

Models of Elastic Diffractive Scattering to Falsify at the LHC

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The 14th EDS Blois Workshop

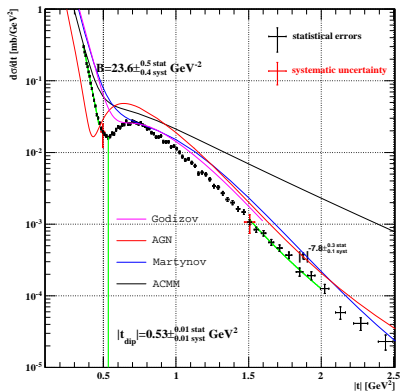
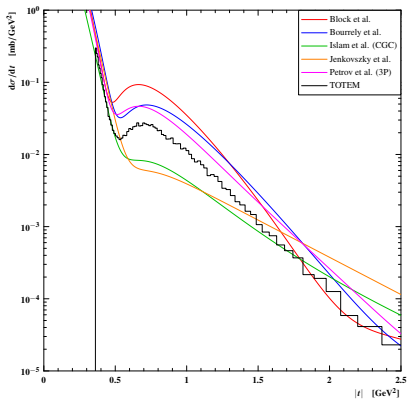
December 15-21, 2011, Qui Nhon, Vietnam

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- M.M. Islam, R.J. Luddy, and A.V. Prokudin, Int. J. Mod. Phys. A 21 (2006) 1
- E. Martynov, Phys. Rev. D 76 (2007) 074030
- R.F. Avila, P. Gauron, and B. Nicolescu, Eur. Phys. J. C 49 (2007) 581
- E. Martynov and B. Nicolescu, Eur. Phys. J. C 56 (2008) 57
- C. Flensburg, G. Gustafson, and L. Lönnblad, Eur. Phys. J. C 60 (2009) 233
- P. Brogueira and J. Dias de Deus, J. Phys. J 37 (2010) 075006
- M.M. Block and F. Halzen, Phys. Rev. D 83 (2011) 077901
- L.L. Jenkovszky, A.I. Lengyel, and D.I. Lontkovskyi,
Int. J. Mod. Phys. A 26 (2011) 4755
- E. Gotsman, E. Levin, and U. Maor, Eur. Phys. J. C 71 (2011) 1553
- M.G. Ryskin, A.D. Martin, and V.A. Khoze, Eur. Phys. J. C 71 (2011) 1617
- S. Ostapchenko, Phys. Rev. D 83 (2011) 014018
- D.A. Fagundes, E.G.S. Luna, M.J. Menon, and A.A. Natale,
arXiv: 1108.1206 [hep-ph]
- A. Godizov, Phys. Lett. B 703 (2011) 331
- A. Donnachie and P.V. Landshoff, arXiv: 1112.2485 [hep-ph]

Models vs. TOTEM. The pp Differential Cross-Section

The TOTEM Collaboration, Europhys.Lett. 95 (2011) 41001



Scattering Amplitude, Born Term (“Eikonal”) and Regge Trajectories

$$T_{12 \rightarrow 12}(s, t) = 4\pi s \int_0^\infty db^2 J_0(b\sqrt{-t}) \frac{e^{2i\delta_{12 \rightarrow 12}(s, b)} - 1}{2i},$$

$$\begin{aligned} \delta_{12 \rightarrow 12}(s, b) &= \frac{1}{16\pi s} \int_0^\infty d(-t) J_0(b\sqrt{-t}) \delta_{12 \rightarrow 12}(s, t) = \\ &= \frac{1}{16\pi s} \int_0^\infty d(-t) J_0(b\sqrt{-t}) \times \\ &\times \left\{ \sum_n \left(i + \operatorname{tg} \frac{\pi(\alpha_n^+(t) - 1)}{2} \right) \Gamma_n^{(1)+}(t) \Gamma_n^{(2)+}(t) s^{\alpha_n^+(t)} \mp \right. \\ &\left. \mp \sum_n \left(i - \operatorname{ctg} \frac{\pi(\alpha_n^-(t) - 1)}{2} \right) \Gamma_n^{(1)-}(t) \Gamma_n^{(2)-}(t) s^{\alpha_n^-(t)} \right\}. \end{aligned}$$

Regge Trajectories and QCD. The BFKL Approach

J. Kwiecinski, Phys. Rev. D **26** (1982) 3293:

$$\alpha_{\bar{q}q}(t) = \sqrt{\frac{8}{3\pi}\alpha_s(\sqrt{-t})} + o(\alpha_s^{1/2}(\sqrt{-t}))$$

R. Kirschner and L.N. Lipatov, Z. Phys. C **45** (1990) 477:

$$\alpha_{gg}(t) = 1 + \frac{12 \ln 2}{\pi}\alpha_s(\sqrt{-t}) + o(\alpha_s(\sqrt{-t}))$$

V.S. Fadin and L.N. Lipatov, Phys. Lett. B **429** (1998) 127,

M. Ciafaloni and G. Camici, Phys. Lett. B **430** (1998) 349:

$$\alpha_{gg}(0) = 1 + \frac{12 \ln 2}{\pi}\alpha_s(\mu) \left(1 - \frac{20}{\pi}\alpha_s(\mu)\right) + o(\alpha_s^2(\mu))$$

C. Lovelace, Nucl. Phys. B **95** (1975) 12:

$$(\alpha_{\phi\phi}^{(k)}(0) + 1)(\alpha_{\phi\phi}^{(k)}(0) + 2)(\alpha_{\phi\phi}^{(k)}(0) + 3) = \frac{16}{3(2k + 1)}$$

D. Heckathorn, Phys. Rev. D **18** (1978) 1286:

$$\alpha_{gg}^{(k)}(0)(\alpha_{gg}^{(k)}(0) + 1)(\alpha_{gg}^{(k)}(0) + 2) = \frac{24N_c}{(2k + 1)(11N_c - 2n_f)}$$

A. Godizov, Phys. Rev. D **81** (2010) 065009:

$$\alpha_{\bar{q}q}^{(k)}(0) = \frac{9(N_c^2 - 1)}{(2k + 1)N_c(11N_c - 2n_f)} - 1$$

Single-reggeon exchanges + eikonalization:

C. Bourrely, J. Soffer, and T.T. Wu, Eur. Phys. J. C 28 (2003) 97

V.A. Petrov and A.V. Prokudin, Eur. Phys. J. C 23 (2002) 135

A. Godizov, Phys. Lett. B 703 (2011) 331

Models without eikonalization:

A. Donnachie and P.V. Landshoff, arXiv: 1112.2485 [hep-ph]

P. Desgrolard, M. Giffon, and E. Martynov, Eur. Phys. J. C 18 (2000) 359

E. Martynov, Phys. Rev. D 76 (2007) 074030

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R.F. Avila, P. Gauron, and B. Nicolescu, Eur. Phys. J. C 49 (2007) 581

E. Martynov and B. Nicolescu, Eur. Phys. J. C 56 (2008) 57

Models Using Reggeon Field Theory

E. Gotsman, E. Levin, U. Maor, Eur. Phys. J. C 71 (2011) 1553

M.G. Ryskin, A.D. Martin, V.A. Khoze, Eur. Phys. J. C 71 (2011) 1617

S. Ostapchenko, Phys. Rev. D 83 (2011) 014018

(a) Elastic amplitude

$$\text{Im} A_{\text{el}} = \overline{\text{---} \text{---}} = 1 - e^{-\Omega/2} = \sum_{n=1}^{\infty} \overline{\text{---} \text{---} \text{---} \text{---} \text{---} \text{---}} \Omega/2$$

(b) Inclusion of low-mass dissociation



$$\text{Im} A_{ik} = \overline{\text{---} \text{---}}^i_k = 1 - e^{-\Omega_{ik}/2} = \sum \overline{\text{---} \text{---} \text{---} \text{---} \text{---} \text{---}} \Omega_{ik}/2$$

(c) Inclusion of high-mass dissociation

$$\Omega_{ik} = \overline{\text{---} \text{---}}^i_k + \overline{\text{---} \text{---}}^i_k \} M + \overline{\text{---} \text{---}}^i_k + \dots + \overline{\text{---} \text{---}}^i_k + \dots$$

Models not appealing to QCD:

R.F. Avila, S.D. Campos, M.J. Menon, and J. Montanha,
Eur. Phys. J. C 47 (2006) 171

P. Brogueira and J. Dias de Deus, J. Phys. J 37 (2010) 075006

“QCD-inspired” models:

M.M. Islam, R.J. Luddy, and A.V. Prokudin,
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arXiv: 1108.1206 [hep-ph]

Models vs. TOTEM. The Total pp Cross-Section

The TOTEM Collaboration, Europhys.Lett. 96 (2011) 21002:

$$\sigma_{tot}^{pp}(7 \text{ TeV}) = (98.3 \pm 0.2^{stat} \pm 2.8^{syst}) \text{ mb}$$

The Model	$\sigma_{tot}^{pp}(7 \text{ TeV})$, mb
P. Desgrolard, M. Giffon, E. Martynov, Eur. Phys. J. C 18 (2000) 359	95
V.A. Petrov, A.V. Prokudin, Eur. Phys. J. C 23 (2002) 135	97 ± 4
C. Bourrely, J. Soffer, T.T. Wu, Eur. Phys. J. C 28 (2003) 97	93
R.F. Avila, S.D. Campos, M.J. Menon, J. Montanha, Eur. Phys. J. C 47 (2006) 171	94
M.M. Islam, R.J. Luddy, A.V. Prokudin, Int. J. Mod. Phys. A 21 (2006) 1	97.5
E. Martynov, Phys. Rev. D 76 (2007) 074030	91
R.F. Avila, P. Gauron, B. Nicolescu, Eur. Phys. J. C 49 (2007) 581	108
E. Martynov, B. Nicolescu, Eur. Phys. J. C 56 (2008) 57	95
C. Flensburg, G. Gustafson, L. Lönnblad, Eur. Phys. J. C 60 (2009) 233	98 ± 9

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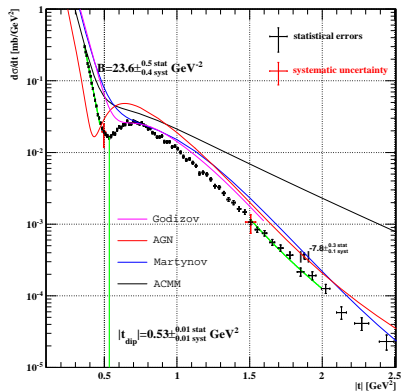
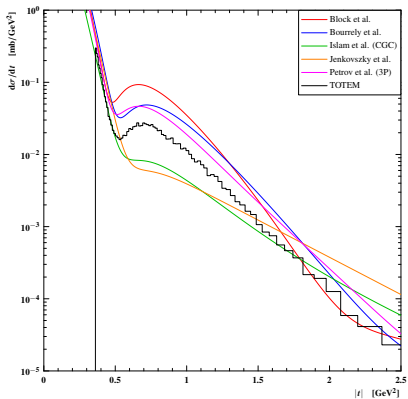
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M.M. Block, F. Halzen, Phys. Rev. D 83 (2011) 077901	95.5 ± 1
L.L. Jenkovszky, A.I. Lengyel, D.I. Lontkovskyi, Int. J. Mod. Phys. A 26 (2011) 4755	98 ± 1
E. Gotsman, E. Levin, U. Maor, Eur. Phys. J. C 71 (2011) 1553	91
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Conclusion

We need deeper interrelation of phenomenological models with QCD. This requires developing powerful non-perturbative QCD techniques which should allow to deal with diffractive (large distance) domain of strong interaction.

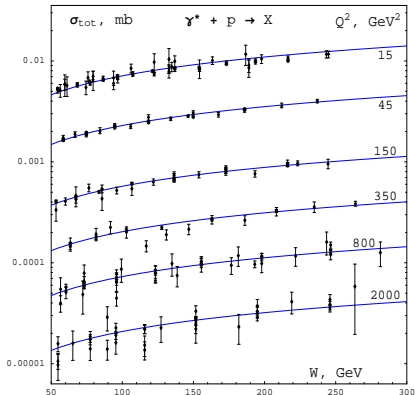
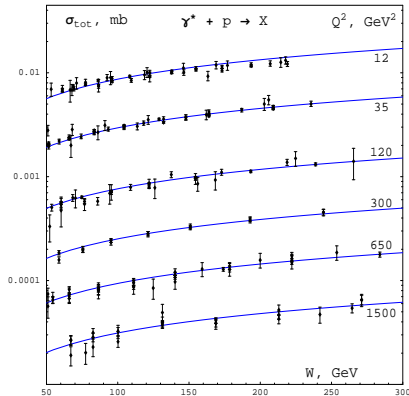
Thank You for Attention!

Extraction of the Pomeron Intercept from the DIS Data

At $50 \text{ GeV} < W < 300 \text{ GeV}$ and $25 \text{ GeV}^2 < Q^2 \ll W^2$ the $\gamma^* p$ total cross-sections can be well-described by simple formula

$$\sigma_{tot}^{\gamma^* p}(W^2, Q^2) \approx \beta(Q^2) W^{2\delta}$$

with $\delta \approx 0.31$



Extraction of the Pomeron Intercept from the DIS Data

This universal parameter can be associated with the intercept of the pomeron trajectory:

$$\alpha_P(0) - 1 = \delta \approx 0.31$$

