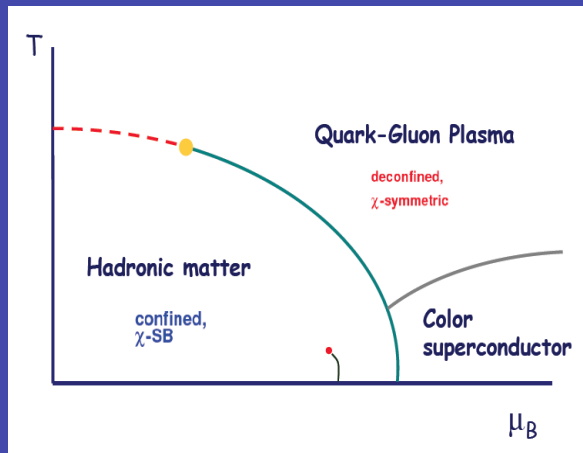
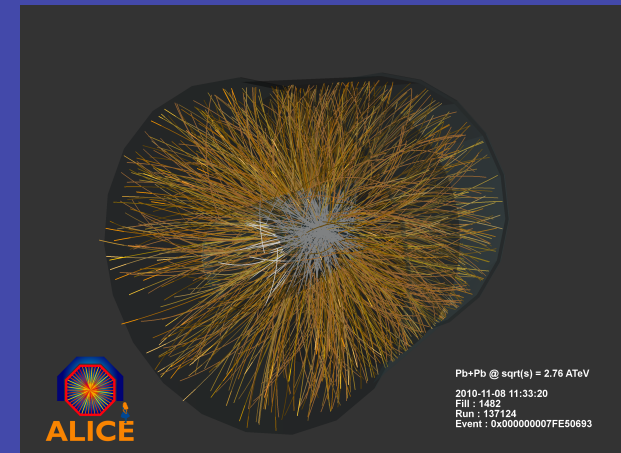


Extreme states of matter and ultra-relativistic heavy ion collisions

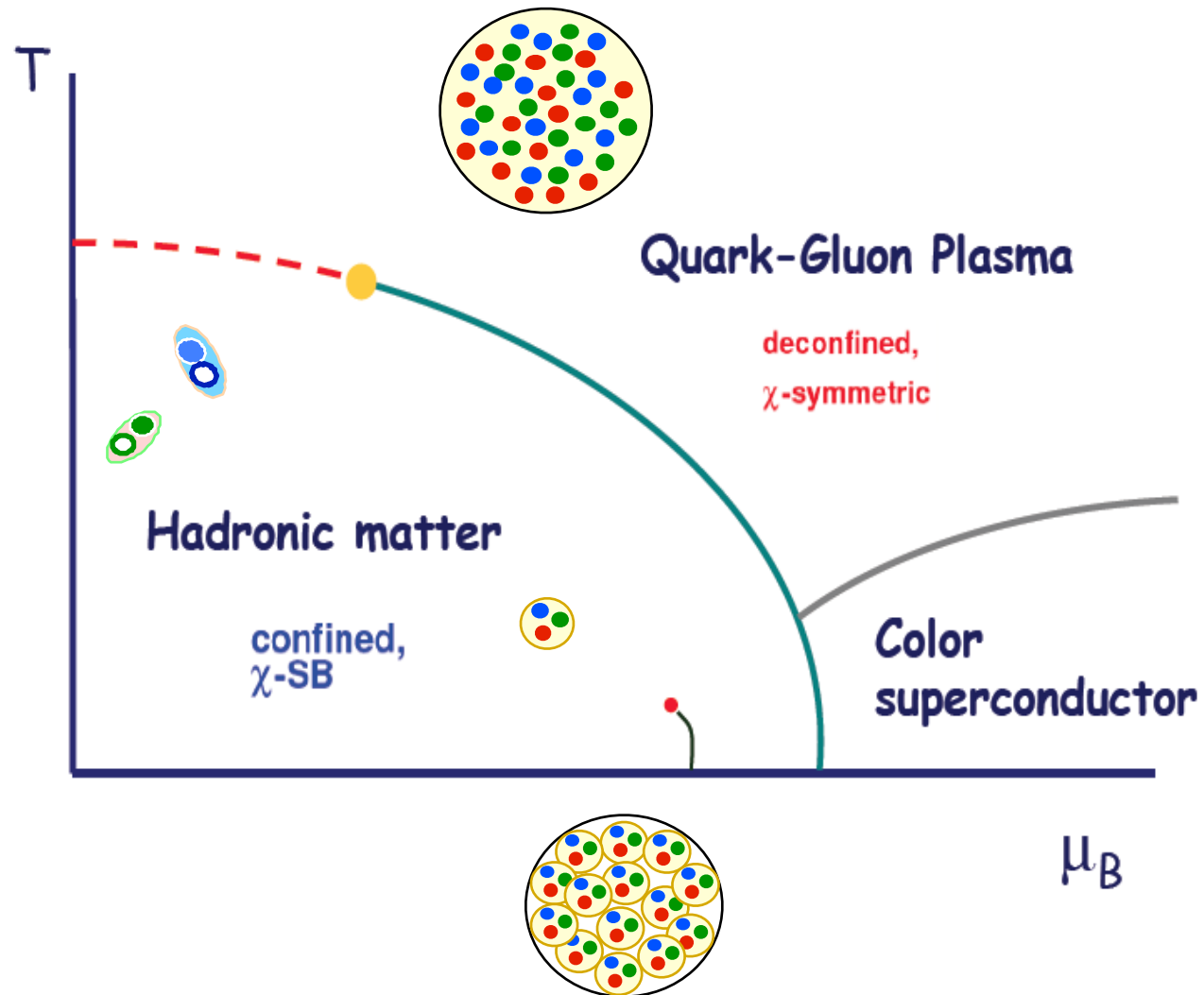


Rencontres du Vietnam
Qui Nhon
Vietnam
20/12/2011

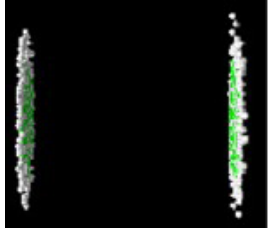


Jean-Paul Blaizot, IPhT- Saclay CEA & CNRS

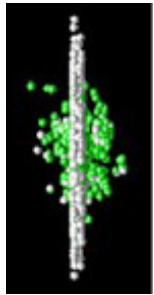
The QCD phase diagram



Colliding heavy nuclei

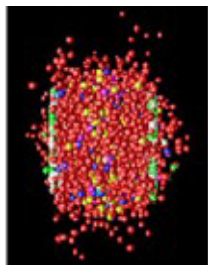


Initial conditions. Fluctuations (geometry, nucleus wave function and its parton content)

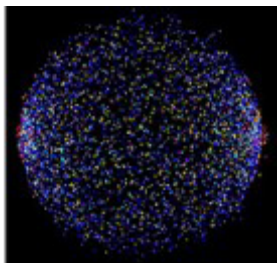


Particle (entropy) production. Involves mostly small x partons ($x = p_{\perp} / \sqrt{s} \sim 10^{-2} - 10^{-4}$ for $p_{\perp} \simeq 2\text{GeV}$)

One characteristic scale: saturation momentum Q_s



Thermalization. Quark-gluon plasma.
Hydrodynamical expansion



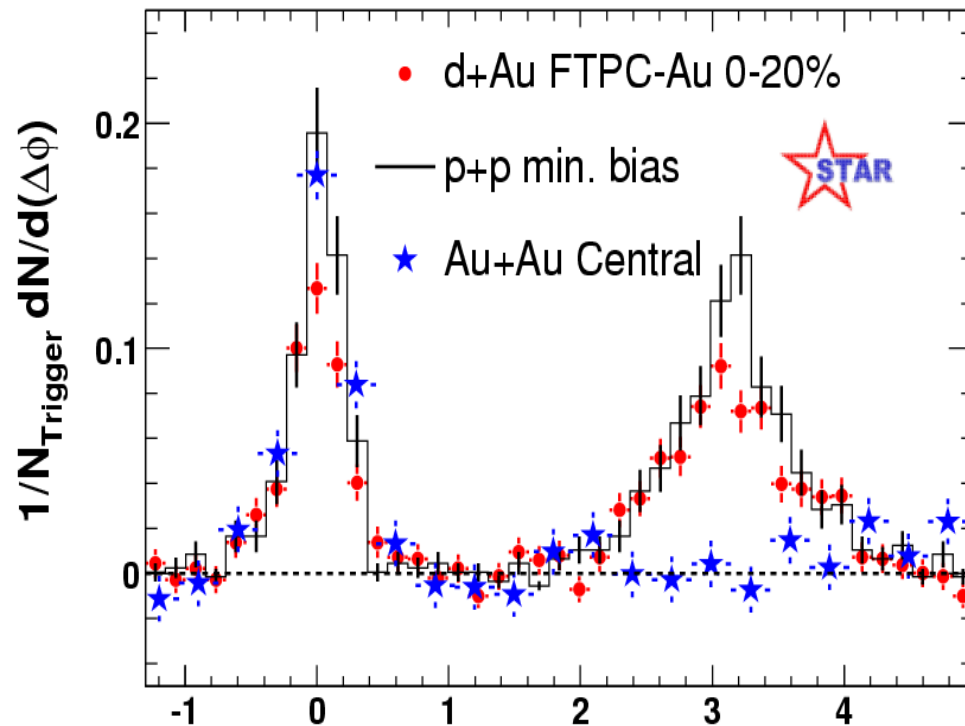
Hadronization. Hydrodynamic expansion continues till freeze-out. Apparent chemical equilibrium at freeze-out

Main surprises from RHIC

(confirmed by LHC)

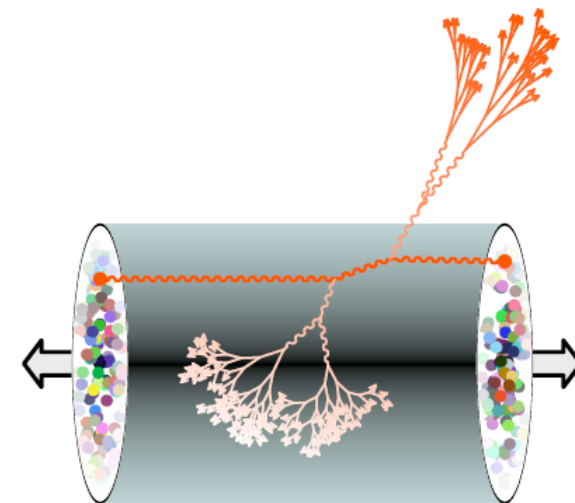
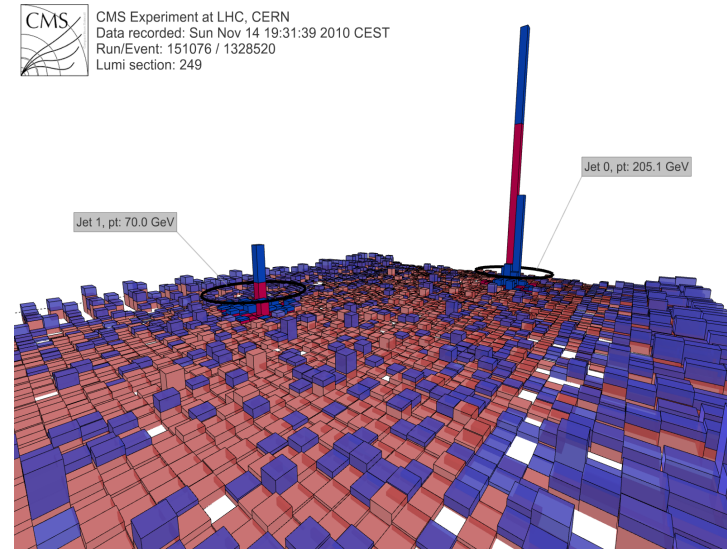
concern matter before freeze-out

Matter is opaque to the propagation of jets

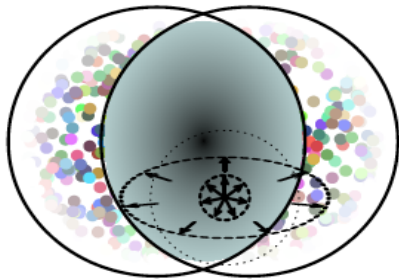
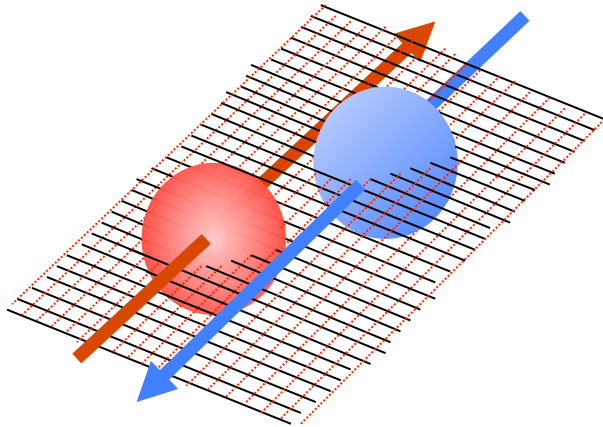


(STAR: Phys.Rev.Lett.91:072304,2003)

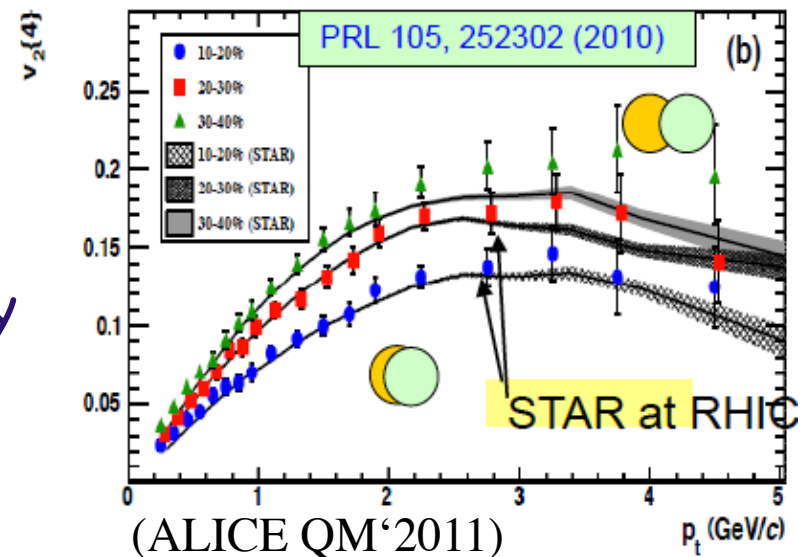
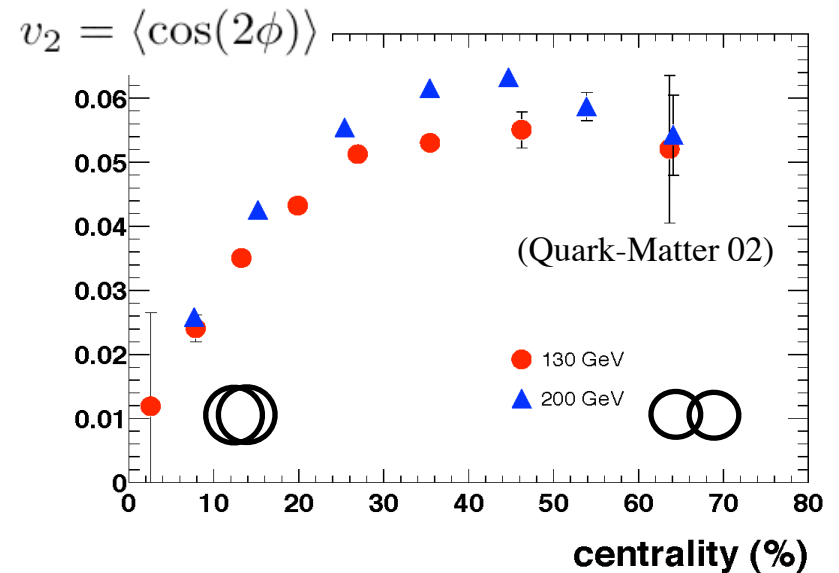
The jet produced near the surface 'escapes' normally. Its partner is absorbed in the produced plasma.



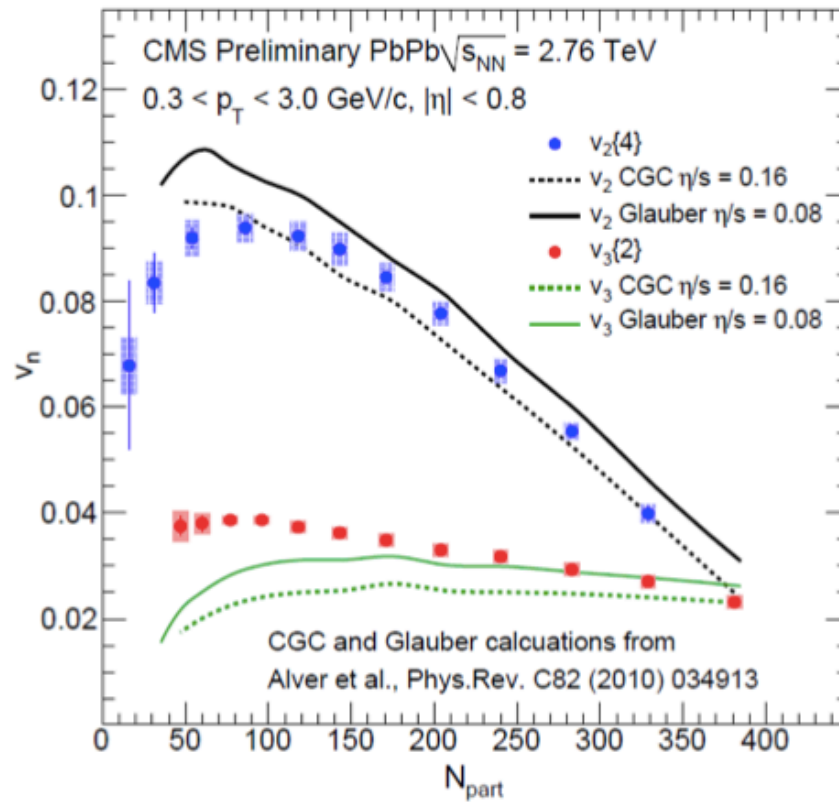
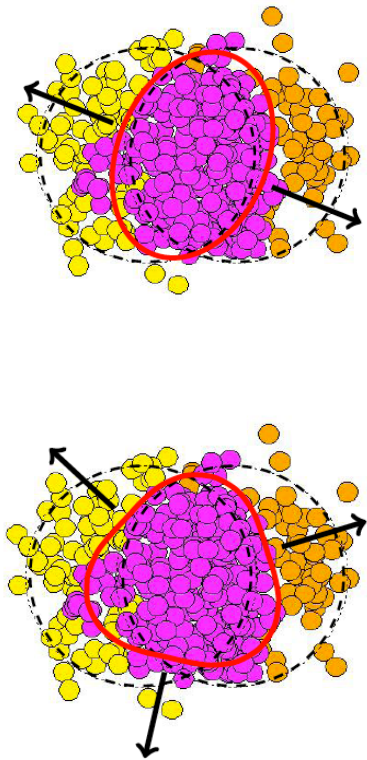
Produced matter flows like a fluid



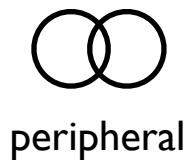
The flow takes place preferentially in the reaction plane rather than away from it.



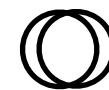
The flow is sensitive to initial nuclear density fluctuations



J. Velkowska,
QM2011
and
Luzum,
arXiv:1011.5173

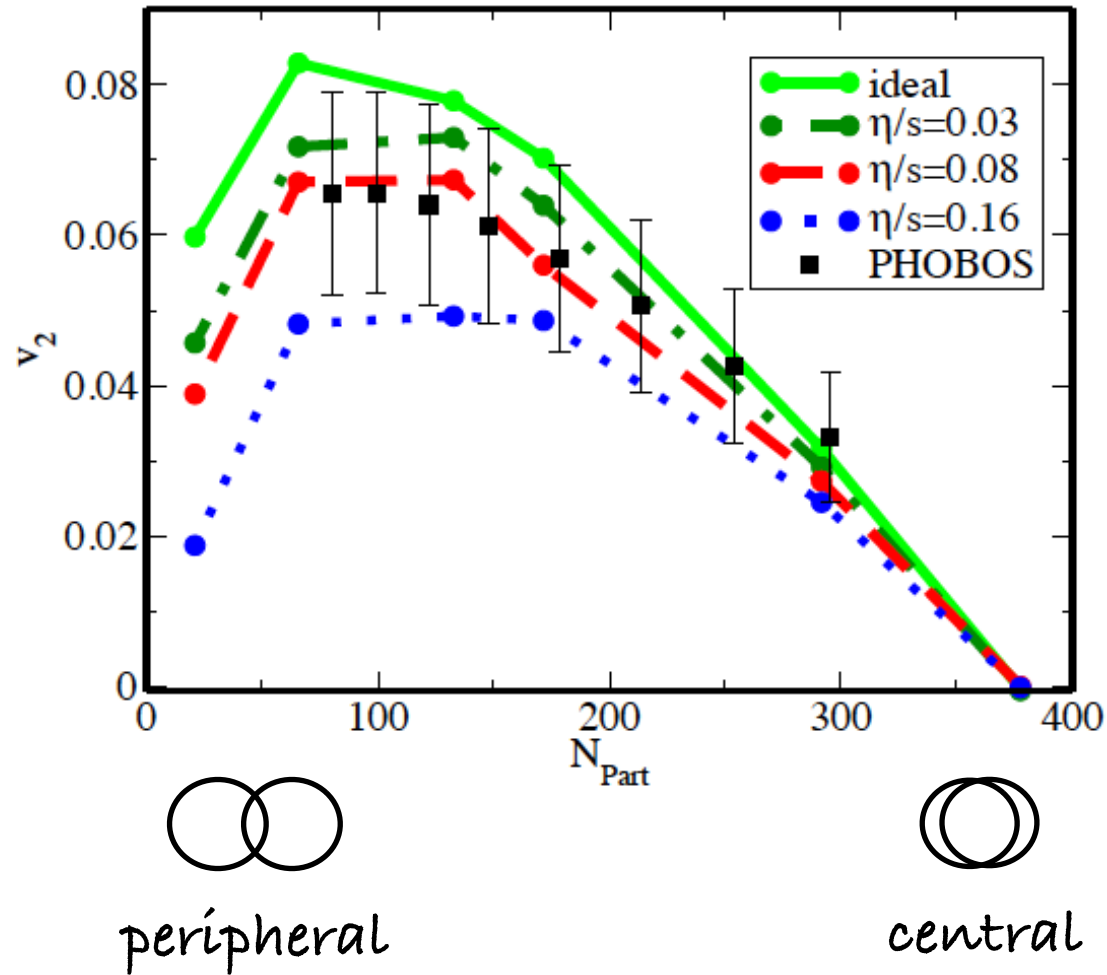


peripheral



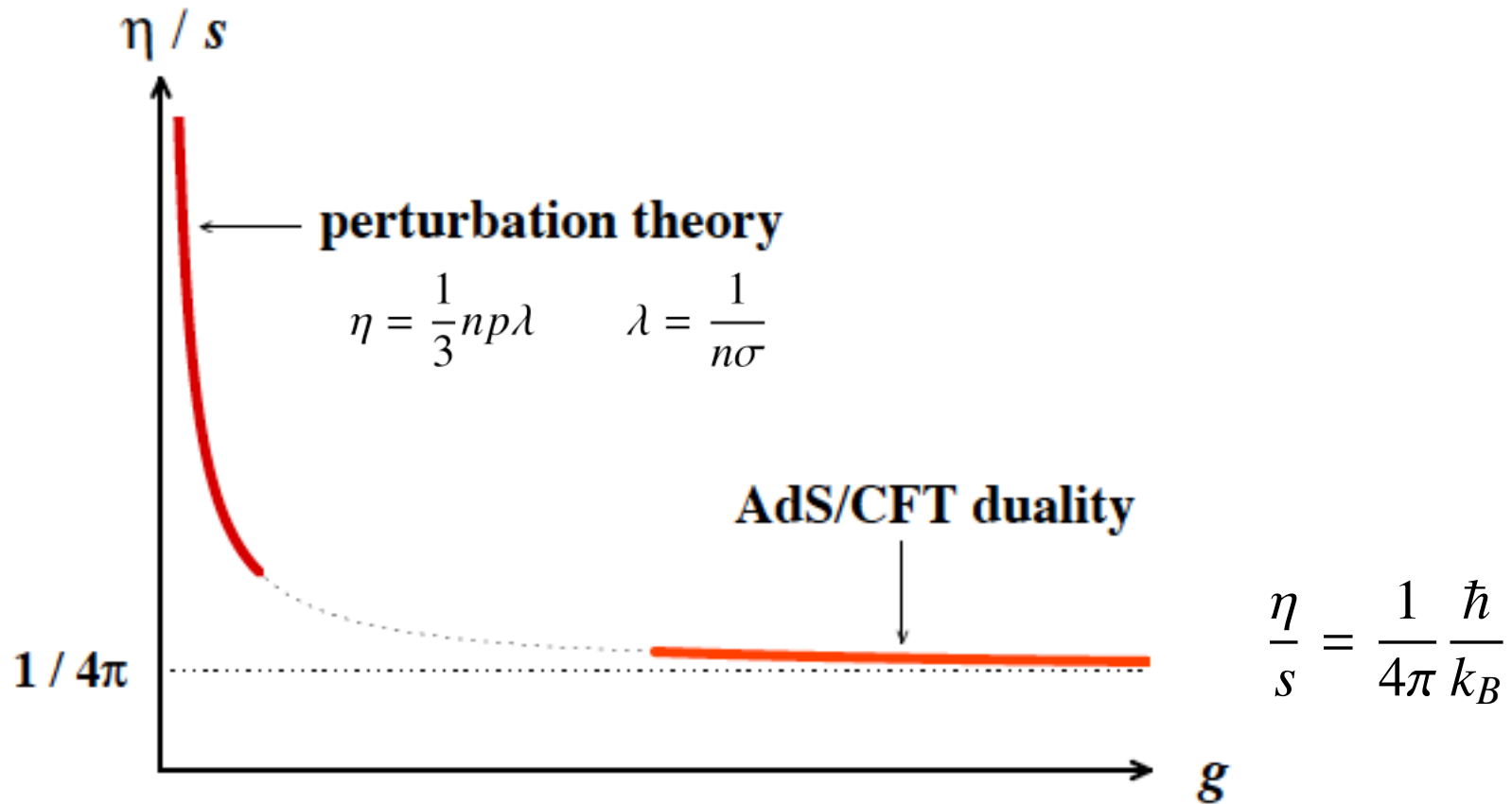
central

The low viscosity of the quark-gluon plasma



(Luzum, Romatschke, 2007)

viscosity at weak and strong coupling



(Policastro, Son, Starinets, 2001)

Theoretical puzzle

Small η/s and short equilibration time seem incompatible with weak coupling

However

- the coupling is not huge ($\alpha_s \sim 0.3 \div 0.4$)

- our present understanding of initial stages of heavy ion collisions is based on weak coupling (for asymptotically large nuclei and large energies)

$$Q_s^2 \approx \alpha_s \frac{xG(x, Q^2)}{\pi R^2}$$

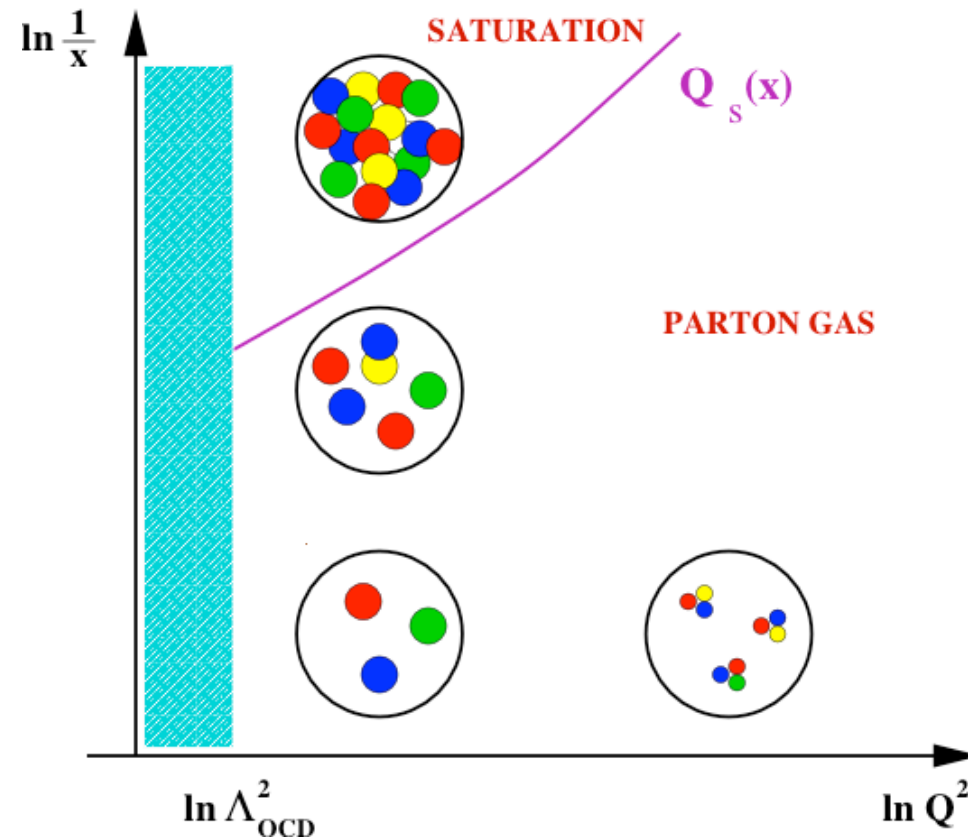
High density partonic systems

$$Q_s^2 \approx \alpha_s \frac{xG(x, Q^2)}{\pi R^2}$$

At saturation,
occupation
numbers are large

$$\frac{xG(x, Q^2)}{\pi R^2 Q_s^2} \sim \frac{1}{\alpha_s}$$

Description in terms of
classical fields possible
(CGC, etc)



The overpopulated quark-gluon plasma

Initial conditions $(t_0 \sim 1/Q_s)$

$$\epsilon_0 = \epsilon(\tau = Q_s^{-1}) \sim \frac{Q_s^4}{\alpha_s} \quad n_0 = n(\tau = Q_s^{-1}) \sim \frac{Q_s^3}{\alpha_s} \quad \epsilon_0/n_0 \sim Q_s$$

Overpopulation parameter $n_0 \epsilon_0^{-3/4} \sim 1/\alpha_s^{1/4}$

In equilibrated quark-gluon plasma

$$\epsilon_{\text{eq}} \sim T^4 \quad n_{\text{eq}} \sim T^3 \quad n_{\text{eq}} \epsilon_{\text{eq}}^{-3/4} \sim 1$$

Mismatch by a large factor (at weak coupling) $\alpha_s^{-1/4}$

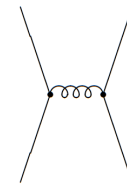
Formation of a Bose-Einstein condensate?

As a result of their interactions, gluons acquire a 'mass', and can condense.

Evidences for this phenomenon in classical scalar field theories (Epelbaum, Gelis, NPA 872 (2011) 210). Non abelian gauge theories ?

Note: when $f \sim 1/\alpha_s$ all dependence on the coupling constant disappears from kinetic equations

$$\partial_t f = C[f] \sim \alpha^2 f^3 \sim \frac{1}{\alpha}$$



(J.-P. B, F. Gelis, J. Liao, L. McLerran and R. Venugopalan, arXiv:1107.5296)

Conclusions

- the field of ultra-relativistic heavy ion collisions is a very rich one, it addresses fundamental issues (hot and dense QCD matter, high density partonic systems, etc)
- exciting developments in recent years, both experimentally and theoretically (RHIC, LHC; CGC, AdS/CFT, etc), many open questions/puzzles (weak vs strong coupling, thermalization, etc)
- future of the field looks bright, with many facilities allowing for such studies: LHC, RHIC2, FAIR, NICA