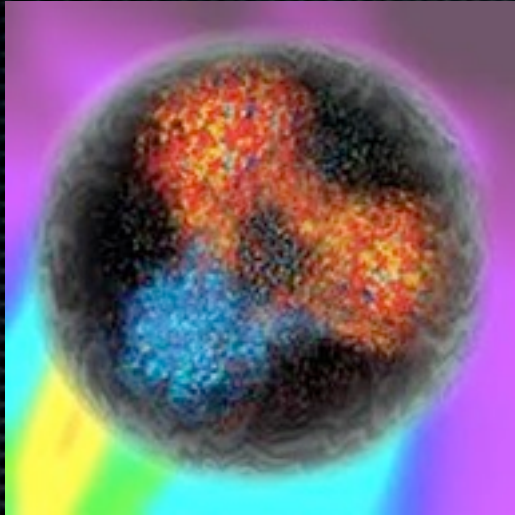


From Running Gluon Mass to Chiral Symmetry Breaking



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11th workshop non-perturbative QCD - Paris 2011

Outline

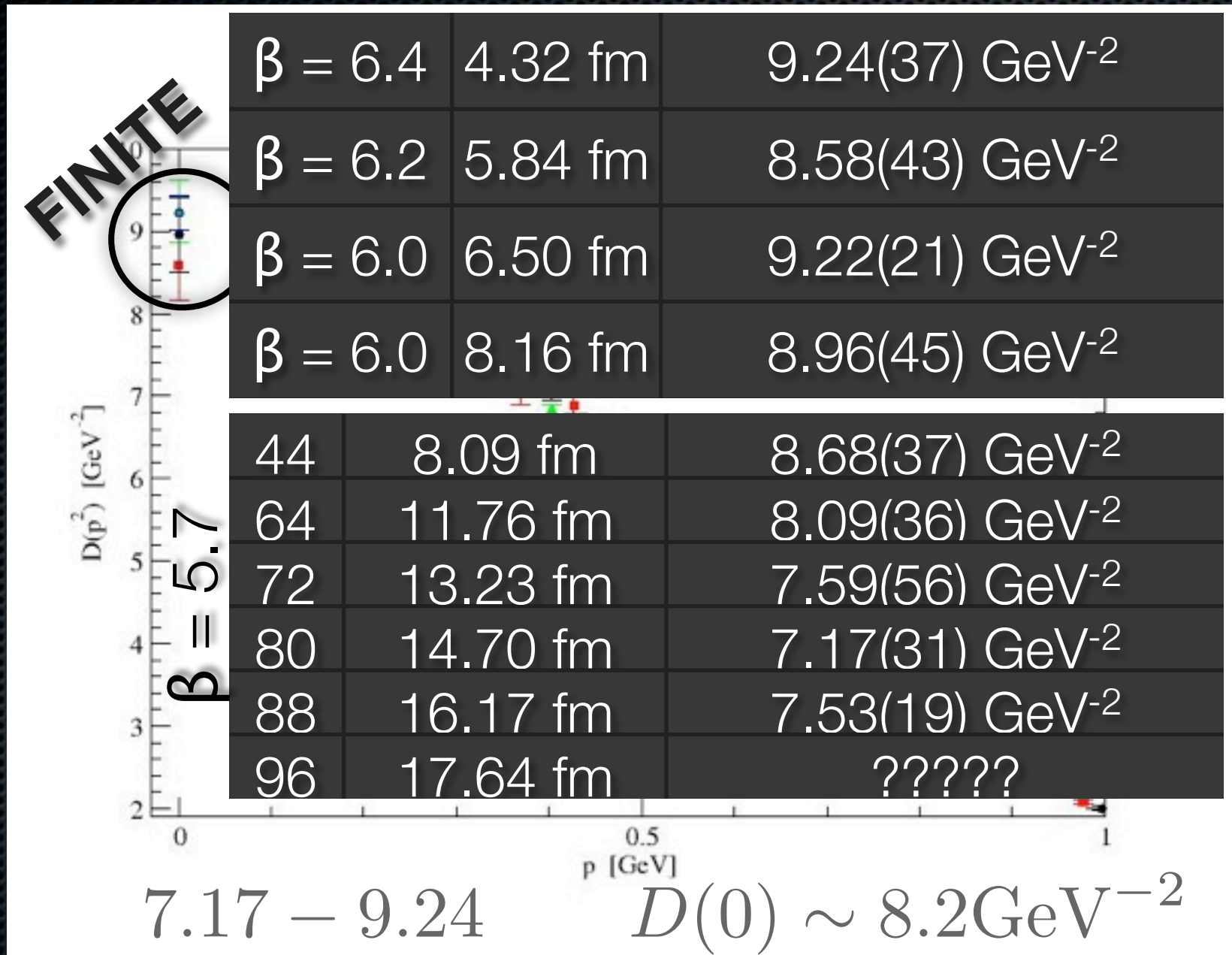
- ✦ **Lattice Gluon Propagator**
- ✦ **What are the implications of having a gluon mass**

Lattice Gluon Propagator

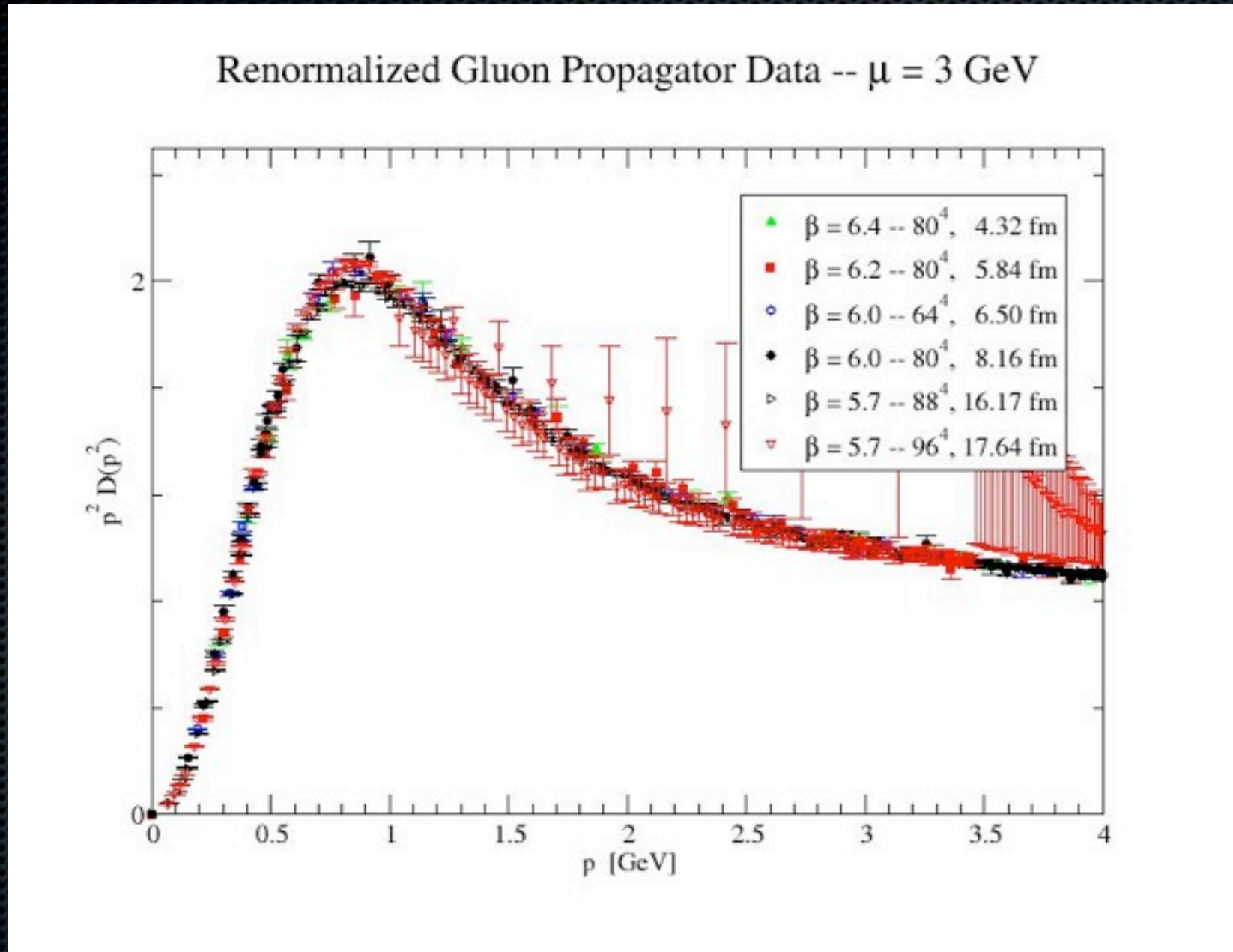
- ✦ Landau gauge
- ✦ Finite volume with periodic boundary conditions
- ✦ be aware of possible finite volume and lattice spacing effects
- ✦ Large discussion about $D(0)$

$$D(p^2) = \delta^{ab} \left(\delta_{\mu\nu} - \frac{p_\mu p_\nu}{p^2} \right) D(p^2)$$

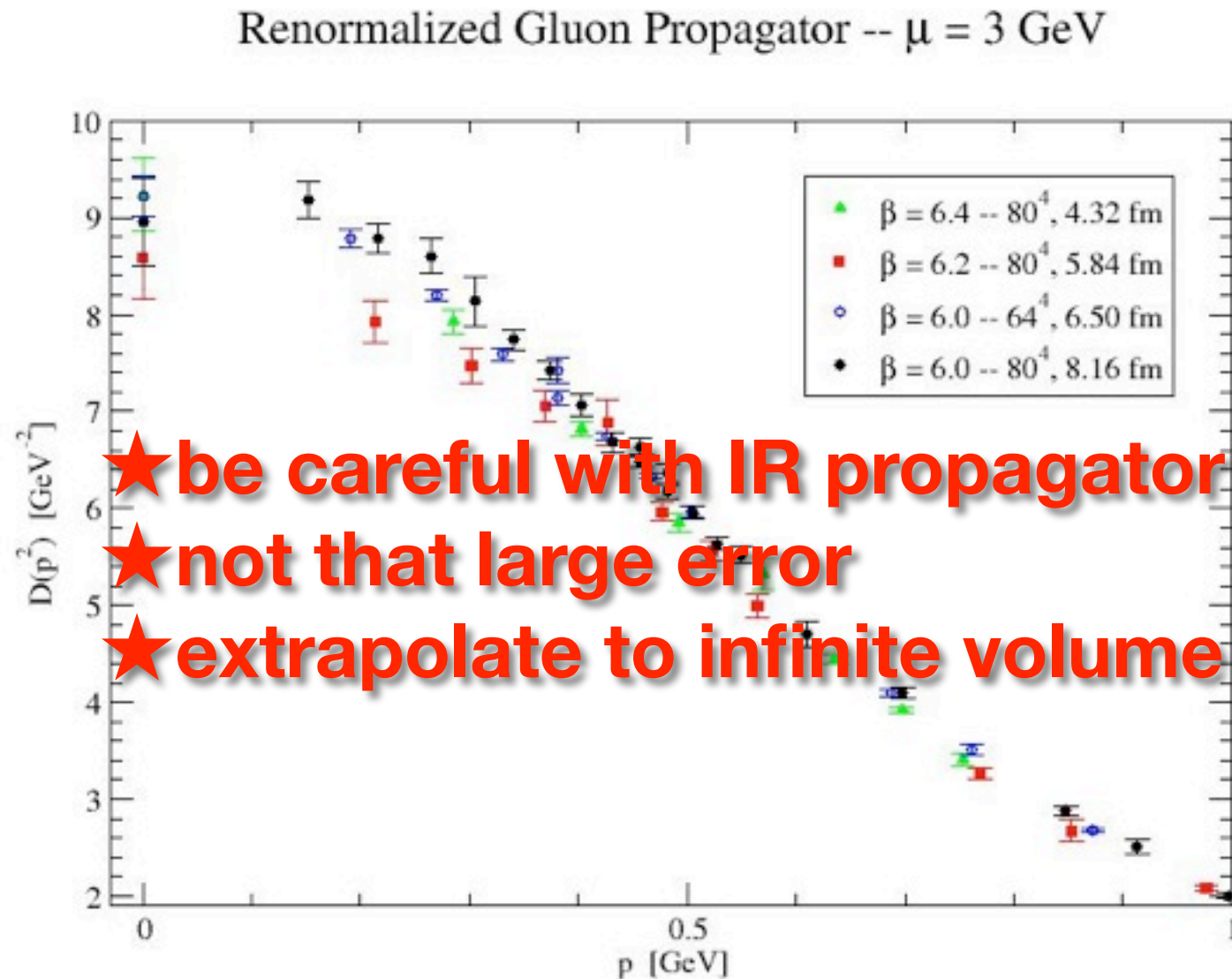
How well do we know $D(p^2)$



How well do we know $D(p^2)$



How well do we know $D(p^2)$



How well do we know $D(p^2)$

$$D(0) \sim 8.2\text{GeV}^{-2}$$

→ finite mass

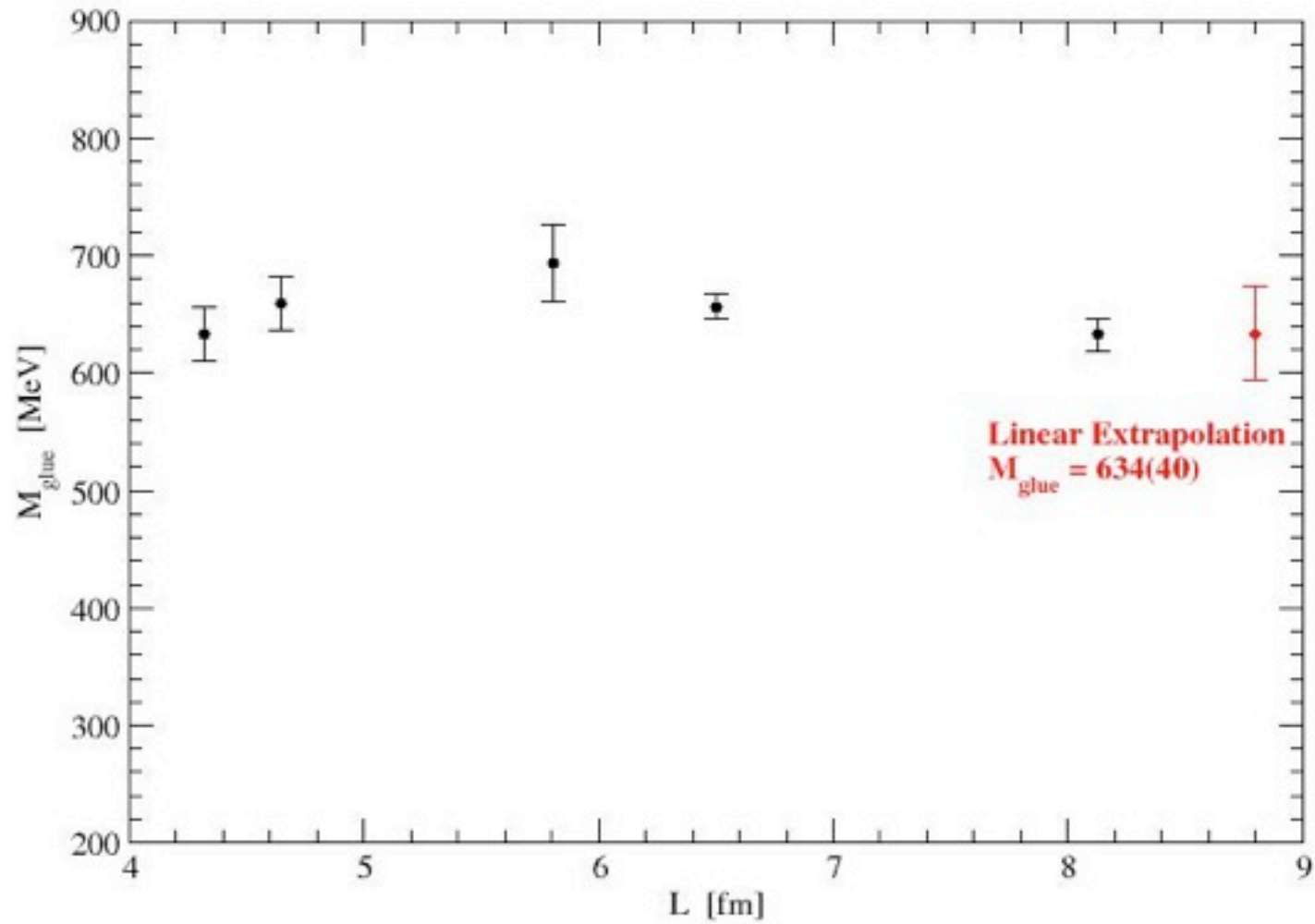
$$D(p^2) = \frac{Z}{p^2 + M^2}$$

$$D(p^2) = \frac{Z(p^2)}{p^2 + M^2(p^2)}$$

don't forget the
log here

how can one measure
the running mass?

Low Energy Propagator



Running Gluon Mass

sliding window analysis:

$$Z(p^2) = Z_0 [A + \ln(p^2 + m_0^2)]^{-\gamma} \quad M^2(p^2) = p^2 \ln p^2$$

model:
$$Z(p^2) = \frac{z_0}{\left[\ln \frac{p^2 + r m_0^2}{\Lambda} \right]^\gamma} \quad M^2(p^2) = \frac{m_0^2}{p^2 + m_0^2}$$

$$z_0 = 1.189(20)$$

$$\Lambda = 1.842(39) \text{ GeV} \quad m_0 = 671(9) \text{ MeV}$$

$$r = 7.49(59)$$

$$p_{max} = 4.2 \text{ GeV}$$

O. Oliveira, P. Bicudo, J. Phys. **G38**, 045003 (2011)

Refined Gribov-Zwanziger Action

- ▶ Gribov copies
- ▶ Restriction of the functional integration space

$$D(p^2) = \frac{p^2 + M^2}{p^4 + (M^2 + m^2)p^2 + 2g^2 N \gamma^4 + M^2 m^2}$$

mass

Describes well the Lattice Data up to 1.5 GeV

$$M^2(p^2) = \frac{m^2 p^2 + 2g^2 N \gamma^4 + M^2 m^2}{p^2 + M^2}$$

meter

D. Dudal, O. Oliveira, N. Vandersickel, PRD81, 074505 (2010)

Fitting the Lattice Data

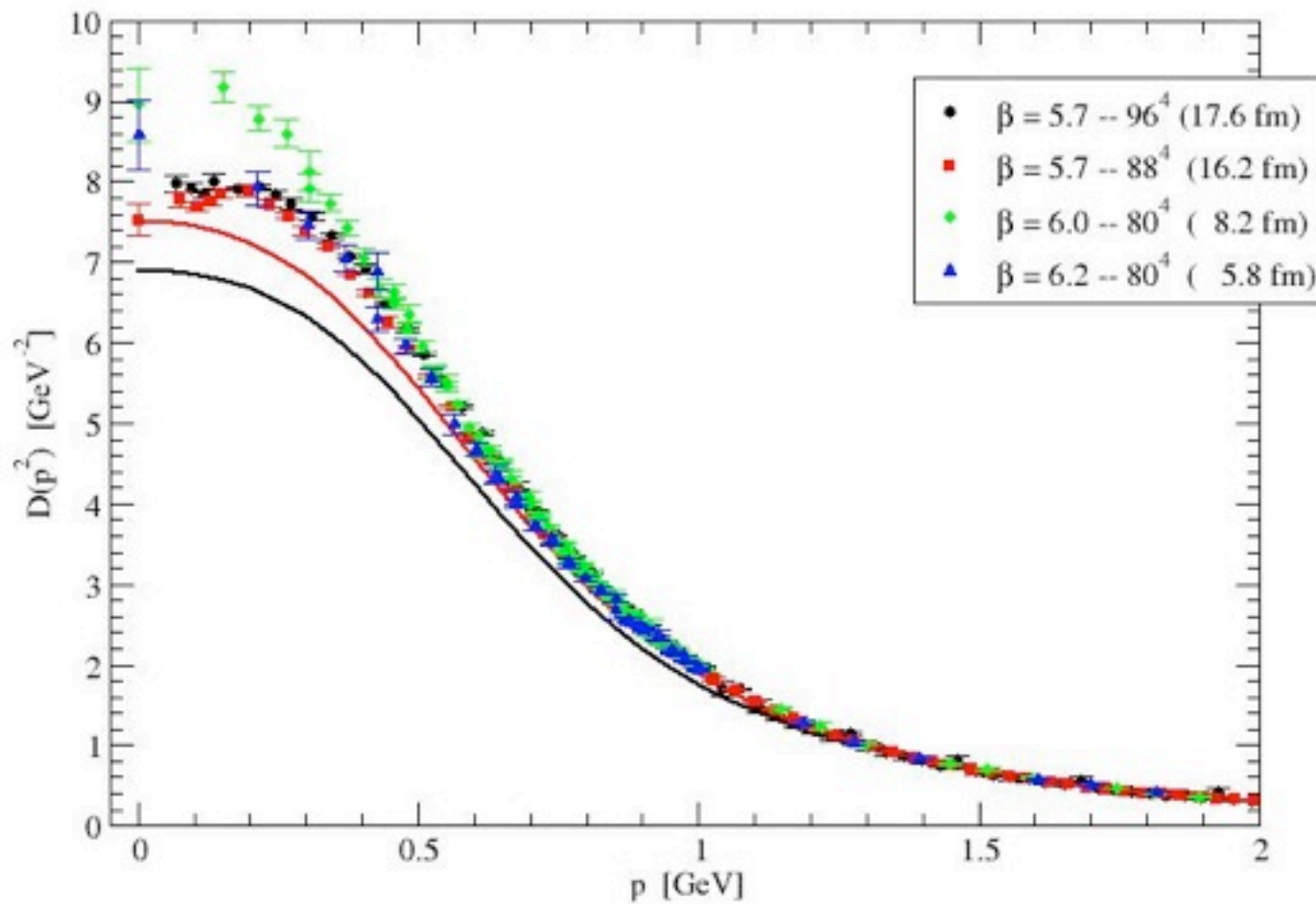
β	La	χ^2/dof	M^2	$M^2 + m^2$	λ^4	$\lambda^4 p_{\text{max}}$
6.0	4.88	1.6	2.81(9)	0.62(3)	0.284(7)	1.2
6.0	6.50	1.80(22)	2.66(6)	0.54(2)	0.288(5)	0.95
	8.13	1.1	2.41(5)	0.47(2)	0.261(4)	1.6
6.2	3.49	1.1	2.5(1)	0.48(4)	0.28(1)	1.4
6.2	4.65	2.10(24)	2.44(7)	0.36(11)	0.280(10)	1.5
	5.81	1.2	2.3(1)	0.42(4)	0.273(9)	1.6

$D(0) = 7.50(90)$

$D(0) = 6.90(85)$

Extrapolated Propagator

Renormalized Gluon Propagator - $\mu = 3 \text{ GeV}$



Gluon Condensate

$$\langle g^2 A^2 \rangle = -\zeta_0 m^2$$

$$\zeta_0 = \frac{9}{13} \frac{N^2 - 1}{N}$$

$$\langle g^2 A^2 \rangle|_3 \text{ GeV} = 2.71(41) \quad 3.21(48) \quad \text{GeV}$$

$$\langle g^2 A^2 \rangle|_{10} \text{ GeV} = 2.45(38) \quad 2.90(43) \quad \text{GeV}$$

$$\beta = 6.0$$

$$\beta = 6.2$$

$$\langle g^2 A^2 \rangle|_{10} \text{ GeV} = 5.1_{-1.1}^{+0.7} \text{ GeV}$$

glue

P. Boucaud et al

PRD**63**, 114003 (2001)

PRD**66**, 034504 (2002)

$$\langle g^2 A^2 \rangle|_{10} \text{ GeV} = 4.4 \pm 0.4 \text{ GeV}$$

quark

PRD**79**, 014508 (2009)

Why a Gluon Mass?

$$8 \otimes 8 = 1 \otimes \underline{8 \otimes 8} \otimes 10 \otimes \overline{10} \otimes 27$$

color octet

singlet = glueballs

lightest glueball - 1.7 GeV
to heavy for a low energy theory

$$\begin{array}{c} \cancel{f_{abc}} \\ d_{abc} \end{array} \left| F_{\mu\nu}^b F^{c\mu\nu} + \bar{q} t^a q \sim \phi^a \right.$$

O. Oliveira, W. de Paula, T. Frederico, arXiv:1105.4899

Effective Degrees of Freedom

Effective Degrees of Freedom

$$\phi^a = \frac{d_{abc} F^b F^c}{\Lambda^3} + \frac{\bar{q} t^a q}{\Lambda^2}$$
$$\sim \frac{\langle F^2 \rangle}{\Lambda^3} + \frac{\langle \bar{q} q \rangle}{\Lambda^2}$$

$$\alpha_s \langle F^2 \rangle = 0.04 \text{ GeV}^4 \quad \langle \bar{q} q \rangle = (-270 \text{ MeV})^3 \quad \Lambda = 0.3 \text{ GeV}$$

$$\frac{\text{glue}}{\text{quark}} \sim 7$$

Effective Lagrangian

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu}^a F^{a\mu\nu} + \sum_f \bar{q}_f \{i\gamma^\mu D_\mu - m\} q_f \\ + \frac{1}{2} (D^\mu \phi^a) (D_\mu \phi^a) - V_{oct}(\phi^a \phi^a) + \mathcal{L}_{GF} + \mathcal{L}_{gh} + \mathcal{L}_{new}$$

$$G_4 \sum_f \left[\bar{q} t^a q \right] \phi^a,$$

$$G_5 \sum_f \left[\bar{q} q \right] \phi^a \phi^a + F_1 \sum_f \left[\bar{q} t^a q \right] d_{abc} \phi^b \phi^c \\ + F_2 \sum_f \left[\bar{q} t^a \gamma^\mu q \right] D_\mu \phi^a + F_3 \sum_f \left[\bar{q} t^a \gamma^\mu D_\mu q \right] \phi^a$$

Mass Generation

$$\frac{1}{2} (D^\mu \phi^a) (D_\mu \phi^a)$$

$$\frac{1}{2} g^2 \phi^c (T^a T^b)_{cd} \phi^d A_\mu^a A^{b\mu}$$

$$\langle \phi^a \rangle = 0 \quad \text{and} \quad \langle \phi^a \phi^b \rangle = v^2 \delta^{ab}$$

$$m_g^2 = N g^2 v^2$$

$$G_5 \bar{q} q \phi^a \phi^a \quad - \frac{G_5}{g^2} \frac{N^2 - 1}{N} m_g^2$$

Chiral Symmetry Breaking

$$\langle \bar{q} q \rangle = -\frac{4NN_f}{16\pi^2} M m_g^2 \left\{ \left(\frac{M}{m_g} \right)^2 \ln \frac{M^2}{M^2 + \bar{\omega}^2} + \left(\frac{\bar{\omega}}{m_g} \right)^2 \right\}$$

No Gluon Mass

No quark Mass

No chiral condensate

does not necessarily hold in all conditions

▶ other degrees of freedom

▶ $\langle \phi^a \rangle = 0$ $\langle \phi^a \phi^b \rangle = v^2 \delta^{ab}$

Parameters Estimation

$$M = 330 \text{ MeV} \quad m_g = 634(40) \text{ MeV} \quad \langle \bar{q} q \rangle = (-270 \text{ MeV})^3$$

$$\langle \bar{q} q \rangle = -\frac{4NN_f}{16\pi^2} M m_g^2 \left\{ \left(\frac{M}{m_g} \right)^2 \ln \frac{M^2}{M^2 + \bar{\omega}^2} + \left(\frac{\bar{\omega}}{m_g} \right)^2 \right\}$$

$$\bar{\omega} = 879 \text{ MeV}$$

$$m_g = \sqrt{N} g v$$

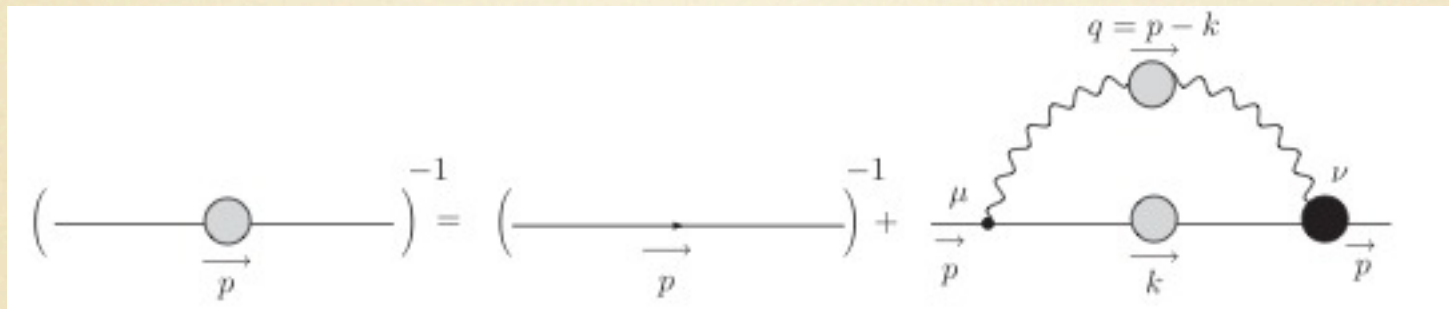
$$g v = 366 \text{ MeV}$$

$$M = -\frac{G_5}{g^2} \frac{N^2 - 1}{N} m_g^2$$

$$\frac{G_5}{g^2} = -0.31 \text{ GeV}^{-1}$$

Schwinger-Dyson Equations

A.C. Aguilar, J. Papavassiliou, PRD**83**, 014013 (2011)



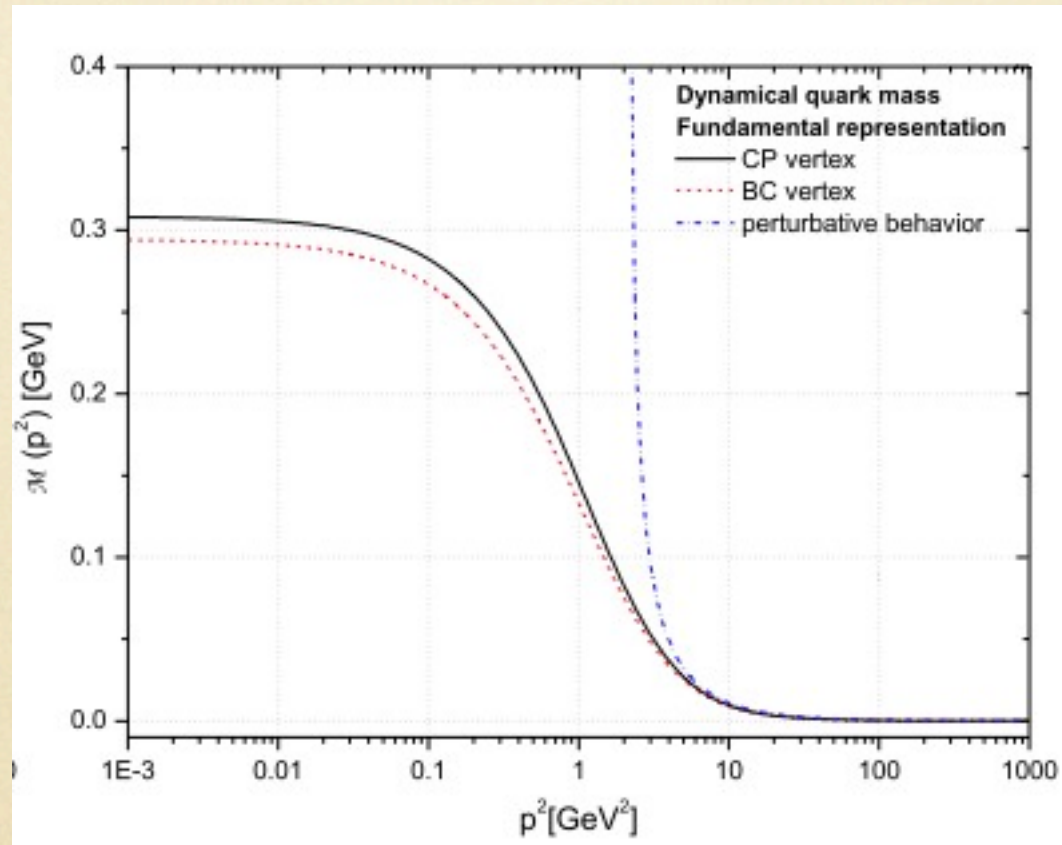
gluon propagator from the lattice $\beta = 5.7$ for $80^4 = (14.7 \text{ fm})^4$

two different ansatzes for quark-gluon vertex

ghost improved Ball-Chiu (BC) + Curtis-Pennington

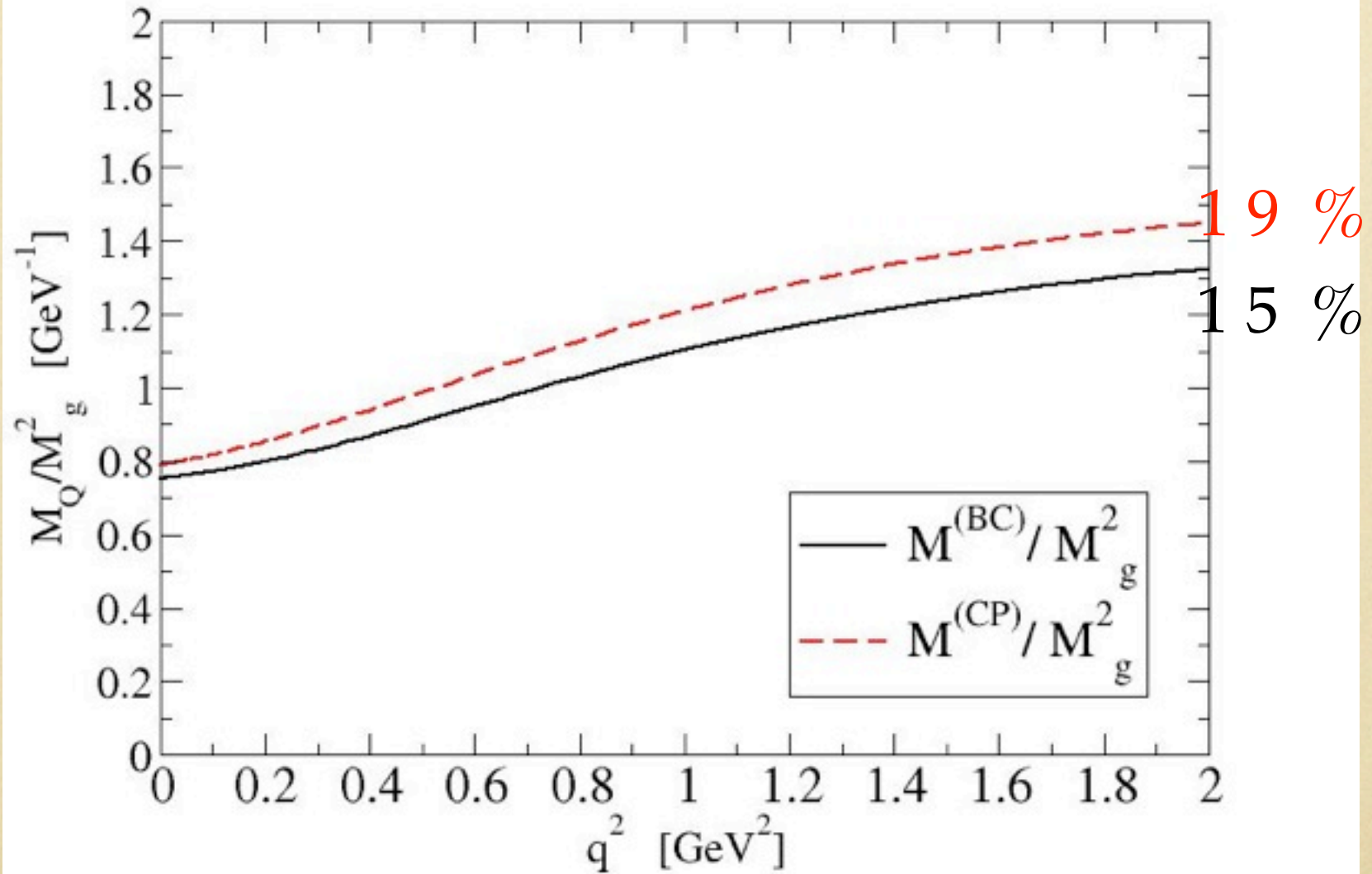
different renormalization point - $\mu = 4.3 \text{ GeV}$

Schwinger-Dyson Equations



dynamical generation of fermion mass = chiral symmetry breaking

Schwinger-Dyson Equations



Results and Conclusions

- ✓ Gluon Propagator is well known both in UV and IR
- ✓ Higher Statistics to get its details

- ✓ Gluon massive in IR
- ✓ Gluon mass implies chiral symmetry breaking
- ✓ It remains to investigate possible phenomenological implications of the low energy effective model

Thank You !!!