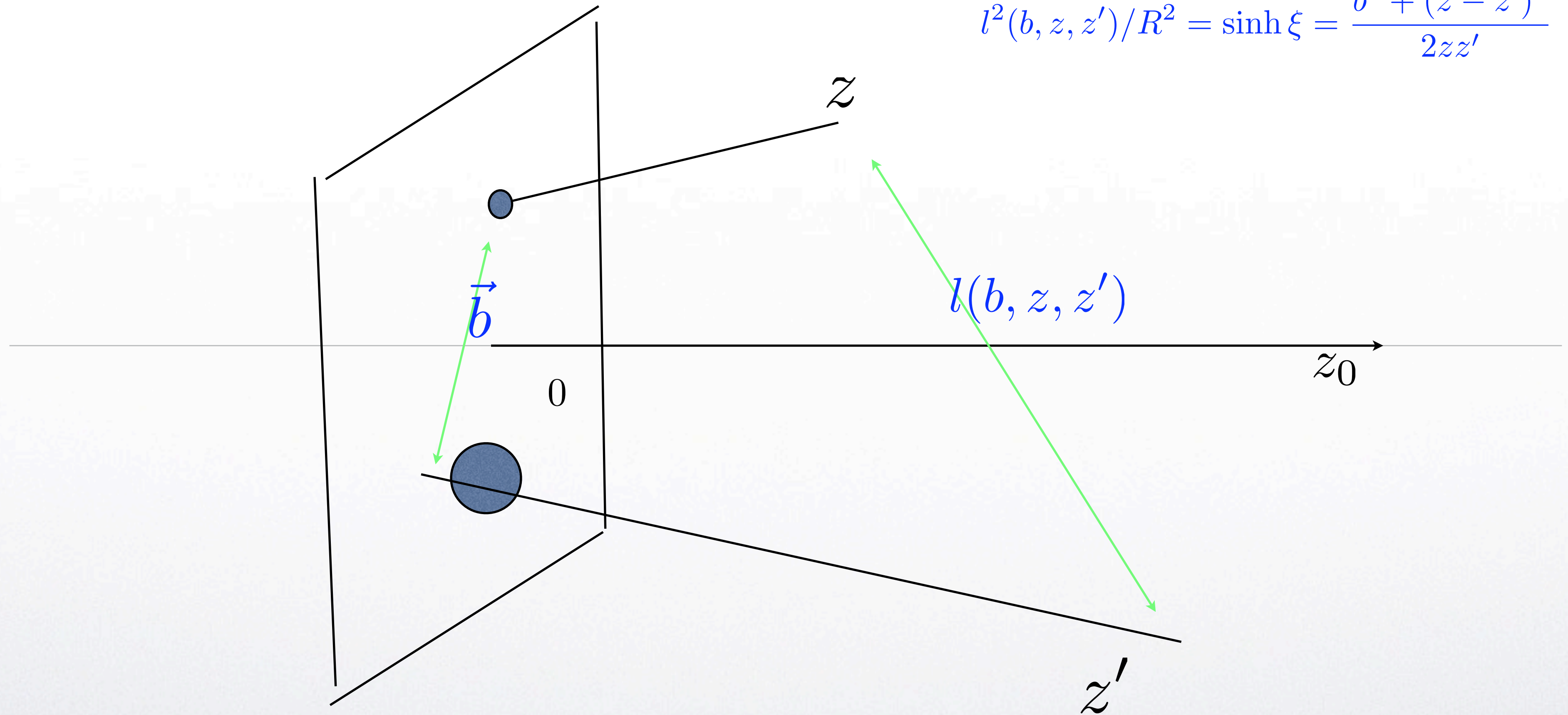
$$\chi(s, b, z, z') = \frac{g_0^2 R^4}{2(zz')^2 s} \mathcal{K}(s, b, z, z')$$

\mathcal{K} is the BPST Pomeron kernel.

Geometry of Transverse AdS-3

chordal distance: $l(b, z, z')$

$$l^2(b, z, z')/R^2 = \sinh^2 \xi = \frac{b^2 + (z - z')^2}{2zz'}$$



CONCLUDING COMMENTS

- Remarkable strong/weak correspondence
- Natural synthesis of near conformal & breaking:
 - IR (confinement) and UV (asymptotic freedom)
- By Non-pert. QCD-2013:
 - xsection for Double Diffractive Higgs observed,
 - Pomeron-Pomeron Higgs fusion better understood.
 - Constrained by
 - total cross section
 - DIS at small x , etc.

