# Angular Correlations of Hadrons Measured at the LHC

Rencontres du Vietnam

Frontiers of QCD: From Puzzles to Discoveries

> December 15-21, 2011 Qui Nhon, Vietnam

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20-12-2011





# Angular Correlations

#### ATLAS-CONF-2011-074

$$C(\Delta\phi\Delta\eta) \equiv \frac{N_{\rm mixed}}{N_{\rm same}} \frac{{\rm d}^2N_{\rm same}/{\rm d}\Delta\phi{\rm d}\Delta\eta}{{\rm d}^2N_{\rm mixed}/{\rm d}\Delta\phi{\rm d}\Delta\eta}$$

Contributions to the two-particle  $\Delta \Phi$ ,  $\Delta \eta$  angular correlation come from anisotropic flow, jets, resonances, HBT, etc



see also CMS HIN-11-006

# Angular Correlations



For very peripheral collisions or when triggered with a high-pt particle the dominant contribution to the two particle angular correlations is due to jet-correlations More central heavy-ion collisions look very very different! anisotropic flow







initial spatial geometry not a smooth almond event-by-event (for which all odd harmonics and sin  $n(\Phi-\Psi_R)$  are zero due to symmetry) may give rise to higher odd harmonics and symmetry planes in momentum space (detailed probes of initial conditions)

### measure anisotropic flow

$$\langle v_n \rangle = \langle \langle e^{in(\phi_1 - \Psi_n)} \rangle \rangle$$

 since the common symmetry planes cannot be measured event-by-event, we measure quantities which do not depend on their orientation: two and multi-particle azimuthal correlations

$$\langle \langle e^{in(\phi_1 - \phi_2)} \rangle \rangle = \langle \langle e^{in(\phi_1 - \Psi_n - (\phi_2 - \Psi_n))} \rangle \rangle$$
  
=  $\langle \langle e^{in(\phi_1 - \Psi_n)} \rangle \langle e^{-in(\phi_2 - \Psi_n)} \rangle \rangle$   
=  $\langle v_n^2 \rangle$ 

 assuming that <u>only</u> correlations with the symmetry planes are present - not always a very good assumption (contributions from jets, resonances, etc)

# Angular Correlations



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Two particle azimuthal correlations can be described efficiently with the first 6  $v_n$  coefficients

# Can we isolate the flow?

• if nonflow is negligible flow "factorizes"  $\rightarrow$   $\leftarrow$ 

$$\langle \langle e^{in(\phi_1 - \phi_2)} \rangle \rangle = \langle \langle e^{in(\phi_1 - \Psi_n - (\phi_2 - \Psi_n))} \rangle \rangle$$
  
=  $\langle \langle e^{in(\phi_1 - \Psi_n)} \rangle \langle e^{-in(\phi_2 - \Psi_n)} \rangle \rangle$   
=  $\langle v_n^2 \rangle$ 

- test with particles separated in rapidity
- test with particles separated in pt
- flow is a collective effect
  - multi-particle correlations
    - Lee-Yang Zeroes, cumulants, q-vectors, etc

#### does it factorize?



ALICE arXiv:1109.2501

- yes it does (to a large extent for more central collisions)
- how large is the flow where factorization "breaks"?
  - to quantify that one needs other techniques (multi-particle)

# multi-particle correlations

 for detectors with uniform acceptance the 2<sup>nd</sup> and 4<sup>th</sup> order cumulant are given by:

> Borghini, Dihn and Ollitrault, PRC 64, 054901 (2001)

$$c_n\{2\} \equiv \langle \langle e^{in(\phi_1 - \phi_2)} \rangle \rangle = v_n^2 + \delta_2$$

$$c_n\{4\} \equiv \langle \langle e^{in(\phi_1 + \phi_2 - \phi_3 - \phi_4)} \rangle \rangle - 2 \langle \langle e^{in(\phi_1 - \phi_2)} \rangle \rangle^2$$

$$= v_n^4 + 4v_n^2 \delta_2 + 2\delta_2^2 - 2(v_n^2 + \delta_2)^2$$

$$= -v_n^4$$

we got rid of two particle nonflow correlations! we can remove nonflow order by order

## v<sub>2</sub> from cumulants



cumulants show behavior as expected when correlations are dominated by collective flow

### v<sub>2</sub> from multi-particle correlations



behavior as expected when correlations are dominated by collective flow (difference between 2 and multi-particle estimates mainly due to e-by-e fluctuations in the flow

The Perfect Liquid



The flow increases about 30%. The system produced at the LHC behaves as a very low viscosity fluid, constrains dependence of  $\eta$ /s versus temperature

## v<sub>2</sub> as function of p<sub>t</sub>



Elliptic flow as function of transverse momentum does not change much from RHIC to LHC energies, can we understand that?





## The Perfect Liquid?

in calculations the RHIC  $v_2$  results are close to the ideal hydrodynamical limit.

these calculations place
an upper limit on η/s
which is smaller than ~
4 x AdS/CFT bound



Based on R. Lacey et al., PRL 98 (2007) 092301





u<sub>1</sub> > u<sub>2</sub> > u<sub>3</sub> shear viscosity will make them equal and destroy the elliptic flow v<sub>2</sub> higher harmonics represent smaller differences which get destroyed more easily, and which, if measurable, makes them more sensitive probes to η/s

## Shear Viscosity

#### Music, Sangyong Jeon



#### initial conditions

#### ideal hydro $\eta/s=0$ viscous hydro $\eta/s=0.16$



Larger η/s clearly smoothes the distributions and suppresses the higher harmonics (e.g. v<sub>3</sub>)

Hydro: Alver, Gombeaud, Luzum & Ollitrault, Phys. Rev. C82 (2010) 17

# the v<sub>n</sub>'s

The  $v_3$  with respect to the reaction plane determined in the ZDC and with the  $v_2$ participant plane is consistent with zero as expected if  $v_3$  is due to fluctuations of the initial eccentricity

The  $v_3\{2\}$  is about two times larger than  $v_3\{4\}$  which is also consistent with expectations based on initial eccentricity fluctuations



ALICE Collaboration, arXiv:1105.3865 PRL 107 (2011) 032301

We observe significant  $v_3$  and  $v_4$  which compared to  $v_2$  has a different centrality dependence (already strong constrain for  $\eta/s$ )

# the v<sub>n</sub>'s



#### ATLAS-CONF-2011-074

# Elliptic and Triangular Flow



Qui, Shen and Heinz, arXiv:1110.3033

The centrality dependence and magnitude are better described by predictions using MC Glauber with  $\eta/s=0.08$ 

# Flow Analysis Methods

flow analysis methods have different sensitivity to nonflow and fluctuations

> Borghini, Dihn and Ollitrault, PRC 64, 054901 (2001) Bilandzic, Snellings and Voloshin, PRC 83, 044913 (2011)

 $v_n^2 \{2\} = \bar{v}_n^2 + \sigma_v^2 + \delta$  $v_n^2 \{4\} = \bar{v}_n^2 - \sigma_v^2$  $v_n^2 \{6\} = \bar{v}_n^2 - \sigma_v^2$  $v_n^2 \{8\} = \bar{v}_n^2 - \sigma_v^2$ 

excellent opportunity to study flow fluctuations and from these get a handle on initial conditions!

#### v<sub>2</sub> versus centrality



Two particle  $v_2$  estimates depend on  $\Delta \eta$ Higher order cumulant  $v_2$ estimates are consistent within uncertainties Two particle  $v_2$  estimates are corrected for nonflow based on HIJING The estimated nonflow correction for  $\Delta \eta > 1$  is included in the systematic uncertainty

### v<sub>2</sub> Fluctuations



Fluctuations are significant and for more central collisions not in agreement with the eccentricity fluctuations in MC-Glauber and MC-KLN CGC

## Conclusions

- Anisotropic flow measurements provided strong constraints on the properties of hot and dense matter produced at RHIC and LHC energies and have led to the new paradigm of the QGP as the so called perfect liquid
  - At the LHC we observe even stronger flow than at RHIC which is expected for almost perfect fluid behavior
- The first measurements of v<sub>3</sub> and higher v<sub>n</sub>'s have recently been made at RHIC and at the LHC and indicate that these flow coefficients behave as expected from fluctuations of the initial spatial eccentricity (geometry!) and a created system which has a small η/s
  - provide new strong experimental constraints on η/s and initial conditions

# Thanks

## v<sub>2</sub> scaling?

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# v<sub>3</sub>(m,p<sub>t</sub>) and the scaling



The behavior of  $v_3$  as function of  $p_t$  for pions, kaons and protons shows the same features as observed for  $v_2$ (the mass splitting, the crossing of the pions with protons at intermediate  $p_t$ )

# Geometry and Harmonics



For central collisions at intermediate pt the higher harmonics v3 and v4 cross v2 and become the dominant harmonics Why do they cross??

For more central collisions this occurs already at lower pt



#### Flow Fluctuations

when (2-particle) nonflow is corrected for or negligible!

in limit of "small" (not necessarily Gaussian) fluctuations

 $v_n^2 \{2\} = \bar{v}_n^2 + \sigma_v^2$  $v_n^2 \{4\} = \bar{v}_n^2 - \sigma_v^2$  $v_n^2 \{2\} + v_n^2 \{4\} = 2\bar{v}_n^2$  $v_n^2 \{2\} - v_n^2 \{4\} = 2\sigma_v^2$ 

$$v_n\{4\} = 0$$
$$v_n\{2\} = \frac{2}{\sqrt{\pi}}\bar{v}_n$$

### Flow Fluctuations

Example: input  $v_2 = 0.05 + 0.02$  (Gausian), M = 500,  $N = 1 \times 10^6$ 



Gaussian fluctuation behave as predicted also for Lee Yang Zeroes and fitting Q distribution (more on that later)